



# **Institute Laboratory Assessment**

## **Interim Review**

**Advanced Meson Science Laboratory**

**MASAHIKO IWASAKI**

The Chief Scientist

February 26, 2010

RIKEN Nishina Center  
for Accelerator-Based Science



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# 1. Reviewers List

| Name                  | Affiliation   |
|-----------------------|---|
| Tadafumi Kishimoto    | Director, Research Center for Nuclear Physics,<br>Osaka University, Japan                 |
| Tomofumi Nagae        | Professor, Graduate School of Science,<br>Kyoto University, Japan                         |
| Kusuo Nishiyama       | Honorary Professor,<br>Institute of Materials Structure Science<br>KEK, Japan             |
| Robert F. Kiefl       | Professor, Department of Physics and Astronomy,<br>University of British Columbia, Canada |
| Jean-Michel Poutissou | Associate Director,<br>TRIUMF, Canada   |

## 2. Review Program

Laboratory Advanced Meson Science Laboratory  
(Chief Scientist: Masahiko Iwasaki)

Date February 26<sup>th</sup> 2010, 9:30~17:00

Venue RIKEN Wako Institute

| Time            | Content   | Other details               |
|-----------------|---|-----------------------------|
| 09 : 30~10 : 00 | Introduction to RIKEN<br>(Dr. En'yo, the Director of Nishina Center for Accelerator-Based Science ) | Wako Campus<br>Nishina Hall |
| 10 : 00~11 : 30 | Presentation by the Chief Scientist   | Wako Campus<br>Nishina Hall |
| 11 : 30~12 : 10 | Question and Answer   |                             |
| 12 : 10~13 : 00 | Lunch (Reviewers only)  | Wako Campus                 |
| 13 : 00~14 : 00 | Discussion by the review members with the Chief Scientist and RNC Director, Executive Directors     | Wako Campus<br>Nishina Hall |
| 14 : 00~14 : 50 | Interview by the reviewer with the laboratory staffs  | Wako Campus<br>Nishina Hall |
| 14 : 50~15 : 00 | Coffee Break  |                             |
| 15 : 00~16 : 40 | Closed discussion by the reviewers  | Wako Campus<br>Nishina Hall |
| 16 : 40~17 : 00 | General briefing to the RNC Director, Board of Chief Scientist Assembly                             |                             |
| 17 : 30~        | Dinner hosted by Dr. En'yo  |                             |



RIKEN Wako Campus

## 3. Review Materials

### 3.1. Introductory Remarks and Acknowledgements

Advanced Meson Science Laboratory is presently belonging to RIKEN Nishina Center for Accelerator-Based Science (RNC). We have two major research subjects. One is the study of meson – nucleon / nucleus interaction for fundamental hadron physics using mesons as a probe. The other is muon science for more generic study using muons, covering wide area such as condensed matter physics with  $\mu$ SR and nuclear-atomic related study on  $\mu$ CF. We have also several accelerator-related studies, such as Mössbauer study as a part of the condensed matter physics.

As described above, we are covering wide area of accelerator-based studies. The diversity of research fields is originated from the establishment of our laboratory. Our laboratory started as Muon Science Laboratory, initiated by former chief scientist Kanetada (Ken) Nagamine, who established RIKEN-RAL pulsed-muon facility at Rutherford Appleton Laboratory (UK) back to 1990. This institution-based international cooperation to the experimental facility, including its construction and operation beyond the national boundary, is quite unique. The facility, named as RIKEN-RAL Branch (RRB), is operated as a research center of muon science, which is open to public through experimental proposals. The experimental program advisory committees (PAC) are funded both in Japan (hosted by RIKEN) and in UK (by RAL) sharing the beam time of this facility. Since then, this activity is taken as one of remarkable success to both Japan and UK.

In 2002, at the initial period of the second term of Japan-UK research contract (10 years from 2000), Masahiko Iwasaki was assigned to be a chief scientist of this laboratory, succeeding in the RRB operation mission after Nagamine's retirement from RIKEN. On the other hand, a new-generation pulsed-muon facility (MUSE) construction was approved as one of the important components of a J-PARC accelerator complex (Tokai, Japan), and KEK muon group was assigned as the construction team. This gave us a strong motivation to extend, enrich and deepen our research field to new direction, without limiting ourselves only to the RRB operation. Meanwhile, our group was requested to maintain and enhance our activity at RAL from the Japan muon / meson science community. Thus we motivated ourselves that the muon science activity at RIKEN-RAL remain one of the main subjects of this laboratory, not only by a mission.

When Iwasaki was appointed as a chief scientist, he was also requested from RIKEN to open new research program initiated by himself. In RIKEN, there were two other large-scale physics programs. One is new radioisotope beam factory (RIBF) covering low energy nuclear physics focusing on element genesis. At that time, four chief scientists were already committed to the RIBF project. The other is the spin physics using RHIC at BNL (RBRC) covering very high-energy nuclear physics focusing on origins of proton spin, property of QGP phase, and theoretical researches related to RBRC. The RBRC has been

operated under a representative chief scientist, Hideto En'yo. Thus, it is required to select a subject orthogonal to RIBF and/or RBRC related ones.

Before the appointment, Iwasaki already initiated and conducted successfully several hadron experiments at the KEK 12 GeV proton synchrotron (PS). Thus, it was natural to choose high-energy nuclear (hadron) physics as an expansion of the laboratory's identity, focusing on strangeness-related hadron physics as a new research field for our laboratory. In this framework, it is to be underlined that we succeeded to attract researchers coming from all over the world, forming international collaborations, with the RIKEN group having leading role. We are, as well, actively participating to international collaborations in other laboratories, such as KEK, GSI and LNF-INFN. Our laboratory has gained an international dimension in both directions: attracting scientists in Japan, and allowing Japanese scientists to take part to experiments performed elsewhere. Presently we are preparing several new experiments using exotic beams of kaon or anti-proton to study the meson property in nuclei at J-PARC hadron hall. To study pion-nuclear interaction or pion property in nuclear media, RIBF is another ideal place, where we are also preparing new high-resolution experiment. Another important activity, this laboratory has, is forming new scientists: we have followed many young students in their Master and Ph. D. theses, with excellent outcome on both sides.

When RNC was established in 2006, the former director, Yasushige Yano, invited all the accelerator-based researchers to be a member of the new Center, defining its mission to endorse and promote all the accelerator-based researches without limiting only to RIBF. Under this spirit, "Nishina Center for Accelerator-Based Science" is named to represent our activities and ourselves. As a result, RNC's mission is full of diversity of researches at RIBF, RBRC, RRB, J-PARC, and more.

Let us briefly describe the future direction of our laboratory. As a part of internal discussions of the institutional based conceptual panning of RNC, where we belong to, we are proposing two outstanding projects. One is hadron physics project to explore the "origin of the hadron (matter) mass", which requires substantial RNC's institutional-based contribution to J-PARC Hadron hall. To realize the project we have been discussing with Radiation laboratory of RNC. The other is the muon science project to realize new muon anomalous magnetic moment ( $g-2$ ) measurement to confirm the physics beyond standard model, using ultra-cold muon source what we have developing at RIKEN-RAL. To realize the ultra-cold muon source, we have been collaborating with KEK and TRIUMF.

This article covers the scientific achievements of Advanced Meson Science Laboratory from 2002, when Masahiko Iwasaki was appointed as a chief scientist, to present. He wishes to address his gratitude to all the contributions to prepare this article by all the members of Advanced Meson Science Laboratory, including who have left the laboratory. Special thanks should be given to Kanetada Nagamine the former chief scientist, Hideto En'yo the director of the RNC, and former director Yasushige Yano. He is also grateful to his research colleagues in RIKEN and RAL, as well as other collaborating institutions over the

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world including students for their continuous and extensive research works. All the helps and encouragements from administrative staffs in RIKEN, from the top executives to the secretaries of the laboratory, are sincerely appreciated. Admixture of those contributions to the chief scientist and Advanced Meson Science Laboratory makes these 8 years to be fully fruitful and enjoyable.

February 8<sup>th</sup>, 2010

Masahiko Iwasaki

## 3.2. Research Topics

### 3.2.1. Hadron Physics

Property of particles in matter / media attracts interest in many ways. It is believed that the property of particles is not independent from the surroundings through the interaction between the two. Present scenario “how the mass of the basic (elementary) particle is formed” is based on the idea that the vacuum is also a kind of matter where many interesting processes are taking place, as for example the higgs condensate. This condensation is believed to have been realized during the cooling down of the universe at around  $T \sim 100$  GeV after the Big Bang. Therefore, to detect higgs is one of the major goals of the LHC project at CERN.

To explain the origin of the mass of the hadrons, however, this higgs mechanism is not sufficient to account for the so called constituent quark mass  $m_q \sim 300$  MeV/c<sup>2</sup> and the baryon formation. To explain that, one needs another phase transition of the vacuum called quark-antiquark pair condensation at around  $T \sim 200$  MeV. Thus, the vacuum expectation value of  $\langle \bar{q}q \rangle$  is non-zero due to the spontaneous chiral symmetry breaking of the vacuum, and this  $\langle \bar{q}q \rangle$ - condensation is the major source of masses of low lying hadrons such as protons, neutrons, pions, etc, and, consequently of the matter and...of ourselves.

The present scenario of the mass formation of the particles depends on how particles interact with surrounding space where higgs and quark-antiquark pairs are condensed. Thus, in the hadron sector, the in-medium particle properties are fundamentally related to the chiral symmetry breaking mechanism. The  $\langle \bar{q}q \rangle$  expectation value (chiral order parameter) is a function of temperature and chemical potential (density). Therefore, there is currently great experimental interest to study the effect of chiral symmetry breaking and its partial restoration in the nuclear media.

In a nucleus, it is known that mesons, especially pions, play important role to bind nucleons. The existence of the mesons was originally predicted by Yukawa as bosons, which form the nuclear field as a solution of the Klein-Goldon equation of the zero total energy, naamely ( $\sim \exp(-m_\pi r)/r$ ). In the standard model, the basic gauge particles of the strong interaction were replaced by gluons, but gluons cannot propagate directly between two nucleons in a nucleus due to the confinement nature, so the meson field approach is still valid in the phenomenological way to represent the nuclear interaction. In this manner, nuclei consist of nucleons and virtual meson fields. Naturally, nuclei are expected to be strongly absorptive space for mesons such as to form stable bound states in it. Therefore, experimental search for these states were untouched for long time. However, it is quite interesting subject how mesons behave and how their properties may change in the nuclear media.

Triggered by the series of recent experimental studies of the mesonic atom, the importance of the experimental search for the mesonic nuclear bound states was shed into new light. Thus, the extensive

experimental studies were started only very recently. It is still controversial whether there exist deeply-bound mesonic states further below atomic states (in other word, whether one can detect them experimentally) inside the nuclei. It is to be underlined that this type of studies are having an importance which extends from particle physics to the astrophysics, since part of the processes undergoing into the (neutron) stars might be explained by these type of processes and can only be studied by this approach.

Therefore, we have been approaching this difficult questions using various machines. In this interim period, we performed experiments at KEK 12 GeV-PS, GSI, DAFNE (LNF-INFN) and preparing several experiments at 50 GeV-PS of the J-PARC facility.

### 3.2.1.1. Individual topics of Hadron Physics

#### Kaonic Atoms

Let us first overview the recent studies of the  $\bar{K}N$  interaction by the x-ray measurement of the kaonic atoms. The  $\bar{K}N$  interaction can be studied by the level shifts and widths of the kaonic atom x rays, and one can study the chiral symmetry in nuclei and give constraints on chiral perturbation theory as it is the case in the pionic atom study. However, present kaonic atomic data are still insufficient for the understanding of the  $\bar{K}N$  interaction in detail.

In general, if the interaction is attractive, the effect of the strong interaction to the atomic level is much more complicated compared to the repulsive case. To demonstrate the situation, simple calculation result of the level shift of the kaonic hydrogen atom is plotted in Fig. 1, as a function of the real part of the potential between kaon and proton. The s-wave atomic levels are calculated by solving the Schrödinger equation in a Coulomb field and Yukawa potential as a local part,  $(V + iW)\exp(-r/\lambda)/r$ , where  $\lambda$  is the range parameter to be 1 fm for the simplicity, and  $V$  and  $w$  are the constants representing real and imaginary part of the local part. As shown in the figure, the atomic level shifts have beat pattern when the attractive force is strong enough compared to the imaginary part. It can be understood as interference between nuclear pole and atomic one. At the node, a bound state to the local potential is formed as a nuclear state, and the atomic level shift changes its sign and moves upward for more attractive interaction. This upward shift is not the result of the repulsion. The atomic ground state change its nature to the nuclear (local) one, and the wave function confined in a local potential beyond the node. After our previous work on kaonic hydrogen atom [1],  $\bar{K}N$  interaction

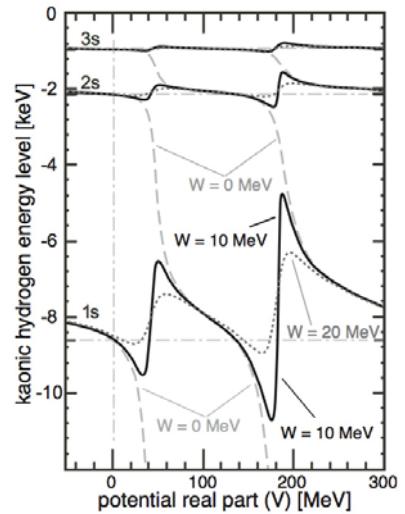


Fig. 1 A diagram of the level shift of the s-wave kaonic-hydrogen atom as a function of real part of the local potential. Three curves of the imaginary part,  $W = 0$  (dashed), 10 (solid) and 20 (dotted line), are plotted.

was confirmed to be strongly attractive.

There is another problem called “kaonic helium puzzle”. Calculated shift of the kaonic helium 2p level should be almost zero (well below 1 eV) while three data points are consistent and having an average value as large as  $\sim 40$  eV (upward). In the x-ray spectra, kaonic helium x-rays were clearly observed. Therefore, a possibility of existence of the deeply bound kaonic states in nuclei is discussed by Wycech as early as 1986 [2], based on the kaonic  $^4\text{He}$  atom data. This prediction triggers a new precision x-ray experiment of the kaonic helium atom at KEK, KEK PS-E570. We observed the kaonic helium x-ray transition very clearly with good resolution, which is achieved by using SDD as x-ray detectors. An x-ray spectrum is shown in Fig. 2. In the spectrum, the Ti and Ni fluorescence x-ray peaks were recorded simultaneously using the self-trigger mode of the SDD signal, and used as in-situ absolute energy-calibration sources to reduce the systematic error. As a result, the shift was consistent to be zero,

$$-\Delta E_{2p} = -(E_{obs}^{3d \rightarrow 2p} - E_{EM}^{3d \rightarrow 2p}) = -2 \pm 2(\text{stat.}) \pm 2(\text{syst.}) \text{ eV.}$$

So the “kaonic helium puzzle” was resolved.

Unfortunately, the atomic shift is sensitive only when the pole position due to the strong interaction is quite close to the atomic one. Thus, to explore the interaction inside the nucleus, one needs a systematic study. To have more direct information, we are presently preparing new experiment, J-PARC E17, using  $^3\text{He}$  as a target.

We have as well had an active role in experiments on kaonic atoms performed at the DAFNE accelerator at Frascati, being paer of the DEAR and SIDDHARTA Collaboration. As such, we performed measurements on kaonic nitrogen, kaonic hydrogen and kaonic helium [4-7]. Presently, we are actively participating to the upgrade of the SIDDHARTA experiments, towards measuring more exotic atoms (kaonic deuterium and heavier targets).

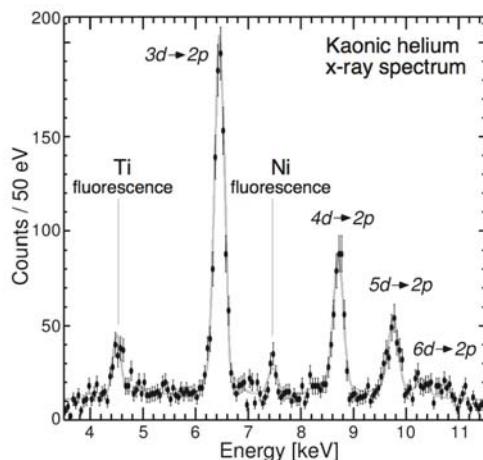


Fig. 2. An x-ray spectrum obtained by the most recent KEK experiment KEK PS-E570.

## References

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- [2]. S. Wycech: Nucl. Phys. A 450, 399C-402C (1986).
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## Pionic Atoms / Pions in Nuclei

We have made precision spectroscopy of pionic lead and tin atoms, and extracted information on the in-medium interaction between pion and nucleus, which leads to the exclusive quantitative evaluation of the chiral symmetry restoration in the nuclear matter [1]. Our collaboration, which mainly consists of the RIKEN and the University of Tokyo group, conducted the experiments starting from R&D to the precision spectroscopy. The experiment was carried out in GSI, Darmstadt, utilizing ( $d, {}^3\text{He}$ ) pion transfer reaction at the zero-momentum transfer kinematics to enhance pionic atom formation cross section.

Our first discovery was pionic 2p state in the lead 207 nucleus [2], where the negative pion is accommodated in a delicate balance between the Coulomb attraction and the strong repulsion. The pion is almost touching the nuclear surface, and in the past studies, such a deeply bound pionic state is regarded as too short-lived for observation as a distinct peak. Following the discovery, we have performed experiments to measure  $1s$  pionic lead 205 [3] and tin 115, 119, and 123 isotopes [1] as shown in Figure 1 for the case of pionic tin isotopes.

We have analyzed the experimental spectra elaborately and extracted in-medium isovector interaction between pion and nucleus. In combination with experimental information on the pionic hydrogen and

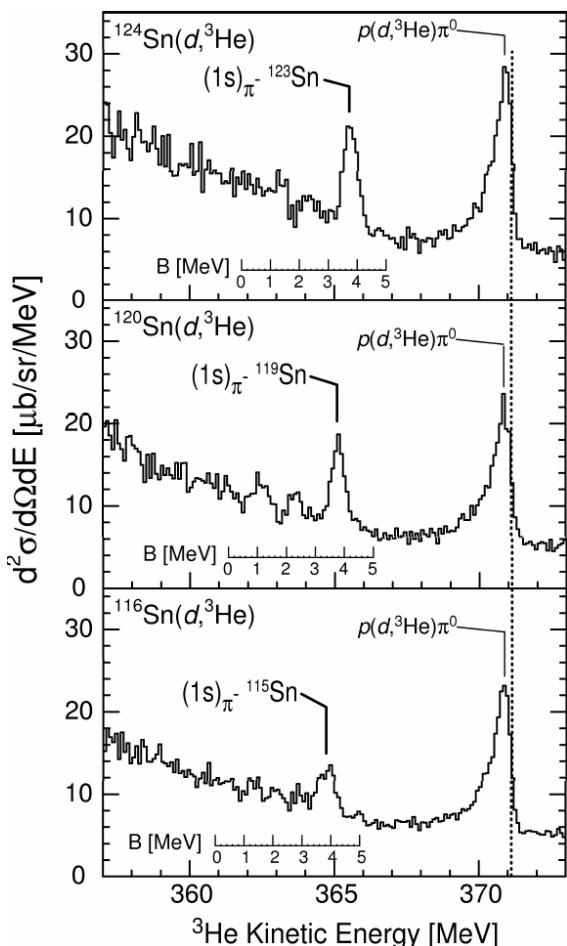


Figure 1 : Experimental spectra of pionic tin isotopes [1].

deuterium, which gives the interaction strength in vacuum, we have evaluated the in-medium interaction modification. The modification is originating in the partial restoration of the chiral symmetry in the nucleus, and we have quantitative result for the first time concluding the chiral order parameter reduction in the nuclear matter to be 33 %, which is consistent with the theoretical prediction of 30 % as shown in Figure 1.

With respect to the approach to the understanding of the chiral symmetry, one of the merits of pionic atom spectroscopy superior to others is its small ambiguity. Since pionic atom is a meta-stable quantum state, the ambiguities in the measurement are minimized in principle in contrast to many experiments where in-medium mesons are in motion.

Presently, we have been preparing for a sophisticated experimental setup of the pionic atom spectroscopy at the RIBF in the RIKEN. We expect about twice better experimental resolution with much smaller systematic errors [4].

## References

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## $\Lambda$ in nuclei

We have been extensively studying non-mesonic weak decay(NMWD) of  $\Lambda$  hypernuclei( $\Lambda+N \rightarrow N+N$ ). This NMWD is the very unique strangeness-changing baryon-baryon weak interaction process, which can occur only in nuclei. Concerning the ratio of two possible NMWD modes,  $\Gamma n/\Gamma p \equiv \Gamma(\Lambda +n \rightarrow n+n)/\Gamma(\Lambda +p \rightarrow n+p)$ , theoretical calculation based one one-pion exchange model predicts  $\Gamma n \ll \Gamma p$ , whereas recent experiments for  $A=5,12$  suggests large ratios close to unity. However these reported results have large errors of  $30 \sim 100\%$ . Up to now, most of the experiments concerning this ratio measured only protons from  $\Lambda p \rightarrow n p$  process and  $\Gamma n$  was determined by the subtraction of all the other decay processes. Thus the obtained results must be much affected by small changes of the assumptions on final state interaction(FSI) effect and by the possible existence of the two-nucleon induced NMWD process,  $\Lambda NN \rightarrow NNN$ .

In order to measure this ratio unambiguously, we choose light s-shell  $\Lambda$  hypernuclei,  $^5_{\Lambda}\text{He}$ , so as to minimize the FSI effect in E462 experiment. In addition, we measured both of n+p- or n+n-pairs emitted from  $\Lambda +p \rightarrow n+p$  or  $\Lambda +n \rightarrow n+n$  NMWD process. When we select two-nucleon pairs which has back-to-back angular correlation, we can measure  $\Gamma n/\Gamma p$  ratio directly only from the ratio of n+p- to n+n-double coincidence pair numbers. The result of this measurement is free from the strength of FSI effect and also from the possible  $\Lambda NN \rightarrow NNN$  contribution. The experiment was extended to the heavier p-shell nuclear target,  $^{12}\text{C}$ , to study the mass-number dependence in E508.

Fig. 1 shows the setup of the decay coincidence system in E462/E508. Coincidence arms are sensitive to the all the particles emitted from the major decay modes of  $\Lambda$  hypernuclei. The top and bottom coincidence arms are placed to maximize the acceptance for the back-to-back n+n and n+p pairs.

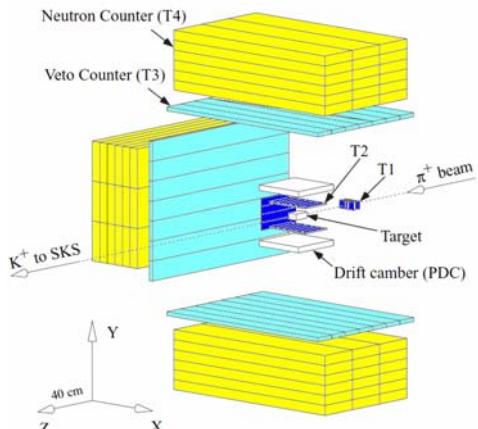


Fig.1 Experimental setup of KEK-PS E462/E508. Hypernuclei are produced by ( $\pi^+, K^+$ ) reaction on the target. Decay counter system is sensitive to all the decay particles from the major decay modes of  $\Lambda$  hypernuclei, neutron/  $\pi^0/\pi^\pm$ /protons. Also we can measure two decay particles in coincidence.

Fig. 2 shows the angular correlation of the n+p- and n+n-pairs from the NMWD of  ${}^5\Lambda$ He (double coincidence acceptance is corrected). In both of n+p/n+n, we observed clear back-to-back correlation. Also the sum energy distributions of two nucleons show peak at the Q-value of NMWD.

From the back-to-back coincidence ratio, we obtained the ratio of two-nucleon coincidence, Nnn/Nnp as

$$N_{nn}/N_{np} = 0.45 \pm 0.11 \pm 0.04 \quad (\text{for } {}^5\Lambda\text{He}) \quad [1]$$

with condition of  $\cos \theta_{NN} < -0.8$  and  $T_N > 30$  MeV.

The same ratio was successfully measured also in E508 experiment as

$$N_{nn}/N_{np} = 0.51 \pm 0.13 \pm 0.05 \quad (\text{for } {}^{12}\Lambda\text{C}) \quad [2].$$

Recent theoretical calculation considering heavier meson and/or direct quark exchange mechanism predicts the  $\Gamma n/\Gamma p$  ratio close to the measured ratio. Both of the dominance of  $\Lambda + p \rightarrow n + p$  decay process and the significant contribution of  $\Lambda + n \rightarrow n + n$  process are established for the first time.

In the recent detailed analysis for the  ${}^{12}\Lambda$ C decay, we found that not only the back-to-back two-nucleon emission yield but also the single nucleon spectra (both protons and neutrons) are quite successfully reproduced when we assume  $29 \pm 13\%$  contribution of  $\Lambda NN \rightarrow NNN$  decay in all the NMWD [3].

In the J-PARC experiment we are planning to carry out two  $\Lambda$  hypernuclear weak decay experiments with SKS at K1.8 beam line: (1) precise measurement of  $\Gamma n/\Gamma p$  ratio for  ${}^4\Lambda$ He (E22) and (2) high-statistics and low-proton detection threshold measurement of the NMWD of  ${}^{12}\Lambda$ C (E18) so as to study the detail of two-nucleon induced NMWD ( $\Lambda NN \rightarrow NNN$ ).

## References

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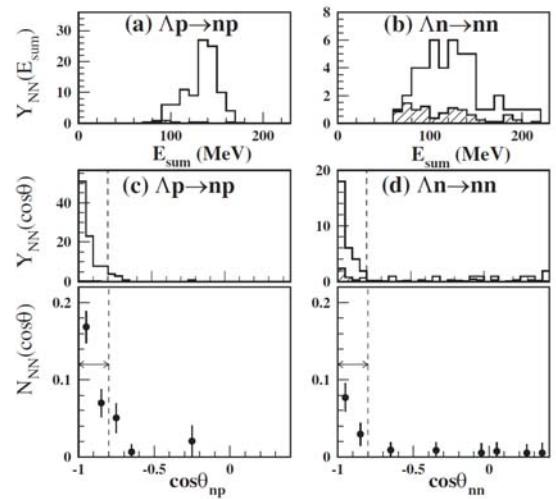


Fig.2 Two-nucleon angular correlation (top) and energy sum (bottom) from the NWMD of  ${}^5\Lambda$ He n+p (left) and n+n (right)

- (a)/(b): the sum energy of n+p and n+n, respectively
- (c)/(d): show the opening angle plots for n+p/n+n-pairs  
(top) bare number plot  
(bottom) acceptance corrected, normalized by NMWD

## Kaons in Nuclei

We have performed experimental exploration of theoretically predicted deeply bound kaonic nuclear states in  ${}^3\text{He}$  nucleus. Akaishi and Yamazaki first calculated large binding energy and narrow width for the total isospin  $T=0$  component ( $K^-ppn$ ) to be 108 MeV and 20 MeV (FWHM), respectively, reflecting rather strong attraction in the isospin  $I=0$  ( $K^-p$  and  $\bar{K}^0n$ ) channel [1]. The estimation was based on the information of kaonic hydrogen atom level shifts [2,3], low energy kaon scattering data, and assumption that  $\Lambda(1405)$  should be a  $\bar{K}N$  bound state.

One of the most interesting features of the kaonic nucleus is that the strong attraction of the kaon contract the surrounding nucleons, which implies extremely high density (several times larger than normal nuclear density) matter formation. Measurement of the kaon properties at such high energy density will provide precious information on the origin of hadron masses, the chiral symmetry breaking and its partial

restoration.

The experimental principle adopted uses stopped  $K^-$  on suprefluid helium target, and we focus on emitted neutron momentum measurement by the time-of-flight (TOF) method. The last orbit of kaonic  ${}^4\text{He}$  atom is 2p and the branching ratio from the last orbit to the nuclear kaon bound state accompanied with a neutron emission was estimated to be 1 % at minimum.

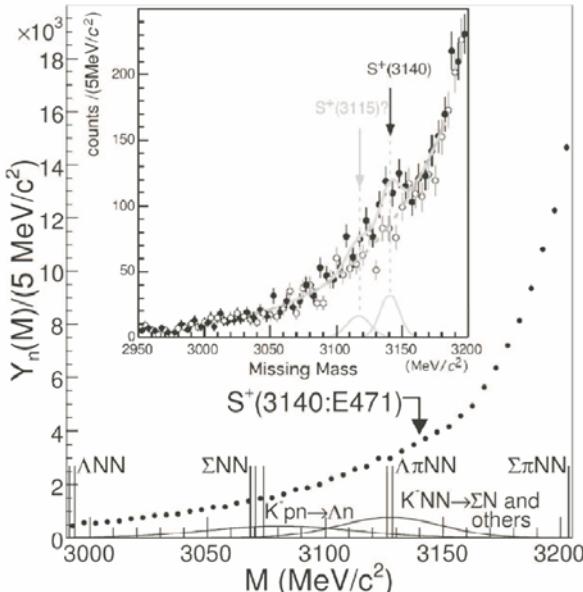
The exploration was performed from 2002/Sept. till 2005/Dec. as series of experiments at the KEK-PS (E471, E549, E570) with (almost) common experimental setup. An important modification to be noted is the installation of

Figure 1 : Mass spectrum of the reaction product in (stopped  $K^-, n$ ) reaction.

equipments dedicated to the emitted proton TOF and trajectory measurement for the detection of the kaonic nucleus with  $T=1$  state ( $K^-pnn$ ) in the latter stage (E549 and E570) of the series. Thus obtained experimental spectra are summarized in Figure 1 and Figure 2 for the neutron [4,5] and the proton [6,7], respectively.

Let us start with the neutron spectra. Figure 1 shows recently updated results of E549. The abscissa is the mass of the reaction product  $M$  shown with thresholds for possible decay channels. The statistical error bars are shown but are to small to be seen. The inset shows results in the previous experiment E471. Here, we do not discuss their consistency but we were able to conclude that the small structure seen in the inset, E471, near  $M = 3140 \text{ MeV}/c^2$  should be due to statistical fluctuations. The spectral shape is rather smooth and elaborate analysis shows upper limit of the kaonic nucleus formation.

Let us set our sight to the proton spectra. Figure 2 shows results of E549 and E471. We recognize their



large inconsistency. Even if the experimental conditions adopted are not equal, the above difference is now considered to be a consequence of erroneous time-walk correction (time vs. pulse height correlation) included in the early experiment E471, and the peak structure seen near  $M \sim 3115$  MeV/c $^2$  is an artifact. Again, upper limit is given for the kaonic nucleus ( $K^- pnn$ ) formation.

After the completion of above series of experiments, the KEK-PS was shut down to switch to a new facility J-PARC. Presently, we are preparing for an experiment to search for  $K^- pp$  nuclear system at the K1.8BR beamline of the J-PARC [8]. With the same setup, we are as well planning to perform dedicated studies of the  $\Lambda(1405)$  formation and decay processes, since this state became the cornerstone of the prediction of the possible existence of the deeply bound kaonic nuclei.

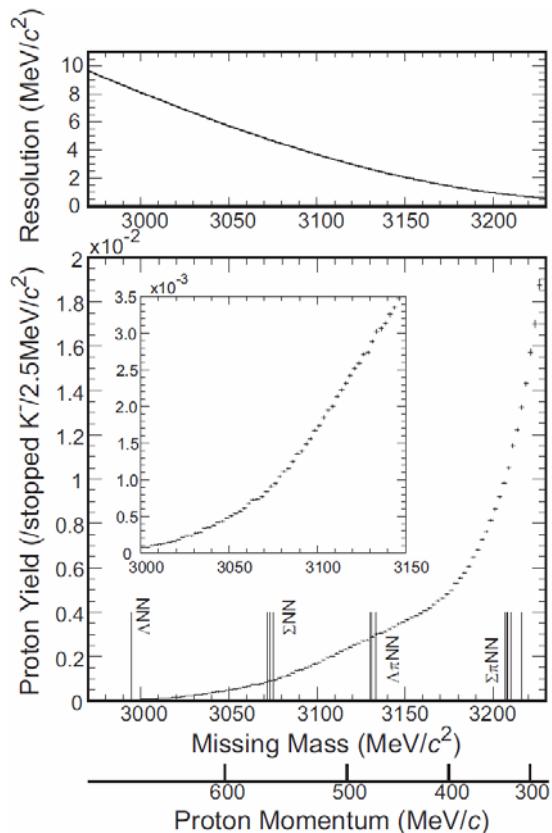
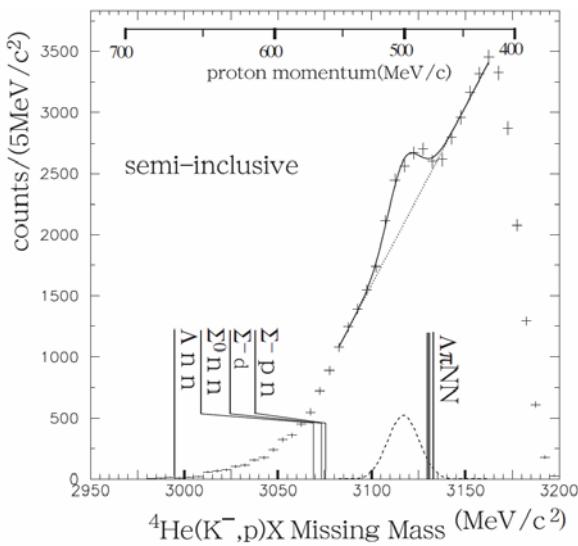


Figure 2 : Mass spectra of the reaction product in (stopped  $K^- ,n$ ) reaction for E471 (left) and E549 (right).

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## Double-Kaons in Nuclei

Possible existence of anti-kaonic nuclear clusters has been investigated extensively both with theoretical and experimental approaches recently. In view of the strongly attractive  $\bar{K}N$  interaction below threshold, the existence of nuclear clusters with more than one  $K^-$  is predicted also, such as  $\bar{K}\bar{K}NN$  double anti-kaonic nuclear systems[1]. Double anti-kaonic nuclear clusters are predicted to have binding energies up to 300 MeV and nuclear densities exceeding  $\sim 10$  times that of the average one  $\rho(0) = 0.17 \text{ fm}^3$ , thus producing conditions in the phase diagram of hadronic matter for which phase transitions to Kaon-condensation / color-superconductivity or pre-cursor effects for these may be reached at low temperature.

We propose [2] to produce double  $K^-$  simultaneously at close distance in a nuclear target and explore the expected ``strong attraction'' mediated by double anti-kaons in the nuclear environment, leading to cold and dense Fermion matter.

The elementary anti-proton annihilation reaction, which produces the two pairs of ( $K^+ K^-$ ), considered is

$$\bar{p} + p \rightarrow K^+ + K^+ + K^- + K^- - 98 \text{ MeV}, \quad (1)$$

with a negative  $Q$ -value of 98 MeV, so forbidden for stopped antiprotons. However, if multi kaonic nuclear cluster exists with deep bound energy, as suggested by Ref.[1], the following  $\bar{p}p$  annihilation reactions will be possible on He targets[2]:

$$\bar{p} + {}^3\text{He} \rightarrow K^+ + K^0 + ppK^-K^- + B_{KK}^{pp} - 109 \text{ MeV}. \quad (2)$$

This double kaonic nuclear cluster process occurs if the binding energy of the two  $K^-$  in a  $ppK^-K^-$  cluster  $B_{KK}^{pp}$  exceeds 109 MeV.

For reactions in the final state (2), we have no exclusive signal for the production of a  $S = -2$  system, thus we can resort to the detection of the decay of the cluster  $X$  into two  $S = -1$  hyperons, such as two  $\Lambda$  particles. So we can measure the missing mass from the  $K^+K^0$  energies and also the decay products in addition to reconstruct the invariant mass of the double kaonic nuclear cluster. Moreover, all particles in the final state can be charged, if the  $ppK^-K^-$  intermediate state will decay to  $\Lambda\Lambda$ , so we can detect all the particles rather easily. For the decay branching-ratio to the  $\Lambda\Lambda$  final state, we needs detailed theoretical evaluation, although this coherent kaon absorption strength would not be small because of the favored isospin-zero channels.

We submitted the LOI to J-PARC to search for such double anti-kaonic nuclear systems at the existing K1.8BR beamline with the E15 spectrometer, or the newly K1.1 beamline. Now, we are preparing the proposal to J-PARC with more detailed studies, and this will be submitted in the next summer.

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## ϕ-mesons in Nuclei

A proton is believed to be composed of two 'Up' quarks and one 'Down' quark, and having mass of 938 MeV/c<sup>2</sup>. On the other hand, the mass of bare quark, *i.e.*'Up' and 'Down' quark are known to be a few MeV/c<sup>2</sup>, *i.e.* contribution of the constituent quarks to the mass of the proton is only a few % at most. The question needed to be answered is the mechanism which generates more than 90% of the proton mass from vacuum. This mechanism is now known as spontaneous breaking of the chiral symmetry, which creates non zero  $\langle\bar{q}q\rangle$  expectation value in vacuum. This  $\langle\bar{q}q\rangle$ -condensation is the major source of the low lying hadron masses such as protons, neutrons, pions, *etc.* In the theoretical framework, the  $\langle\bar{q}q\rangle$  expectation value (chiral order parameter in other word) is a function of temperature and chemical potential (density).

Various experiments have been performed to detect the restoration of the chiral symmetry. One of the approach is the formation of meson nucleus bound state and measurement of the energy levels of the state which will be a direct connection to the  $\langle\bar{q}q\rangle$  expectation value in nucleus.

Here we are focusing on the  $\phi$  meson in nucleus. Experimental study of  $\phi$  meson invariant mass spectra via di-electron measurement in the  $pA$  reaction reported about 3 % mass reduction of the  $\phi$  meson in medium-heavy nuclei (Cu), but no mass shift in light nuclei (C) [1]. On the other hand, in this measurement, the natural width broadening of  $\phi$  meson is only found to be  $\Gamma_\phi^{in\ media}/\Gamma_\phi^{free} \sim 3.4$  [1]. This result indicates that the most of the produced  $\phi$  meson in nuclear matter kept its property as a  $\phi$  meson except to its mass. Here we are considering the meaning of 3 % (= 30 MeV/c<sup>2</sup>) mass reduction of  $\phi$  meson in nucleus. Hint comes from the situation of kaons in nuclei. The reference [2] pointed out that mass of the K<sup>-</sup> will be reduced in nuclear matter due to strong attractive potential exist between K<sup>-</sup> and nucleon. This theoretical prediction indicates that the "mass reduction of  $\phi$  meson in nucleus" will be directly connected with the possible existence of attractive potential between  $\phi$  meson and nucleus. The depth of the potential is expected to be at the same order as the one giving the mass reduction which has been measured. Therefore,  $\phi$  meson bound in nuclear state, if it exists, will be a unique tool to investigate properties of  $\phi$  meson in nuclear media.

Here we are focusing on the reaction  $\bar{p} + p \rightarrow \phi + \phi$  channel as a elementary process to produce  $\phi$  meson bound state in nucleus. Using this reaction, we will be able to perform missing mass spectroscopy via the  $X_A^Z(\bar{p}, \phi)\phi \cdot X_{A-1}^{Z-1}$  process. This elementary process has two interesting feature. One is the rather large  $\phi\phi$  production cross section near the production threshold (0.9-1.3 GeV/c), which is about 4  $\mu b$ . The other is the yield of the kaon-associated  $\phi$  production channels,  $\phi K^+K^-$  and  $K^+K^-K^+K^-$ , which are much smaller than that of the double  $\phi$  production channel for the incident  $\bar{p}$  momentum below 1.4 GeV/c[3], which is only less than 1/10 of the  $\phi\phi$  production cross section. Those experimental fact indicates that if we find the  $\phi$  meson as a final state particle, one more  $\phi$  meson will be produced with high probability. Moreover, once we select only the  $\phi$  meson emitted 0° respect to the beam direction in laboratory system, the momentum of the other produced  $\phi$  meson is only about 200 MeV/c, which is almost of the same order of magnitude with the Fermi momentum of nucleon in nucleus. Therefore we

expect marginal  $\phi$  meson sticking probability in nucleus. The experimental method to search for  $\phi$  meson bound state is as follows. First of all  $\bar{p}$  beam with momentum of 1.0 GeV/c will be shooting on to the target nucleus. Then forward going  $\phi$  meson via its decay particles,  $K^+K^-$  will be detected for missing mass spectroscopy. Once the  $\phi$  meson captured by the nucleus, the final state of the  $\phi$ -meson bound state is expected to contain  $K^+\Lambda$ , owing to the final state interaction between captured  $\phi$  meson and nucleon. Therefore, for this measurement, we will required final state  $K^+\Lambda$  together with  $K^+K^-$  pair from forward going  $\phi$  meson decay. To detect forward going Kaons, together with the  $K^+\Lambda$  in the final state, huge reduction of background on missing-mass spectrum efficiently in  $(\bar{p}, \phi)$  spectroscopy will be realized. From this measurement, one can independently deduce the mass shift information. A systematic study over several nuclear targets will yield a unique, definitive and precise determination of the in-medium mass modification of the vector meson  $\phi(\bar{s}s)$ . In spite of the lower cross section of  $p(\bar{p}, \phi)\phi$ , we can expect an excellent ground-state formation event rate of 240 per month using the  $\bar{p}$  beam of  $2 \times 10^6$  per spill on a carbon target. The conceptual design for the spectrometer is shown in Figure 1. The experimental proposal gathering an international team of proponents, has been submitted to J-PARC Program Advisory Committee and detail design and R&D work for the detector is under the way.

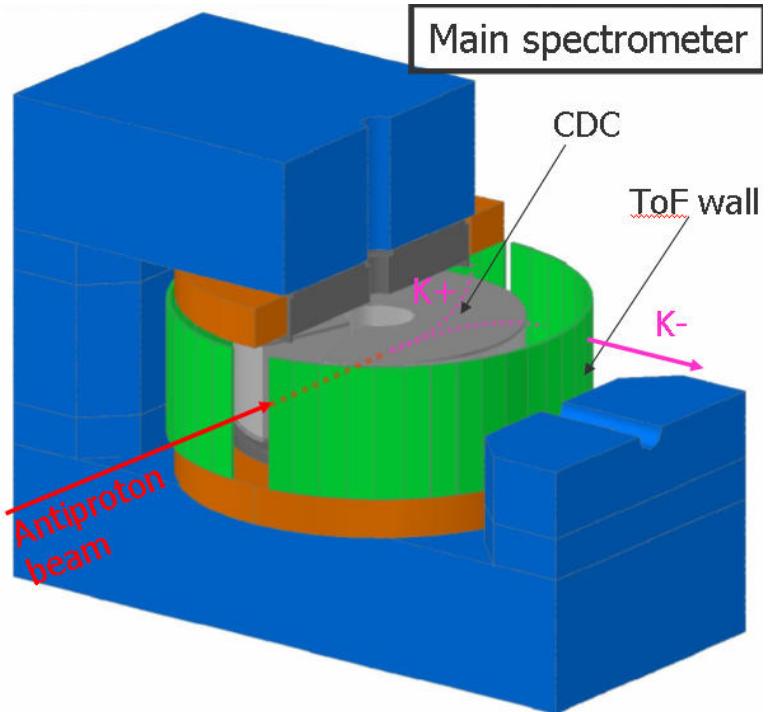


Figure 1 : Conceptual design for the Spectrometer

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### 3.2.1.2. R&D for Hadron Physics

#### Cylindrical Detector System (CDS)

Although there are several experimental reports of search for deeply-bound kaonic nuclear states [1,2,3,4,5,6], the situation is still controversial because there is no conclusive evidence for the observation of such bound states. Actually, in both experimental and theoretical sectors, the obtained binding energy and decay width of the  $K^-pp$  state vary from a few 10 MeV to around 100 MeV.

In order to clarify this controversial issue, the J-PARC E15 experiment was proposed to search for the simplest kaonic nuclear bound states, namely  $K^-pp$ , via the in-flight  $^3He$  ( $K^-,n$ ) reaction using 1.0 GeV/c  $K^-$  beam[7]. This experiment has the advantage that the exclusive measurement can be performed by a simultaneous measurement of missing mass spectrum using the primary neutron and invariant mass spectroscopy via the expected decay  $K^-pp \rightarrow \Lambda p \rightarrow p\pi^-\bar{p}$ .

The E15 spectrometer consists of four parts, namely Beam-line spectrometer, Cylindrical Detector System (CDS) with liquid  $^3He$  target system, Beam Sweeping Magnet and Neutron TOF wall. The decay particles from the expected decay  $K^-pp \rightarrow \Lambda p \rightarrow p\pi^-\bar{p}$  are detected by the CDS, and the Neutron TOF wall detects forward neutrons whose flight length are about 15 m. Incident kaons which pass through the target are bent by the Beam Sweeping Magnet which is placed just after the CDS. The expected spectrometer performance for the  $K^-pp$  measurement is  $9.2 \text{ MeV}/c^2 (\sigma)$  for the missing-mass resolution via neutron detection, and  $16 \text{ MeV}/c^2$  for the invariant-mass resolution via the  $K^-pp$  decay, where we assume the  $K^-pp$  binding energy to be  $100 \text{ MeV}/c^2$ . In the following sentences, a preparation status of the CDS is described.

In the CDS, all detectors are configured cylindrically. Trajectory of particles is reconstructed with the Cylindrical Drift Chamber (CDC) which operates in a magnetic field of 0.5 T provided by the solenoid magnet. The Cylindrical Detector Hodoscope (CDH) which surrounds the CDC is used for the trigger and the particle identification counter. The solid angle of the CDS from the center is about 7.4 sr.

The CDC consists of two aluminum end-plates of 20 mm thickness, a CFRP tube with 1 mm thick as a inner wall and six aluminum blocks which are placed outside the tracking volume. The CDC uses gold-plated tungsten of  $30\mu\text{m}\phi$  for the sense wires and gold-plated aluminum of  $100\mu\text{m}\phi$  for the field and guard wires. These wires are supported by feedthroughs with a bush fixing the wire position. The total length along beam axis of tracking volume is about 840 mm. The argon(50%)-ethane(50%) mixed gas is used at 1 atm. The CDC has 15 layers of hexagonal cells with typical drift length of 9 mm, which are grouped into 7 super layers (A1,U1,V1,A2,U2,V2,A3).

The information of the longitudinal position is obtained by 8 stereo layers with tilt angle of typically 3.5 degree. The number of readout channels is 1816 and the total number of wires in CDC is 8064.

The CDH consists of 36 segments, and each segment is individually mounted on the inner wall of the solenoid magnet. Hamamatsu type R7761 fine-mesh Photo Multipliers (PMT) with 1.5 inch diameter. The measured time resolution of the CDH with cosmic rays is typically 71 psec, achieving the design goal.

The performance test of CDC was done using cosmic ray. Figure 1 shows a typical residual distribution obtained from the test. Applying a gaussian fit to this spectra, we obtained  $\sigma = 206 \mu\text{m}$ . In order to estimate the intrinsic resolution of the CDC, we performed a simple simulation, where straight tracks are generated and make a hit with a given intrinsic resolution and then the data are analyzed using the same routine as that used in the cosmic-ray test. Figure 2 shows the results of the simulation, together with the result of the cosmic-ray test. By comparing the results, the intrinsic resolution for a straight track is determined to be  $200 \mu\text{m}$ .

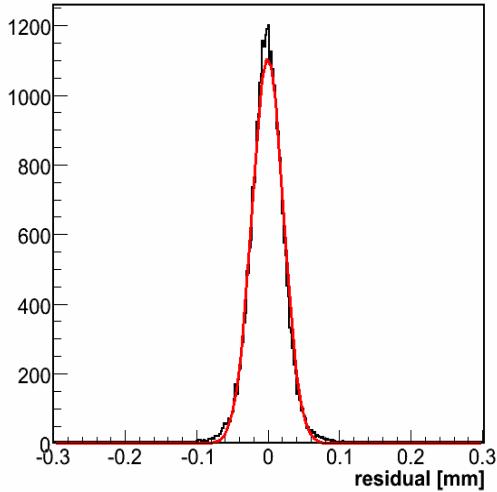


Figure 1 : A typical residual distribution obtained from the cosmic-ray test.

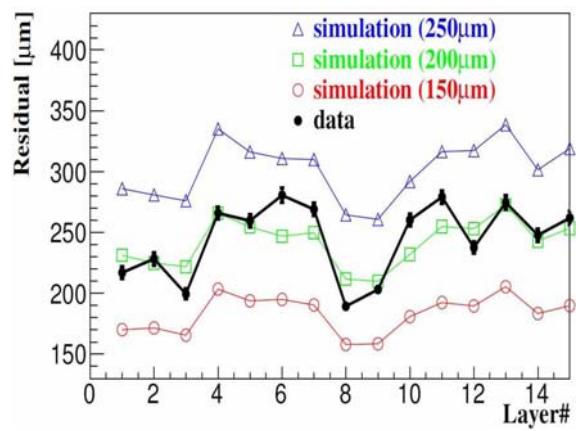


Figure 2 : Layer dependence of CDC resolution. The black line shows the measured resolution from the cosmic-ray test. The red, green and blue lines correspond to the results of simulations, assuming an intrinsic resolution of 150, 200 and 250 mm, respectively.

In order to estimate the performance for CDS, we performed a Monte Carlo simulation. Basic parameters used for the MC simulation are (1) the intrinsic resolution of CDC  $\sigma = 200 \mu\text{m}$ , which is determined from the cosmic-ray test, and the field strength inside the magnet 0.5 T. To evaluate the invariant mass resolution for  $\Lambda$  and  $K^-pp$  states,  $K^-pp$  are generated in the following way.

- The binding energy and decay width of  $K^-pp$  are set to 100 MeV and 0, respectively.
- $K^-pp$  is generated uniformly in the center-of-mass frame.
- $K^-pp$  decays into  $\Lambda$  and proton in an isotropic way.

Figure 3 and Figure 4 shows reconstructed mass spectrum of (a)  $\Lambda$  and (b)  $K^-pp$ . The expected invariant mass resolution is  $2.2 \text{ MeV}/c^2$  for  $\Lambda$  and  $15 \text{ MeV}/c^2$  for  $K^-pp$ .

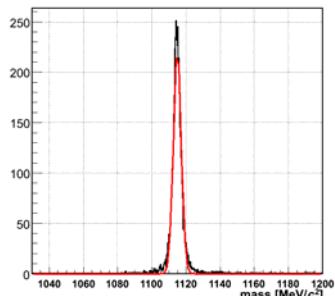


Figure 3 :  $\Lambda$  invariant mass

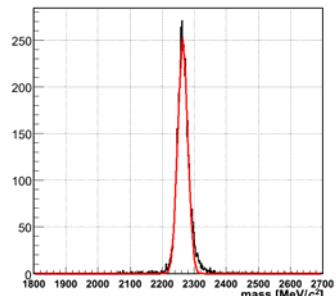


Figure 4 :  $K^-pp$  invariant mass

Finally, The CDC and CDH were already successfully installed in the solenoid magnet as shown in Figure 5. The excitation of the solenoid and the performance study of CDC and CDH in the magnetic field will be performed in March to April 2010.

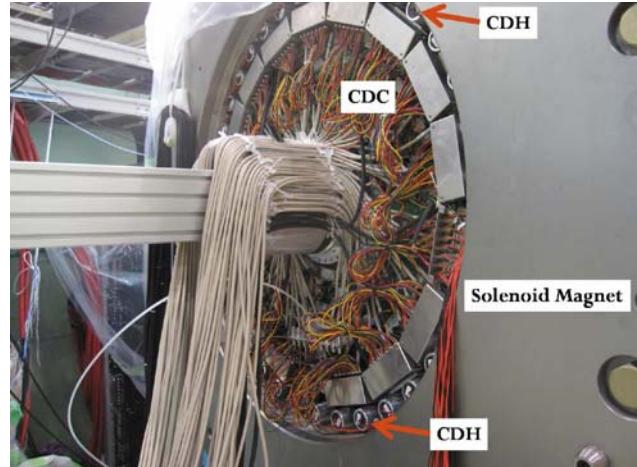


Figure 5 : A picture of CDS. CDC and CDS are installed in the solenoid magnet

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``A search for deeply-bound kaonic nuclear states by in-flight  ${}^3\text{He}(\text{K}^-, \text{n})$  reaction", (2006)

## Development of a TGEM-TPC

The experiment J-PARC E15 searches for the simplest kaonic nuclear bound state,  $\text{K}^-\text{pp}$ , by in-flight  ${}^3\text{He}(\text{K}^-, \text{n})$  reaction. The key point of such a measurement is to identify precisely  $\Lambda$  and  $\Sigma$  decays with the secondary vertex reconstruction, because expected decay modes of  $\text{K}^-\text{pp}$  are  $\text{p}\Lambda/\text{p}\Sigma^0$  and  $\text{p}\pi\Sigma$ . To realize these measurements we are developing a Thick Gas Electron Multiplier (TGEM) Time Projection Chamber (TPC) as an inner tracker for the E15 upgrade. The readout of a TPC with the TGEM has many advantages: robustness and cost-effectively fabrication of TGEM, unnecessary of the support frame, good single-point accuracy and multi-track resolution in projection, substantial reduction of ion feedback, and so on. The requirements for the detector are spatial resolutions of within 1mm in z-direction, and low material budget in the detector acceptance.

The TPC has a cylindrical design with an inner diameter of 170 mm and an outer diameter of 280 mm, filled with P10 gas at atmospheric pressure. The drift length is 30~cm with the field-cages of a double-sided flexible Printed Circuit Board (PCB) with staggered strip electrodes. The schematic view and photograph of the TPC are shown in Figure 1. We use a double-TGEM structure for amplification, and signals are read out with 4~mm long and 20~mm wide pads printed on a standard PCB. The TGEM is

made from double-clad FR4 plate, using standard PCB techniques, with mechanically drilled holes and chemically etched rims. We use 400  $\mu\text{m}$  FR4 plates with holes of  $\phi \sim 0.3$  mm and rims of  $\phi \sim 0.36$  mm, typically.

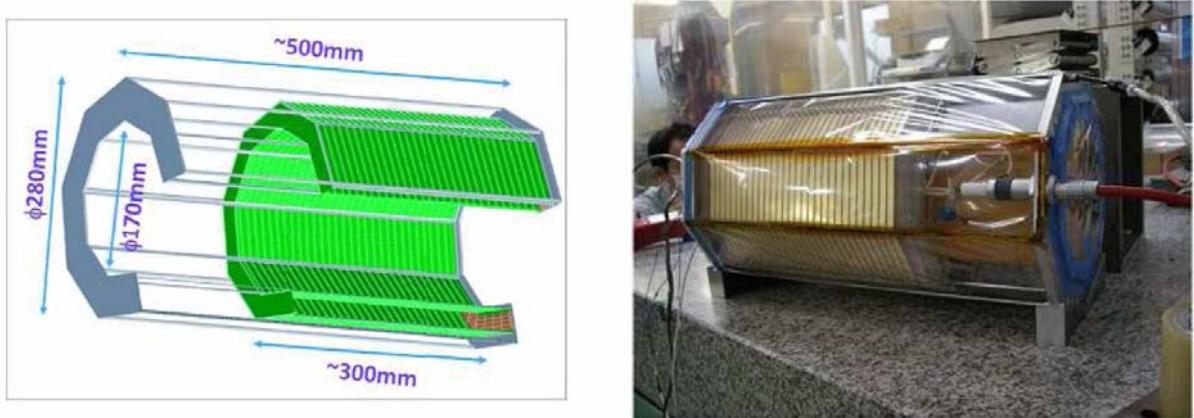


Figure 1 : A schematic view and photograph of the TPC.

To study the performances of the TGEM, we use a prototype TGEM with an active area of  $10 \times 10 \text{ cm}^2$  produced by REPIC Corp., Japan. The requirements for the TPC-TGEMs are effective gain of approximately  $10^4$  and stability of the gain and energy resolution. We tried many types of TGEM, i.e. 200/400  $\mu\text{m}$  thickness,  $\phi$  300/500  $\mu\text{m}$  hole size, 0/30/50/100  $\mu\text{m}$  rim size, and so on. All types of TGEM with thickness of 400  $\mu\text{m}$  achieve the maximal effective gain of above  $10^4$ , but TGEMs with larger rims require higher voltage. And TGEMs with no-rims and small-rims (30  $\mu\text{m}$ ) work rather stably, however, TGEMs with large-rims (100 & 50  $\mu\text{m}$ ) are in-stable. It is highly possible that this instability is caused by charge-up of insulator not metalized, but that's not understood well so far. Farther studies for basic TGEM behavior and performance are in progress now.

To avoid the effects of rims, we are developing a new resistive-electrode TGEM (400  $\mu\text{m}$  thickness) which has electrodes coated with graphite paint and drilled holes of  $\phi \sim 0.3$  mm without rims, which has an advantage of being fully spark-protected. In addition, we are developing a new hybrid-TGEM, i.e. TGEM with carbon and copper electrodes on each side without rims, recently. These new types of TGEM would solve the difficulties of conventional TGEM.

## Development of Cryogenic Targets

A super-fluid liquid  ${}^4\text{He}$  target was developed to experimental search for deeply bound kaonic nuclear states by the  ${}^4\text{He}(K^-_{\text{stopped}}, N)$  reaction at the K5 beamline of the KEK 12-GeV proton synchrotron (E471, E549)[1]. One of the most critical background sources is the negative pion absorption,  $\pi^- NN \rightarrow NN$ . A most effective way to suppress this background is to reduce the amount of material around the liquid  ${}^4\text{He}$  target. It was accomplished by two approaches. One was to utilize the properties of super-fluid helium of extremely high thermal conductivity and low vapor pressure. The other was to develop a thin-walled

CFRP vacuum chamber and a PET-based target cell. The wall thickness of the CFRP vacuum chamber was only 0.9 mm with the inner diameter of 310 mm. The dimensions of the target cell were 200 mm inner diameter and 150 mm length with the volume of 4.7 l (2.175 g/cm<sup>2</sup> in the beam direction at <sup>4</sup>He temperature below 2 K), and was equipped 75  $\mu$ m thick beam windows made of a PET film. During the experimental period, liquid <sup>4</sup>He was stably-held at 1.3 K in the target cell. The performance of cryostat was the heat load to the cell and reservoir of 0.16 W and a liquid-helium consumption rate of 45 l/day.

A liquid <sup>3</sup>He target system has been developed for successive experiment of the deeply bound kaonic nuclei by in flight <sup>3</sup>He(K<sup>-</sup>, n) at J-PARC (E15). Figure 1 shows the design of the liquid <sup>3</sup>He target which is based on the techniques developed for the super-fluid <sup>4</sup>He target. The target was designed by L-form to place a target cell in the center of the cylindrical detector system. Liquid <sup>4</sup>He is supplied to the <sup>4</sup>He-separator (~4 K) inside the target system from a 1000 liter dewar placed outside by a transfer tube. From the separator, only liquid <sup>4</sup>He drop into the <sup>4</sup>He-evaporator which is placed under the separator. The inside of evaporator is decompressed down to ~2 mbar to keep a temperature of 1.3 K. The <sup>3</sup>He gas is cooled down and liquefied in the 3He-heat exchanger which is placed under the <sup>4</sup>He-evaporator. The target cell almost 1 m away from cryostat is filled with liquid <sup>3</sup>He liquefied in the heat exchanger. The cooling tests were performed at K5 experimental area in KEK-PS north counter hall. The setup of cooling system was the almost same as production run setting for E15, and 250 liters <sup>3</sup>He gas was used as the target material on the test. As a result, the target cell was cooled down and kept at 1.3 K. Since the pressure in the evaporator and cell is also measured at the same time, we were able to study the correlation of the pressure and temperature. The result is shown in Figure 2, where black and white dots are corresponding to the cell and to the evaporator. Although there are small shift due to the error of the measuring device, both measurements are in good

J-PARC E15 LHe-3 Target

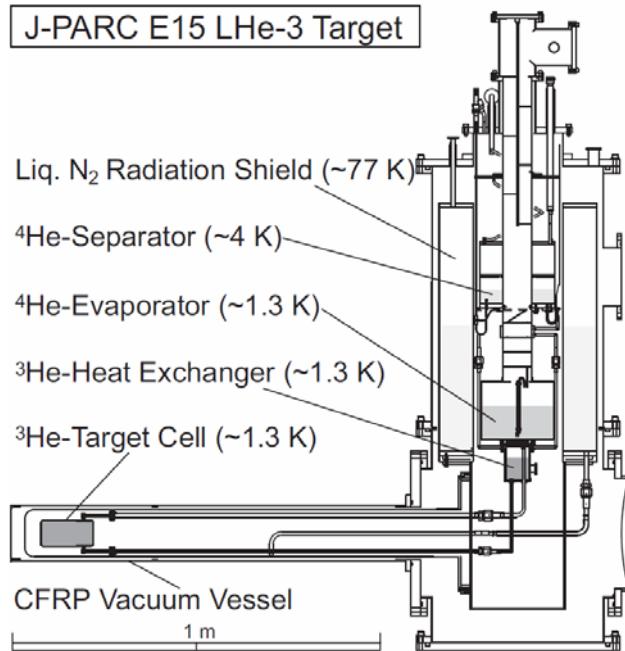
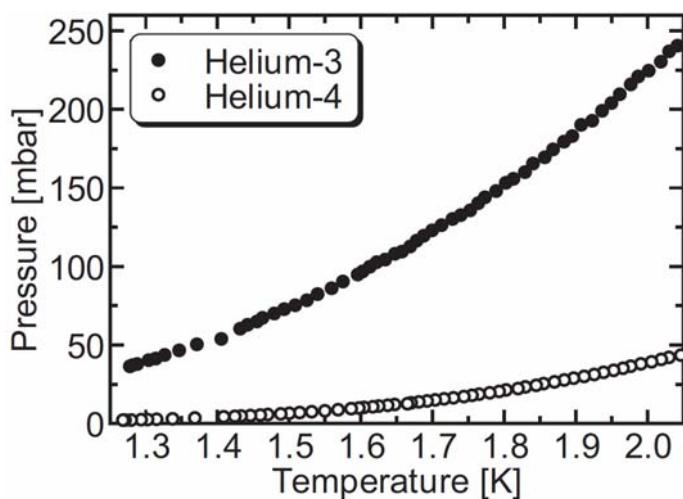
Figure 1: A schematic side view of the E15 liquid <sup>3</sup>He target.

Figure 2 : The measurements of the vapor pressure curve by the correlation of the pressure and temperature in the cell and evaporator

agreement with vapor pressure curve. It is clear that liquid  $^3\text{He}$  exists in the target cell. In addition, the total of the heat load to the cell, heat exchanger and evaporator is 0.19 W. High-performance cooling system for E15 liquid  $^3\text{He}$  target was completed.

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### 3.2.1.3. Theoretical Activities

#### Electron population during the cascade of kaonic nitrogen atoms and the charged kaon mass

The Particle Data Group assigned  $493.677 \pm 0.013$  MeV to the charged kaon mass as a world average. However, there exists a serious disagreement between the most two recent mass measurements using kaonic atom x-rays, in which the deduced masses differ about 60 keV although their individual uncertainties are about 6-7 keV. In order to settle this discrepancy, new precise charged kaon mass measurement using kaonic nitrogen atom x-rays in a gaseous target is planned at the DAΦNE.

In order to achieve the required accuracy, the electron screening effect on kaonic x-ray energy needs to be estimated correctly, which is determined by the balance between Auger electron emission and electron refilling during the atomic cascade process. Thus, the cascade calculation of kaonic nitrogen atoms involving electron refilling process is performed and the electron population during the cascade of kaonic nitrogen atom is theoretically investigated to be useful for the future kaon mass measurement at DAΦNE. It is found that the one 1s-electron remains with the probability of 4% at the moment of the  $6 \rightarrow 5$  kaonic x-ray emission in the gaseous target at a density  $\rho = 3.4\rho_{\text{NTP}}$ , which corresponds to the experimental condition at DAΦNE.

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#### Calculation of $^3\text{He}(\text{in-flight K}, \text{n})$ spectrum for deeply-bound K- pp state

In the study of KN interaction in nuclei, it is important to verify the presence of the deeply-bound kaonic nuclei. A three-body KNN bound state with a  $[\text{K}^\otimes \{\text{NN}\}_{l=1}]_{l=1/2}$ ,  $J^\pi = 0^-$  configuration, which is

called "K<sup>-</sup>pp" symbolically, is suggested to be the lightest and the most fundamental K<sup>-</sup> nucleus. Many theoretical works for K<sup>-</sup>pp system support the existence of K<sup>-</sup>pp bound state, but the predicted binding energies and widths are not converged because of an ambiguity of KN interaction, together with a different procedure for a three-body calculation involving decay processes. Moreover, several experimental observations of K<sup>-</sup>pp state have been reported, but they disagree with each other. Therefore, a new experimental search for the deeply-bound K<sup>-</sup>pp state using <sup>3</sup>He(in-flight K<sup>-</sup>, n) reaction have been proposed at J-PARC (J-PARC E15 experiment).

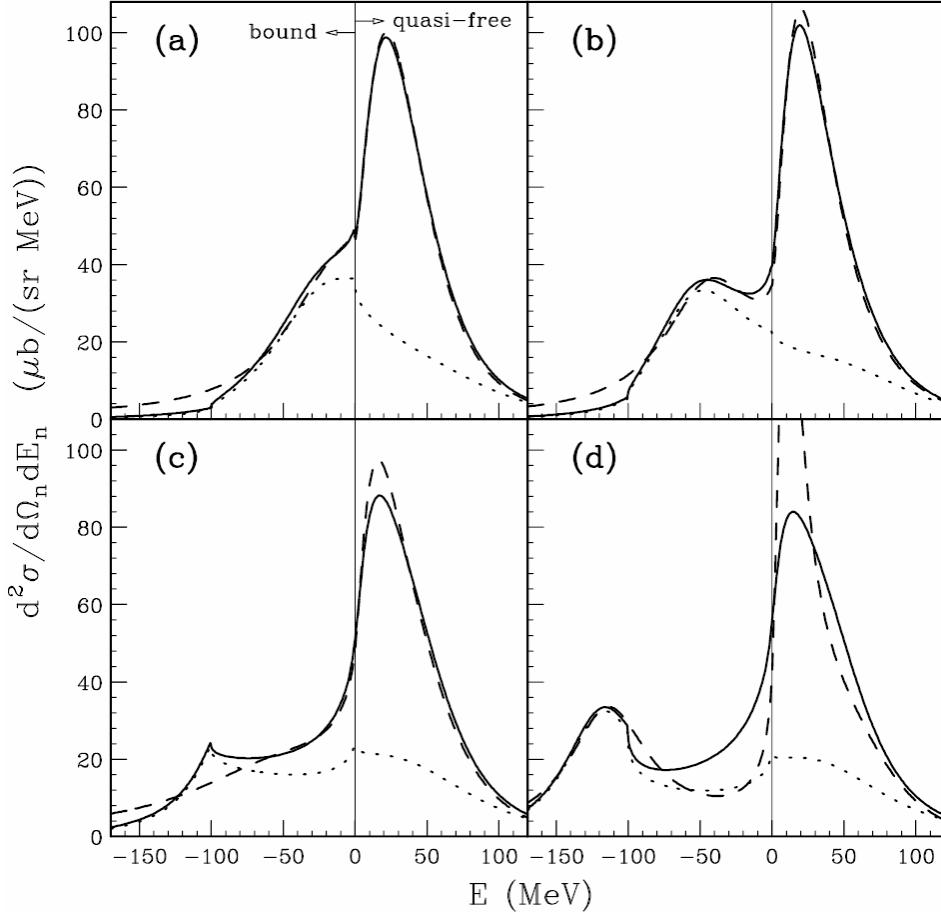


Figure 1 : Calculated inclusive spectra of the <sup>3</sup>He(in-flight K<sup>-</sup>, n) reaction at  $p_{K^-} = 1.0 \text{ GeV}/c$  and  $q_{\text{lab}} = 0^\circ$  as a function of the energy E of the K<sup>-</sup>pp system measured from K<sup>-</sup>+p+p threshold for several phenomenological K<sup>-</sup>"pp" optical potentials. The solid and dashed curves denote the inclusive spectra with the energy-dependent  $U^{\text{opt}}(E)$  and energy-independent  $U^{\text{opt}}_0$  potentials, respectively. The dotted curve denotes the  $L=0$  component in the inclusive spectrum for  $U^{\text{opt}}(E)$ . The vertical line at E = 0 MeV indicates the K<sup>-</sup> + p + p threshold, and the left- and right-hand sides of this line are the K<sup>-</sup> bound and quasi-free scattering regions, respectively.

The formation of a deeply-bound K<sup>-</sup>pp state by the <sup>3</sup>He(in-flight K<sup>-</sup>, n) reaction is theoretically investigated in a distorted-wave impulse approximation using the Green's function method. The expected inclusive and semi-exclusive spectra at  $p_{K^-} = 1.0 \text{ GeV}/c$  and  $\theta_{\text{lab}} = 0^\circ$  are calculated for the forthcoming J-PARC E15 experiment. We demonstrate these spectra with several phenomenological K<sup>-</sup>"pp" optical potentials  $U^{\text{opt}}(E)$  which have an energy-dependent imaginary part multiplied by a phase space suppression factor, fitting to recent theoretical predictions or experimental candidates of the K<sup>-</sup>pp bound state. The results show that a cusp-like peak at the  $\pi\Sigma N$  threshold is an unique signal for the K<sup>-</sup>pp bound state in the

spectrum including the  $[K^- pp] \rightarrow Y + N$  decay process from the two-nucleon  $K^-$  absorption, as well as a distinct peak of the  $K^- pp$  bound state. The shape of the spectrum is explained by a trajectory of a moving pole of the  $K^- pp$  bound state in the complex energy plane. The importance of the  $[K^- pp] \rightarrow Y + N$  spectrum is emphasized in order to extract clear evidence of the  $K^- pp$  bound state.

## References

- [1] T. Koike and T. Harada, Phys. Rev. C80(2009)055208
- [2] T. Koike and T. Harada, Nucl. Phys. A804(2008)231
- [3] T. Koike and T. Harada, Phys. Lett. B652(2007)262

## p $\Xi^0$ force studied with lattice QCD

Study of the baryon-baryon interaction is an important subjects in the nuclear physics. The present hyperon-nucleon (YN) and hyperon-hyperon (YY) interactions have large uncertainties despite these interactions playing important roles in high density nuclear systems such as interior of neutron stars. For example, no reliable phase shift analysis of the  $\Lambda N$  scattering has as yet been performed experimentaly so that different phaseshifts have been predicted by different theoretical works. In principle, the interaction should be understood in terms of the dynamics of quarks and gluons, namely quantum chromodynamics (QCD).

We study the p $\Xi^0$  force by using quenched lattice QCD. The Bethe-Salpeter amplitude is calculated for the lowest scattering state of the system so as to obtain the p $\Xi^0$  potential. The numerical calculation is performed with  $\beta = 5.7$ , the lattice spacing of  $a=0.1416(9)$  fm, on the  $32^3 \times 32$  lattice. The spatial lattice volume is  $(4.5 \text{ fm})^3$ . Two kinds of ud quark mass are used, corresponding to  $m_\pi \approx 0.37 \text{ GeV}$  and  $0.51 \text{ GeV}$ . Tables 1-2 compare the hadron masses calculated from the lattice QCD with the experimental values. Figure 1 shows the effective central potential, obtained from the wave function at the time slice  $t-t_0=6$  with the hopping parameter  $\kappa_{ud}=0.1678$ , corresponding to  $m_\pi \approx 0.37 \text{ GeV}$ . The scattering length is obtained from Lüscher's formula[1]. As is seen in Figure 2, the p $\Xi^0$  interaction is both attractive at  ${}^1S_0$  and  ${}^3S_1$  channels, and the interaction in the  ${}^3S_1$  is more attractive than in the  ${}^1S_0$ . These attractive forces become stronger as the u,d quark mass decreases.

| $\kappa_{ud}$ | $N_{\text{conf}}$ | $m_\pi$  | $m_p$  | $m_K$    | $m_{K^*}$ |
|---------------|-------------------|----------|--------|----------|-----------|
| 0.1678        | 1283              | 368(1)   | 813(4) | 554(5)   | 884(2)    |
| 0.1665        | 1000              | 511.2(6) | 861(2) | 605.3(5) | 904(2)    |
|               | Exp.              | 135      | 770    | 494      | 892       |

Table 1: Meson masses in the unit of MeV. The numbers in parenthesis show the errorbar in the last digit.

| $\kappa_{ud}$ | $N_{\text{conf}}$ | $m_N$   | $m_\Lambda$ | $m_\Sigma$ | $m_\Xi$ |
|---------------|-------------------|---------|-------------|------------|---------|
| 0.1678        | 1283              | 1167(7) | 1266(6)     | 1315(6)    | 1383(6) |
| 0.1665        | 1000              | 1300(4) | 1354(4)     | 1357(4)    | 1419(4) |
|               | Exp.              | 940     | 1116        | 1190       | 1320    |

Table 2: Baryon masses in the unit of MeV. The numbers in parenthesis show the errorbar in the last digit.

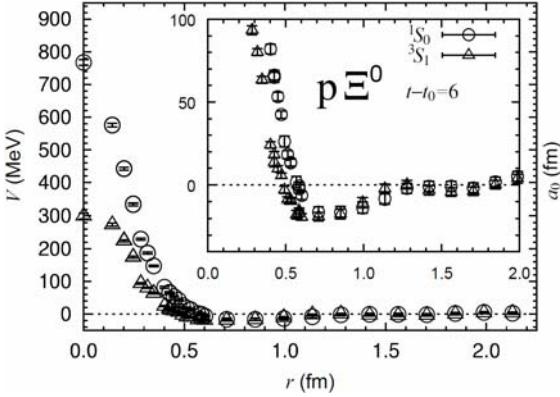


Figure 1: The effective central potential for  $p\Xi^0$ , in the  ${}^1S_0$  (circle) and  ${}^3S_1$  (triangle), obtained from the wave function at time slice  $t-t_0=6$ . The hopping parameter  $\kappa_{ud}=0.1678$  is used for the u,d quark. The inset shows its enlargement.

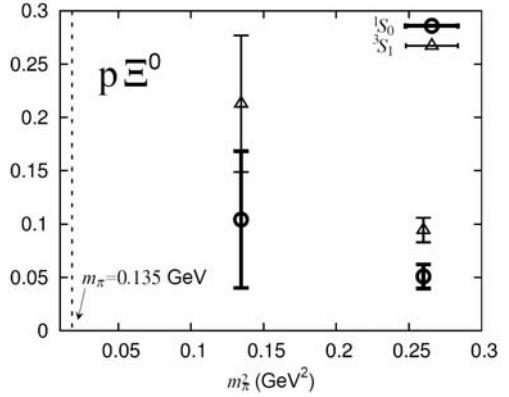


Figure 2 : The scattering lengths for  $p\Xi^0$ , in the  ${}^1S_0$  (circle) and  ${}^3S_1$  (triangle), as a function of  $m_\pi^2$ . The dashed line shows the physical point at  $m_\pi^2 = 0.135 \text{ GeV}^2$ .

## References

- [1] H.Nemura, N.Ishii, S.Aoki and T.Hatsuda, arXiv:0806.1094 [nucl-th].
- [2] M. Lüscher, Nucl. Phys. B **354**, 531 (1991).
- [3] S. Aoki, *et al.* [PACS-CS Collab.], arXiv:0807.1661 [hep-lat].



### 3.2.2. Muon Science

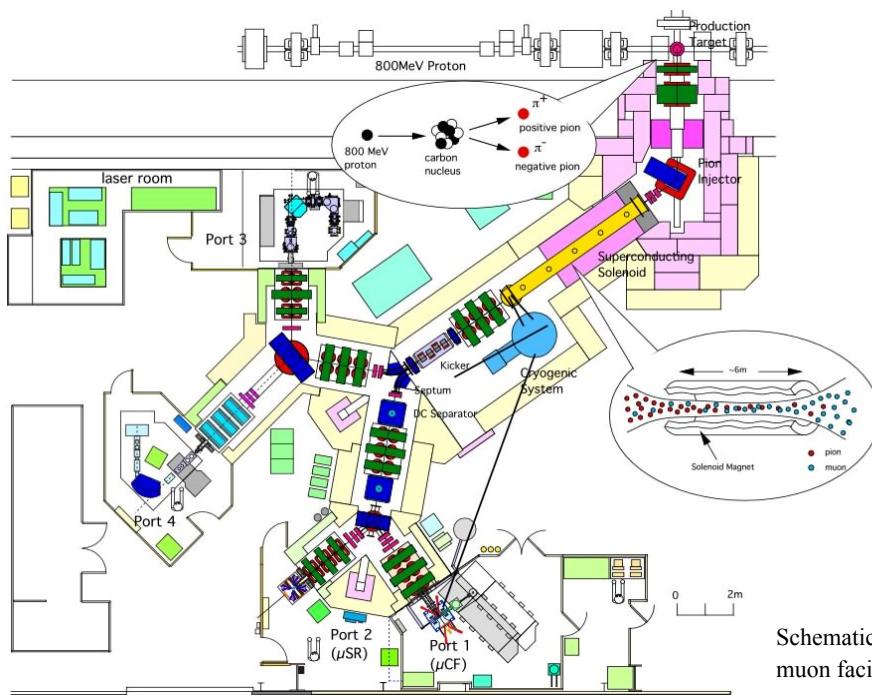
In 2002, when Iwasaki arrive at the chief scientist position of this laboratory, the RIKEN-RAL pulsed-muon facility was already in full operation mode, initiated by former chief scientist, Nagamine. Our muon facility is located inside Rutherford Appleton Laboratory in UK, attached to ISIS proton accelerator, which shared surface-muon (and pion) production target with another ISIS oriented muon facility. This production target is located upstream of the neutron production target. Our channels are located in right-hand side, and ISIS muon channels are located in left.

Our channels have four legs, and a port to each channel. Port-1 is for muon catalyzed fusion study ( $\mu$ CF), port-2 is for condensed matter physics using muon spin rotation method ( $\mu$ SR), port-3 is for R&D work for ultra-slow muon generation, and port-4 is for muonic atom x-ray measurement ( $\mu$ A\*). Recently, to enhance the capability of the facility, we installed pulsed laser system for port-2 to be able to perform pulse-laser driven  $\mu$ SR measurement, collaborating Yamanashi U., KEK and UCR. Using these four ports, we have been performing variety of experiments under international collaborations.

To make our facility competitive until full commissioning of J-PARC muon facility (MUSE), we drew up a short-term strategy, which was reviewed by our international advisory committee in 2009. They strongly endorsed two program pillars:

1. Condensed matter and molecular physics using muons
2. Ultra-slow muon development

For 1), we installed DC kicker system and new  $\mu$ SR spectrometer in port-4 so as to double the capability of our  $\mu$ SR measurement, as a member of RIKEN basic research group “Molecular Ensemble”, which is initiated by Reizo Kato a chief scientist of Condensed Molecular Materials Laboratory. For 2) we are presently developing two orders of magnitude high-intensity Lyman- $\alpha$  VUV laser together with Solid laser laboratory in RIKEN initiated by Satoshi Wada.



Schematic figure of RIKEN-RAL muon facility at ISIS in RAL

## Advanced Meson Science Laboratory Interim Review



Science & Technology Facilities Council  
Rutherford Appleton Laboratory

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Dr Yasushige Yano  
Director  
RIKEN Nishina Center for Accelerator-Based Science  
2-1 Hirosawa, Wako-shi  
Saitama, 351-0198  
Japan

13 January 2009

Dear Dr Yano,

### RIKEN-RAL Muon Facility – International Advisory Committee Review Report

I have pleasure in enclosing the report of the second meeting of the RIKEN-RAL Muon Facility International Advisory Committee. I have also sent a copy of this report to Sydney Gales, Chair of the Nishina Centre Advisory Council.

The Committee continues to endorse the findings of its first report, namely that the RIKEN-RAL Muon Facility is world-class, with unique capabilities achieved through significant technical expertise. The Committee were pleased to observe that, since its first meeting in November 2007, there have been further highly commendable developments in collaborative activity, publications and instrumentation.

The Committee recommends that, in order to continue the excellent scientific and technical outputs of the Facility, its activities are focused on two programme pillars:

1. Condensed matter and molecular physics using muons
2. Ultra-slow muon development

The Committee was highly impressed by the existing achievements in these two areas, and felt that these two pillars allow both development of the existing outstanding science programme, together with providing challenging goals for unique, world-class future developments. The Committee was keen to support the provision of sufficient financial and manpower resources to both of these areas to enable their development.

To allow continued world-class science at the RIKEN-RAL Facility focusing on these two pillars, the Committee recommends an extension of the RIKEN-RAL agreement beyond 2010 by at least another 7½ years to 2018.

In support of this recommendation, the Committee noted that conditions are appropriate for continued investment in the Facility, namely:

- a strong publication record from the Facility (including further publications since the first IAC meeting), particularly in the area of condensed matter and molecular physics;
- plans for an ambitious, unique project for the development of an ultra-slow muon source suitable for a pulsed muon facility;
- a strong user base involving many groups in Japan, desire for use of RIKEN-RAL facilities by UK researchers, and future potential for expansion of the user community in other Asian countries;



- real scientific and technical leadership – both in the area of condensed matter and molecular studies using muons, and in the area of fundamental nuclear physics using muons;
- strong support from the Japanese muon society for the RIKEN-RAL Muon Facility;
- strong support from ISIS for continuing to host the RIKEN-RAL Muon Facility.

I look forward to presenting these recommendations at the RIKEN Nishina Centre Advisory Council meeting, and to talking with you further during my visit.

It was once again a pleasure to chair the RIKEN-RAL International Advisory Committee. We particularly appreciated your hospitality during the review, and the chance to see the impressive radioactive ion beam facilities which are being developed at the Nishina centre.

Yours sincerely

Andrew Taylor

cover page of RIKEN-RAL IAC report January 2009

### 3.2.2.1. Individual topics of Muon Science

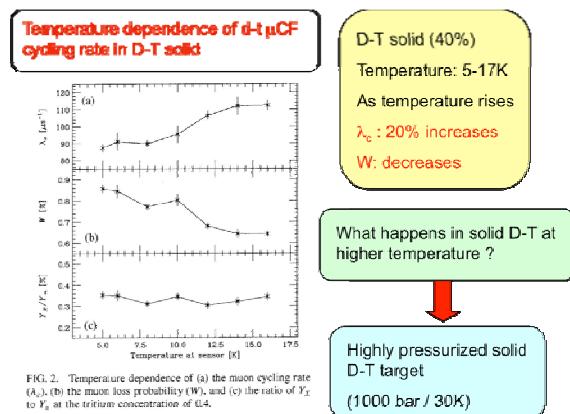
#### Muon Catalyzed Fusion with High-pressure Solid D<sub>2</sub>+T<sub>2</sub> Target toward Realization of Scientific Breakeven at the RIKEN-RAL Muon Facility

##### Abstract

We are aiming at achieving scientific breakeven with muon catalyzed d-t fusion (d-t  $\mu$ CF). The goal is to generate more than 300 d-t fusion neutrons per muon. As the first step, we are manufacturing a high-pressure solid D<sub>2</sub> target, and will commence muon catalyzed d-d fusion study to measure the d-d  $\mu$ CF cycling rate and the muon loss probability around 30K in solid D<sub>2</sub>. Then we will extend the study to d-t  $\mu$ CF with high-pressure solid D<sub>2</sub> + T<sub>2</sub> target to achieve scientific breakeven.

##### Achievement

In order to increase the d-t  $\mu$ CF cycling rate, it is required to (1) increase the dt $\mu$  formation rate, (2) decrease the muon loss probability (by increasing the reactivation rate) and (3) increase the D<sub>2</sub> + T<sub>2</sub> target density. We have observed an anomalous temperature dependence of d-t  $\mu$ CF cycling rate in solid D<sub>2</sub> + T<sub>2</sub> target [1]. The figure shows that the d-t  $\mu$ CF cycling rate ( $\lambda_c$ ) (top) increases by 20% as the solid target temperature increases from 5K to 17K. In addition, muon loss probability (W) (middle) decreases with the temperature. Possible explanations of the temperature dependence are (1) increase of the dt $\mu$  formation rate, (2) increase of the reactivation rate, (3) increase of the quasi-resonant molecular formation (3-body effect) and (4) increase of condensed matter effect (phonon contribution). Though there has been no definitive theory explaining the temperature dependence, we have decided to extend our study on d-t  $\mu$ CF experiment with solid D<sub>2</sub> + T<sub>2</sub> target to higher temperature region (up to 30K) towards realization of scientific breakeven. For this purpose, we are developing a high-pressure solid D<sub>2</sub> + T<sub>2</sub> target with the operation condition of 30K and 1,000atm. As a proto-type, we have manufactured a high-pressure solid D<sub>2</sub> target, and will make d-d  $\mu$ CF experiment to measure temperature dependence of d-d  $\mu$ CF cycling rate and muon loss probability at solid D<sub>2</sub> temperature range from 5K to 30K.



##### Publications

- 1) "Discovery of temperature-dependent phenomena of muon-catalyzed fusion in solid deuterium and tritium mixtures", N. Kawamura, K. Nagamine, T. Matsuzaki, K. Ishida, S.N. Nakamura, Y. Matsuda, M. Tanase, M. Kato, H. Sugai, K. Kudo, N. Takeda, G.H. Eaton, Phys. Rev. Lett. **90**, 043401 (2003).

Figure: Temperature dependence of d-t  $\mu$ CF parameters in solid D<sub>2</sub> + T<sub>2</sub> target

## Observation of the $d\bar{d}\mu$ Formation Rate Dependent on the Ortho-Para Concentration of D<sub>2</sub> Molecules in Muon Catalyzed d-d Fusion

### Abstract

We observed for the first time that the  $d\bar{d}\mu$  formation rate in muon catalyzed fusion is quite different for the ortho-D<sub>2</sub> and para-D<sub>2</sub> target. The effect showed an interesting dependence on the density of the D<sub>2</sub> target.

### Achievement

The muonic molecule formation rate is one of the most important parameters in muon catalyzed fusion. The rate itself limits the efficiency of fusion catalysis. In addition, the process involves rich physics of the few body problems at an unique scale. We have developed a method to prepare nearly pure D<sub>2</sub> and para-rich D<sub>2</sub> and also the detection system for d-d fusion neutrons using the d.c. muon beam at TRIUMF. With these methods, we were able to separate contributions to  $d\bar{d}\mu$  formation from ortho-D<sub>2</sub> and para-D<sub>2</sub> for the first time. The first measurement in solid D<sub>2</sub> has shown surprisingly an effect opposite to theoretical predictions. After systematic measurement for liquid and gas D<sub>2</sub> at various densities and at temperatures ranging between 3 K and 40 K, we have found that the effect is very much dependent on the target density. This indicates either a strong contribution of sub-threshold resonance at higher density or a modification of the resonance energy with the density. We plan a systematic measurement using the new high pressure D<sub>2</sub> target we are developing in order to cover a wider temperature range up to 300K.

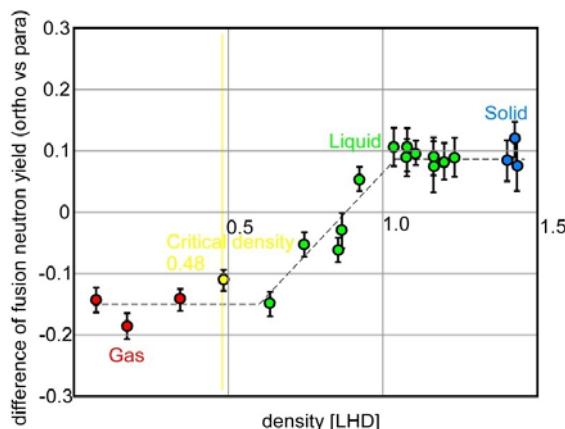


Figure: Density dependence of the difference of d-d fusion neutron yield in ortho D<sub>2</sub> and para D<sub>2</sub> around 35 K.

### Publications

- 1) "Density effect in  $d\bar{d}$  muon-catalyzed fusion with ortho and para-enriched D<sub>2</sub>", H. Imao, et al., Physics Letters B **658**, 120 (2008).
- 2) "Dependence of muon-catalyzed d-d fusion on the ortho-para ratio in solid and liquid deuterium", H. Imao, et al., Physics Letters B **632**, 192 (2006).

## X-ray Measurement from Muonic Atoms of Implanted Stable (Unstable) Nuclei in Solid D<sub>2</sub> layer for Study of Nuclear Charge Density Distribution

### Abstract

A feasibility experiment of muonic X-ray spectroscopy using solid D<sub>2</sub> layer with implanted nuclei has been performed to study nuclear charge density distribution of stable and unstable nuclei. As the demonstration data, muonic X-rays from <sup>148</sup>Sm and <sup>152</sup>Sm were successfully observed, showing characteristic spectra for spherical and deformed nuclei.

### Achievement

Negative muons ( $\mu^-$ ) are injected to solid deuterium layer after ion implantation of nuclei (A), and form muonic hydrogen atoms ( $d\mu^-$ ). Then, the  $d\mu^-$  atoms diffuse, and collide with the implanted nuclei to form muonic atoms ( $\mu^-A$ ) via muon transfer. In the formation process, characteristic muonic X-rays are generated. The muon orbit is closely located to the nuclear surface, and muonic X-ray energies are influenced by nuclear charge density distribution (nuclear shape). By measuring muonic X-ray energies and isotope energy shifts, precise information on nuclear charge density distribution is obtained. The experimental apparatus is equipped with a surface ionization source to produce alkali, alkali-earth and rare-earth (Sm, Nd) ions for implantation to solid D<sub>2</sub> layers. For example, samarium isotopes show very abrupt changes in their nuclear characteristics from spherical to deformed nuclei. <sup>144</sup>Sm is magic in neutrons (N = 82) and display the characteristics of a stiff spherical nucleus, whereas <sup>152</sup>Sm and <sup>154</sup>Sm reveal the characteristic of deformed nuclei whose muonic X-ray spectra are expected to show a 2p hyperfine structure. As shown in figure, a successful observation of muonic X-rays from <sup>148</sup>Sm and <sup>152</sup>Sm was achieved. The obtained muonic X-ray energies, isotope shift and 2p hyperfine splitting agree with the existing data, and will be compared with theoretical calculation including nuclear physic parameters.

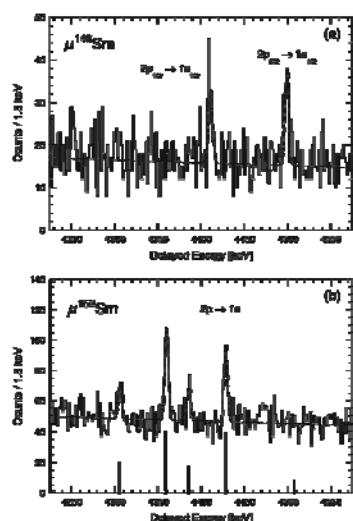


Figure: Delayed energy spectrum of the 2p $\rightarrow$ 1s muonic transitions measured with 1-mm pure D<sub>2</sub> layer implanted of about 1 ppm of <sup>148</sup>Sm and <sup>152</sup>Sm, respectively.

### Publication

- 1) "Muon spectroscopy with trace alkaline- earth and rare-earth isotopes implanted in solid D<sub>2</sub>", P. Strasser et al., Hyperfine Interact **193**, 121 (2009).

## Precision Measurement of the Positive Muon Lifetime ( $\tau_{\mu^+}$ ) with a Highly Intense Muon Beam at RIKEN-RAL

### Abstract

The  $\tau_{\mu^+}$  is associated with the Fermi coupling constant ( $G_F$ ), which is one of the fundamental constants in the Standard Model. We developed a new experimental method for intense pulsed beam and obtained  $\tau_{\mu^+}$  in 51 ppm precision. The result is consistent with previous measurements. The established method is an important milestone to a new measurement with intense pulsed muon beam.

### Achievement

The Standard Model (SM) requires three experimental input parameters ( $\alpha$ ,  $M_Z$  and  $G_F$ ). Muon is the best probe for precision measurement of  $G_F$  in a ppm level, which means that this experiment is the precision test of the SM.

We established the new method named “multi-decay per one time window method”: we used the intense pulsed muon beam for the first time to observe the large number of muon decay ( $\sim 1.15 \times 10^{10}$ ) in a low background ( $B/S \sim 5 \times 10^{-5}$ ) with highly segmented detector (MWPC, 192ch) and developed associated new techniques such as the target system, the clock system and DAQ. In the offline analysis, we also developed a new numerical correction method to solved pile-up problem. The results were

$$\begin{aligned} \tau_\mu &= 2197.01 \pm 0.11 & +0.006 \\ &\quad \text{stat.} & -0.034 &\quad \text{syst.} & (51 \text{ ppm}) \\ G_F &= (1.166372 \pm 0.000029) \times 10^{-5} & \text{stat. + syst.} \end{aligned}$$

This result is consistent with the previous measurements. This “RIKEN-RAL pulsed method” is an important milestone for the future muon lifetime measurement with a very high intense pulsed muon beam.

### Publications

- 1) “Precise muon lifetime measurement with a pulsed beam at the RIKEN-RAL muon facility”,  
D. Tomono, S.N. Nakamura, et. al., Nucl. Phys. B **149**, 341 (2005).
- 2) “Precise measurement of the positive muon lifetime and test of the exponential decay law”,  
D. Tomono, S.N. Nakamura, et. al., Journal of Physics G **29**, 2013(2003).
- 3) “Muon Lifetime measurement at RIKEN-RAL and prospects at an intense muon source”,  
D. Tomono, S.N. Nakamura, et. al., Nuclear Instruments and Methods A **503**, 283(2003).

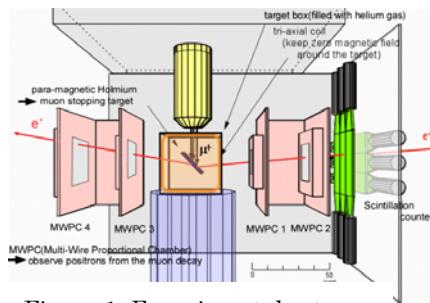


Figure 1. Experimental setup.

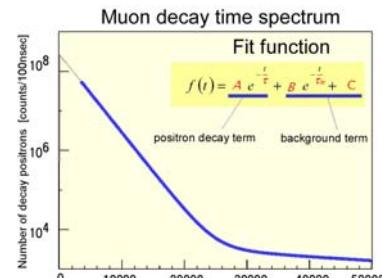


Figure 2. Muon decay time spectrum and fitting function

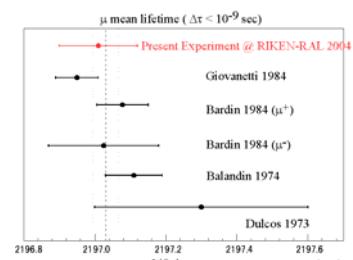


Figure 3. Lifetime precision in the past 30 years.

## Observation of Quasi-Static Internal Fields in the Normal State of the Cu-based Superconductor, $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$

### Abstract

Quasi-static internal fields at the muon site have been observed from precise ZF- $\mu$ SR measurements on the high- $T_c$  superconductor,  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ , in the normal state of around 100 K. The quasi-static internal fields would be due to the theoretically predicted microscopic orbital circular currents caused by the formation of the spin-gap state. Our observation indicates that the orbital circular current would be dynamically fluctuating.

### Achievement

We have carried out precise  $\mu$ SR on the high- $T_c$  superconductor  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  in the normal state around 100 K accumulating higher statistics of muon events rather than usual. Tiny changes of the time spectrum have been revealed to be due to changes of the dynamics of an additional internal field at the muon site. The magnitude of the additional internal field was around a couple of Gauss and the fluctuation frequency was the order of 100 kHz. This quasi-static internal field appears around 100 K as shown in Fig. (a) and (b). The current study would prove the appearance of the theoretically predicted microscopic circular current caused by the formation of the spin-gap state in the high- $T_c$  superconductors and indicate that the predicted circular current is dynamically fluctuating with the low frequency of about 100 kHz.

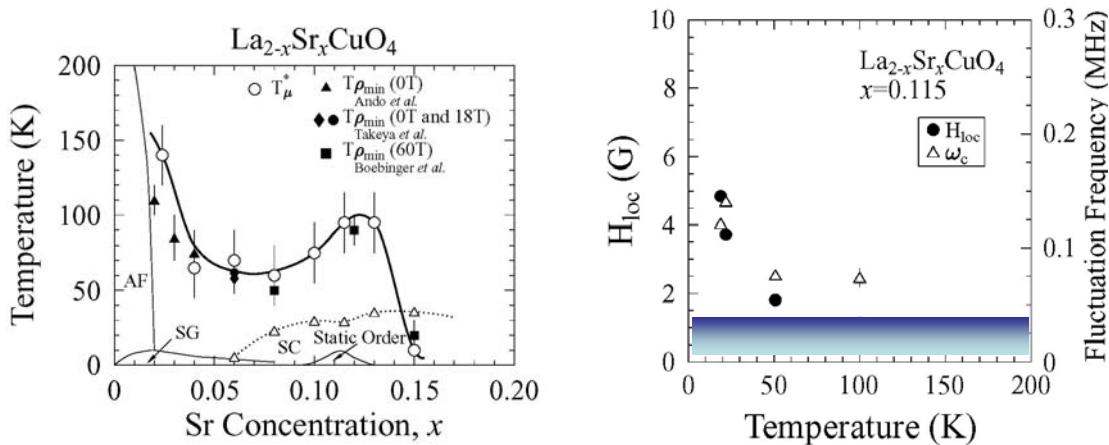


Figure: (a) Phase diagram determined from the current precise ZF- $\mu$ SR study. Open circles show temperatures where the quasi-static internal field appear. (b) Temperature dependences of the magnitude and frequency of the quasi-static internal field in the case of  $x=0.115$ .

### Publications

I. Watanabe *et al.*, J. Phys. Soc. Jpn. **77**, 124716 (2008).

## First Indication of a Relationship between Magnetic Correlations and Superconducting State in Fe-Based Superconductors

### Abstract

A magnetic phase diagram of the Fe-based superconductor,  $\text{SmFeAsO}_{1-x}\text{F}_x$ , has been firstly suggested from ZF- $\mu$ SR measurements on the collaboration between RIKEN and ISIS muon groups. A possible co-existing state between magnetically ordered and superconducting states has been suggested between  $x=0.10$  and  $0.15$ .

### Achievement

The Fe-based superconductor has been launched from Japan to open a new field of superconductors. We have preformed ZF- $\mu$ SR measurements on  $\text{SmFeAsO}_{1-x}\text{F}_x$  in order to investigate the carrier concentration dependence of a magnetically ordered state and the superconducting state. Figure (a) shows the crystal structure of  $\text{SmFeAsO}_{1-x}\text{F}_x$  and Figure (b) shows the magnetic phase diagram suggested from our current study. There is a clear region between  $x=0.10$  and  $0.15$  where the superconducting state and the magnetically ordered state possibly show the co-existing state. This result shows the first indication of the possible co-existing state between the magnetically ordered state and superconducting state in the Fe-based superconductor. We have shown from the current study that a similar discussion on the origin of the superconducting state to that of high- $T_c$  Cu-base superconductors would be applied to Fe-based superconductors as well.

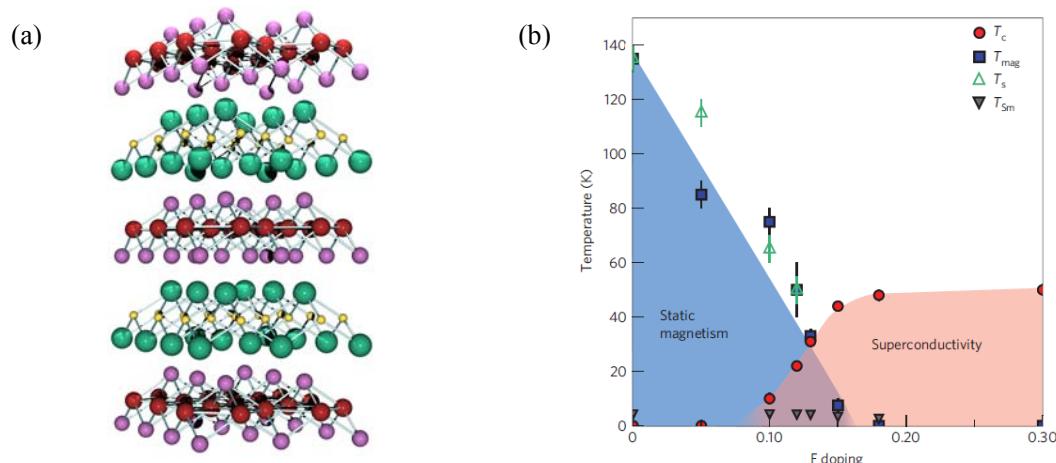


Figure: (a) Crystal structure of  $\text{SmFeAsO}_{1-x}\text{F}_x$ . (b) Magnetic phase diagram suggested from the current ZF- $\mu$ SR study.

### Publications

- 1) A.J. Drew *et al.*, Nature Material **8**, 310 (2009).
- 2) A.J. Drew *et al.*, Phys. Rev. Lett. **101**, 0970101 (2008).

## μSR Study around a Quantum Critical Point --- Magnetism and Superconductivity ---

### Abstract

One of the most exciting topics in modern condensed matter physics is the interplay between different electronic ground states near a quantum critical point (QCP). Especially, studying a relationship between magnetism and superconductivity in the vicinity of QCP is very important to understand the mechanism of superconductivity, i.e., competitive or corporative. We have performed  $\mu$ SR measurements in following superconductors in order to elucidate their mechanisms; (i) FeAs-based superconductors: LaFeAs(O,F), LaFeAs(O,H), (Ba,K)Fe<sub>2</sub>As<sub>2</sub>, Ba(Fe,Co)<sub>2</sub>As<sub>2</sub>; (ii) Antiperovskite-type superconductor: ZnNNi<sub>3</sub>; (iii) Heavy fermion systems: CeCo(In,Cd)<sub>5</sub>, (Ce,La)<sub>2</sub>IrIn<sub>8</sub>, Ce<sub>2</sub>Rh(In,Sn)<sub>8</sub>

### Achievement

Figure shows the temperature dependence of the muon spin relaxation rate  $\sigma_v$ , which corresponds to the inverse squared in-plane magnetic penetration depth  $\lambda$  ( $\sigma_v \propto \lambda^{-2}$ ), in Ba<sub>0.75</sub>K<sub>0.25</sub>Fe<sub>2</sub>As<sub>2</sub> under  $H = 500$  Oe. The lines represent the known *s*- and *d*-wave superconducting gap behaviors and the result of a fit to the two gap model, which is assuming two independent contributions to the total  $\sigma_v$ . It is obvious that two gap model well reproduced the data, while the data shows significant departure from the simple *s*- and *d*-wave curves. Therefore, the two gap scenario is qualitatively consistent with our experimental data to explain the superconducting pairing mechanism.

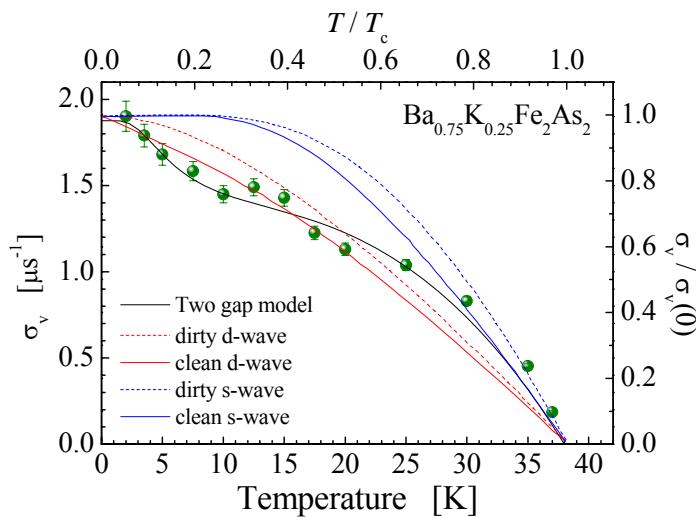


Figure: Temperature dependence of muon spin relaxation rate  $\sigma_v$  versus Temperature  $T$  for Ba<sub>0.75</sub>K<sub>0.25</sub>Fe<sub>2</sub>As<sub>2</sub>. The figure shows experimental data (green circles with error bars) and theoretical fits for different superconducting gap behaviors: Two gap model (solid black line), dirty d-wave (dashed red line), clean d-wave (solid red line), dirty s-wave (dashed blue line), and clean s-wave (solid blue line). The x-axis is Temperature  $T$  [K] from 0 to 40, and the y-axis is  $\sigma_v$  [μs<sup>-1</sup>] from 0.0 to 2.0. The two-gap model fits the data well, while single-gap models deviate significantly at higher temperatures.

### Publications

- 1) "Development of the heavy-fermion state in Ce<sub>2</sub>IrIn<sub>8</sub> and the effects of Ce dilution in (Ce<sub>1-x</sub>La<sub>x</sub>)<sub>2</sub>IrIn<sub>8</sub>", K. Ohishi *et al.*, Phys. Rev. B **80**, 125104/1-7 (2009).
- 2) "Magnetism and Superconductivity in Heavy Fermion Superconductor CeCo(In<sub>0.97</sub>Cd<sub>0.03</sub>)<sub>5</sub>", K. Ohishi *et al.*, Physica B **404**, 754-756 (2009).
- 3) "Quasiparticle excitations in newly discovered antiperovskite superconductor ZnNNi<sub>3</sub>", K. Ohishi *et al.*, Physica C in press.

## Quantum Critical Behavior and Soft Mode in the Randomness-introduced Quantum Spin Systems Probed by LF- $\mu$ SR Measurements

### Abstract

Longitudinal-field muon-spin-relaxation (LF- $\mu$ SR) measurements were carried out on the randomness-introduced quantum spin systems. The relative temperature change of the muon-spin-relaxation rate  $\lambda$  in longitudinal-fields, which corresponds to the wave-vector integration of the generalized dynamical susceptibility, were deduced from LF- $\mu$ SR measurements.

### Achievement

In the mixed system  $Tl_{1-x}K_xCuCl_3$ , the spatial randomness of the local chemical potential is introduced through the difference of the value of the dominant intradimer interaction between  $TlCuCl_3$  and  $KCuCl_3$ . Magnetization measurements suggest that the ground state is a magnetic state with finite susceptibility in the mixed system in zero field, although finite excitation gap remains. To investigate microscopic dynamical magnetic properties in highly random systems, we carried out detailed longitudinal-field muon-spin-relaxation measurements in  $Tl_{1-x}K_xCuCl_3$  with  $x = 0.60$  single crystals. Deduced relative temperature change of the muon spin relaxation rate  $\lambda$  in each longitudinal field is shown in Fig. 1. The temperature where the peak is observed decreases with decreasing the magnetic field. The muon-spin-relaxation rate  $\lambda$  in the longitudinal field corresponds to the wave-vector integration of the generalized dynamical susceptibility. In other word, the longitudinal field ( $H_{LF}$ ) dependence of  $\lambda$  corresponds to the frequency ( $\omega_{LF} = \gamma_\mu H_{LF}$ ) spectrum of spin fluctuations. Therefore, observed peak shift to lower temperatures with decreasing  $H_{LF}$  is the observation of the slowing down of Cu-3d spins fluctuation frequency, and is interpreted as the soft mode of spin waves toward a possible magnetic phase transition.

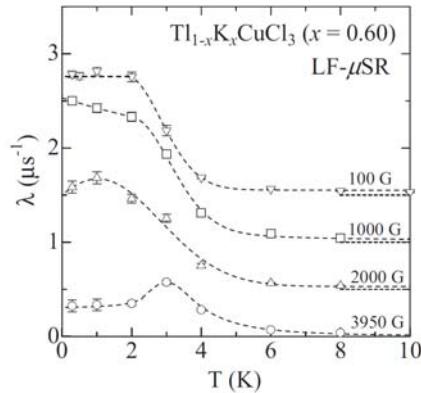


Figure 1: Relative temperature change of  $\lambda$  in each longitudinal field.

### Publications

- 1) "Quantum critical behavior in highly random systems  $Tl_{1-x}K_xCuCl_3$  probed by zero- and longitudinal-field muon-spin-relaxation measurements", T. Suzuki, F. Yamada, T. Kawamata, I. Watanabe, T. Goto, and H. Tanaka, Phys. Rev. B **79**, 104409-(1-5) (2009).
- 2) "Muon spin relaxation detection of the soft mode toward the exotic magnetic ground state in the bond-disordered quantum spin system IPA-Cu( $Cl_{0.35}Br_{0.65}$ )<sub>3</sub>", T. Goto, T. Suzuki, K. Kanada, T. Saito, A. Oosawa, I. Watanabe, and H. Manaka, Phys. Rev. B **78**, 054422-(1-6) (2008).

## μSR Study of the Spin-Lattice Cooperative Phenomenon of RbCoBr<sub>3</sub>

### Abstract

We carried out μSR experiments of RbCoBr<sub>3</sub> whose results are reproduced by a spin-lattice model considering a new concept “lattice frustration”. As a result, it is found that the magnetically ordered state between  $T_{N1}$  and  $T_{N2}$  is a ferrimagnetic phase but not the partial disordered phase suggested by the spin-lattice model. It is likely that the suitable parameters are not determined in this model.

### Achievement

RbCoBr<sub>3</sub> is a very unique frustration compound with spin and lattice frustration suggested by a new spin-lattice model which is introduced in the frustration of not only spin systems but also lattice systems, say “lattice frustration”. In order to investigate the spin-lattice cooperative phenomenon of RbCoBr<sub>3</sub>, we carried out μSR experiments.

Clear muon spin rotation was observed below the partial disordered (PD) phase transition temperature  $T_{N1} \sim 37$  K. Temperature dependence of the precession frequency obtained by Fourier analysis of the spectra is shown in Figure. It is found that there are two frequency components below  $T_{N1}$ . Our calculation of the dipole field from the spins on Co<sup>2+</sup> ions in PD phase shows only one frequency component. On the other hand, it is predicted by the spin-lattice model that a two-sublattice ferrimagnetic (2FR) phase appears in a very small temperature region around  $T_{N1}$ . Our field calculation for the 2FR phase shows two frequency components. Accordingly, the magnetically ordered state between  $T_{N1}$  and  $T_{N2}$  in RbCoBr<sub>3</sub> is the 2FR phase but not the PD phase. There is a possibility that the parameters in the spin-lattice model are not suitable.

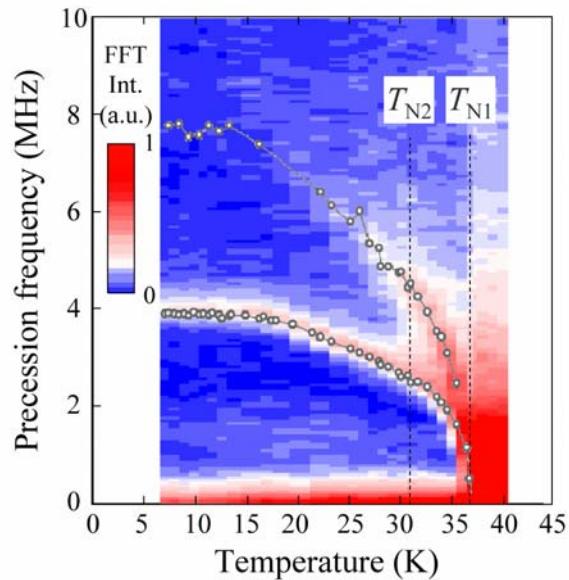


Figure: Temperature dependence of the precession frequency obtained by Fourier analysis. Open circles indicate the peak center frequency.

### Publications

- 1) “μSR study of the spin-lattice cooperative phenomenon of RbCoBr<sub>3”</sub>, T. Kawamata *et al.*, RIKEN Accel. Prog. Rep **42**, 243 (2009)

## High-Pressure $\mu$ SR Studies of Molecular-Based Materials

### Abstract

Magnetic properties of molecular based materials have been investigated by means of ambient- and high-pressure  $\mu$ SR measurements. In the case of an organic metal  $(\text{DMe-DCNQI})_2\text{Cu}$ , pressure-induced magnetic ordering has been revealed and pressure-temperature phase diagram is established. We also developed a high-pressure experimental technique.

### Achievement

An organic metal,  $(\text{DMe-DCNQI})_2\text{Cu}$  ( $\text{DMe-DCNQI} = 2,5\text{-dimethyl-DCNQI}$ ), shows an unusual pressure effect. The hybridization between the wide 1D  $2p\pi$  bands and the narrow  $3d$  bands is a key factor in understanding such unconventional electronic properties of this compound. At ambient pressure, this material shows metallic behavior down to 450 mK. Peculiar to  $(\text{DMe-DCNQI})_2\text{Cu}$ , an insulating phase is induced by the application of pressure higher than 100 bar. This unusual  $P$ - $T$  phase diagram can be reproduced by the chemical pressure effect using selectively deuterated compounds. The fully deuterated sample of  $(\text{DMe-DCNQI})_2\text{Cu}$ , in which the chemical pressure corresponds to 512 bar, exhibits the antiferromagnetic ordering below 8 K.

Recently, we have developed a high-pressure  $\mu$ SR setup for the RIKEN-RAL Muon Facility and successfully observed a sign of pressure-induced magnetic ordering of  $(\text{DMe-DCNQI})_2\text{Cu}$ , which was predicted by the chemical pressure study, using this high-pressure setup.

Figure shows typical zero-field  $\mu$ SR time spectra at an actual pressure of 500 bar. At a low-temperature of 2 K, significant decrease of initial asymmetry due to static magnetic ordering is recognized. This decrease of the asymmetry is observed below 8 K, which indicates this compound is an antiferromagnet with  $T_N = 8$  K. This  $T_N$  value is consistent with the chemical pressure study.

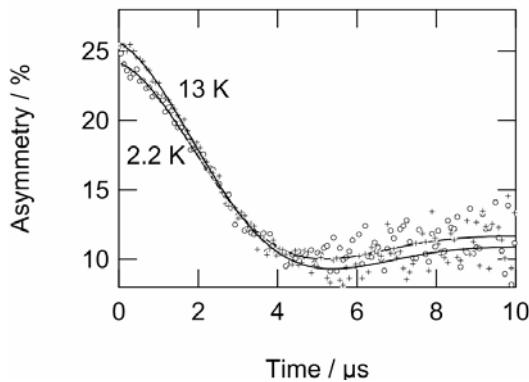


Figure: Zero-field  $\mu$ SR time spectra of  $(\text{DMe-DCNQI})_2\text{Cu}$  at 500 bar for selected temperatures.

### Publications

- 1) "High pressure magnetic study on a molecular conductor,  $(\text{DMe-DCNQI})_2\text{Cu}$ ", Y Ishii et al., RIKEN Accel. Prog. Rep. **42**, 254 (2009)

## **μSR Study of Structure Dependent Electron Radical Dynamics in Poly(3-alkylthiophene)**

### **Abstract**

Longitudinal field muon-spin-relaxation (LF-μSR) measurements have been performed for poly(3-alkylthiophene) to elucidate directly the intra- and inter-chain hopping mechanisms. The present results have revealed the remarkable shift of relative dominance between the intra- and inter-chain charge transport which is depending on their regio-regularity and side chain length.

### **Achievement**

LF dependences of the muon-spin depolarization rate ( $\lambda_1$ ) in regio-regular (RR) and regio-random (Rdm) poly(3-hexylthiophene) (P3HT) and RR-poly(3-octylthiophene) (RR-P3OT) show  $H^{-0.5}$  field dependent characteristic of intra-chain diffusion in low temperatures which is turns over to inter-chain diffusion as characterized by its field-dependence of  $\lambda_1$  of C-H<sup>0.5</sup> curve. The initial sign of change was suggested to be around the low temperature of 25 K, 50 K and 50 K for RR-P3HT (b), Rdm-P3HT (c) and RR-P3OT (d), respectively. In the P3HT samples, it is apparent that inter-chain diffusion process requires the assistance of higher temperature for regio-random system which may be related to different values of bandgap between them. The regio-regular structure is known to have a small bandgap (1.7 eV) which is 0.4 eV smaller than that of the regio-random structure. With the smaller bandgap, thermal excitation energy needed to support the inter-chain charge transport is supposedly smaller. Comparison in the difference side chain length of the alkyl between RR-P3HT and RR-P3OT, it found that inter-chain polaron diffusion process in RR-P3OT, requires the assistance of higher temperature. It may nonetheless be related to the difference distance between the chains. Clearly, longer side chain leads to larger distance between the chains which imply in turn the need of large thermal energy to support the inter-chain hopping process.

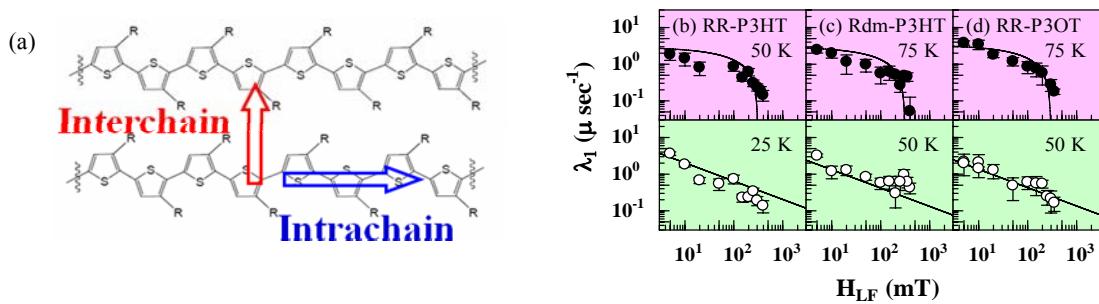


Figure: The scheme of direction of charge motion intra-chain and inter-chain (a) and LF dependences of depolarization rate ( $\lambda_1$ ) in RR-P3HT (b), Rdm-P3HT (c) and RR-P3OT(d).

### **Publications**

- 1) " Intra- and Inter-Chain Polaron Diffusion in Regio-random Polythiophene Studied by Muon Spin Relaxation ", Risdiana et al., Physica B doi:10.1016/j.physb.2010.01.080 (2010).

## Muons for Spintronics: the New Muonium Method Detecting Conduction Electron Spin Polarization (CESP) in n-type GaAs

### Abstract

The spin-dependent exchange scattering between an electron in a triplet-state muonium (Mu) and spin-polarized conduction electrons has been observed in Si doped n-type GaAs under low temperature. This new measurement technique could be applied to other important semiconductors, such as Si, Ge, and C, to study the behavior of their electron spin, where the traditional optical method could not be applied because of the weak spin-orbit coupling.

### Achievement

As shown in Fig. 1, the circularly polarized laser light photoexcites the conduction electron spin throughout the sample by the below bandgap excitation. The muons implanted from the other side of the sample forms the Mu, which then interacts with the CESP by the spin exchange scattering. Since the muon asymmetry signal from the singlet Mu disappears immediately, whereas one from the triplet Mu stays the same, the induced CESP should give us a different  $\mu$ SR time spectrum depending on the direction of the CESP. The typical  $\mu$ SR spectrum has a sharp drop at the timing of the laser illumination due to the additional singlet Mu formed by interaction with the photoexcited electrons with their spin oriented anti-parallel to the muon spin. If the laser light is circularly polarized, the  $\mu$ SR spectra show additional small difference, depending on the light helicity (i.e. the direction of the CESP), as shown in Fig. 2. This effect has been confirmed by changing the temperature and pump light intensity so that our results are consistent with the previous works on the same material.

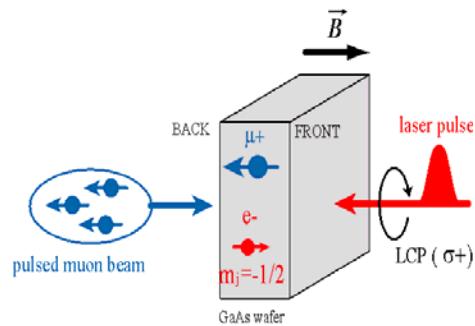
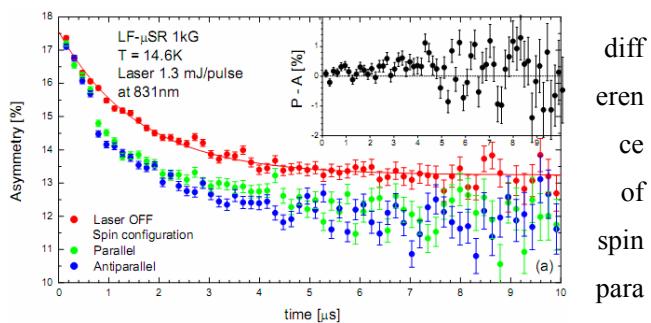


Fig. 1: Schematic view of the laser- $\mu$ SR experiment. “LCP” stands for “Left Circularly Polarized”.

Fig. 2:  $\mu$ SR time spectrum. The inset shows the



(green circles) and anti-parallel (blue circles) configuration.

### Publications

- “Muons for spintronics: Photo-induced conduction electron polarization in n-type GaAs observed by muonium method”, K. Yokoyama, K. Nagamine, K. Shimomura, H.W.K. Tom, R. K. Kawakami, P. Bakule, Y. Matsuda, F.L. Pratt and E. Torikai, Physica B **404**, 856-858 (2009).

## New Measurements of the Chemical Reaction Rate of Muonium with Stimulated Raman-Pumped H<sub>2</sub>\*(v=1)

### Abstract

The reaction rate of Mu+H<sub>2</sub> gives a rigorous test of the chemical reaction rate theory because a highly accurate H<sub>3</sub> potential energy surface is available. We used stimulated Raman pumping (SRP) to produce H<sub>2</sub> in its first vibrational state H<sub>2</sub>\*(v=1) and measured, for the first time, the reaction rate of Mu + H<sub>2</sub>\*(v=1)  $\rightarrow$  MuH + H.

### Achievement

We used the 2nd harmonic output of a Nd:YAG laser (532 nm) to pump the first vibrational state H<sub>2</sub>\*(v=1) (see bottom left figure). In the bottom right figure of the difference spectrum of μSR, the muonium precession amplitude is increasing with time because the precession dumps faster with laser ON compared with laser OFF (thus the observed phase is reversed compared with the raw MuSR spectra). We obtained the muonium reaction rate from the difference in the dumping rate. The obtained reaction constant  $k_{\text{Mu}} = 10.7 \pm 2.4 \times 10^{-13} \text{ cm}^3 \text{s}^{-1}$  is in excellent agreement with theory  $9.8 \times 10^{-13} \text{ cm}^3 \text{s}^{-1}$ . This enhancement of as much as 8 orders of magnitudes compared with the rate from H<sub>2</sub> ground state is due to the decreased barrier height. We are working on reducing the error.

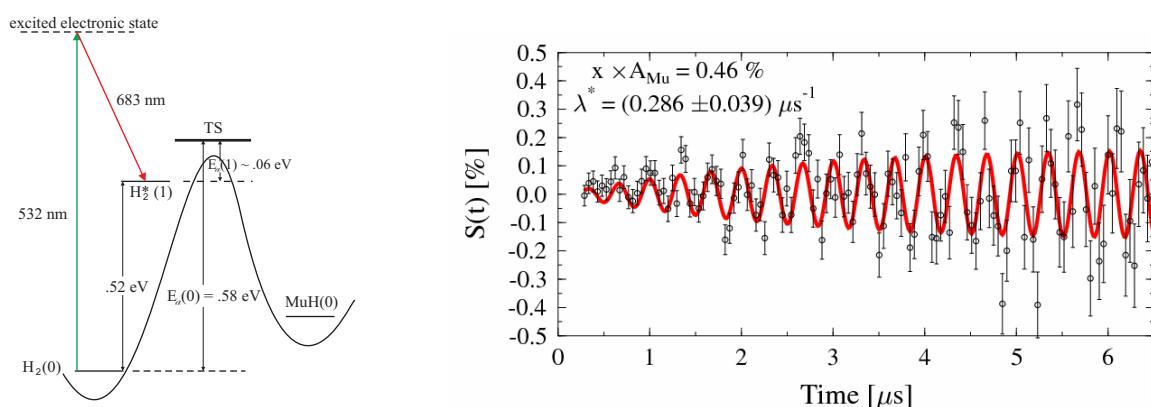


Figure: Scheme of Mu+H<sub>2</sub> reaction and laser excitation (left) and the difference μSR signal with laser OFF subtracted from laser ON (right)

### Publications

- 1) "Toward the First Study of Chemical Reaction Dynamics of Mu with Vibrational-State-Selected Reactants in the Gas Phase: the Mu+H<sub>2</sub>\*(v=1) Reaction by Stimulated Raman Pumping", P. Bakule et al., Physica B 404, 5-7 (2009).

## Domestic Collaborations for $\mu$ SR on the Material Science

### Abstract

We have been continuing to collaborate with more than 50 domestic research groups concerning the material science using  $\mu$ SR. On the basis of these collaborations we have produced more than 170 papers and grown more than 20 master-course students and 10 doctor course students.

### Achievement

Figure shows a map of collaborations on the basis of domestic activities at the RIKEN-RAL Muon Facility. We have been collaborating with more than 50 groups in Japan. Those collaborations have created more than 170 papers in a wide range of the material science. Also we have grown master-course students and doctor-course students through those collaborations. We have grown up new users of  $\mu$ SR from the zero-base level and made user cores in several universities. Those new user cores are now important as core groups who can support  $\mu$ SR activities in the J-PARC in near future.

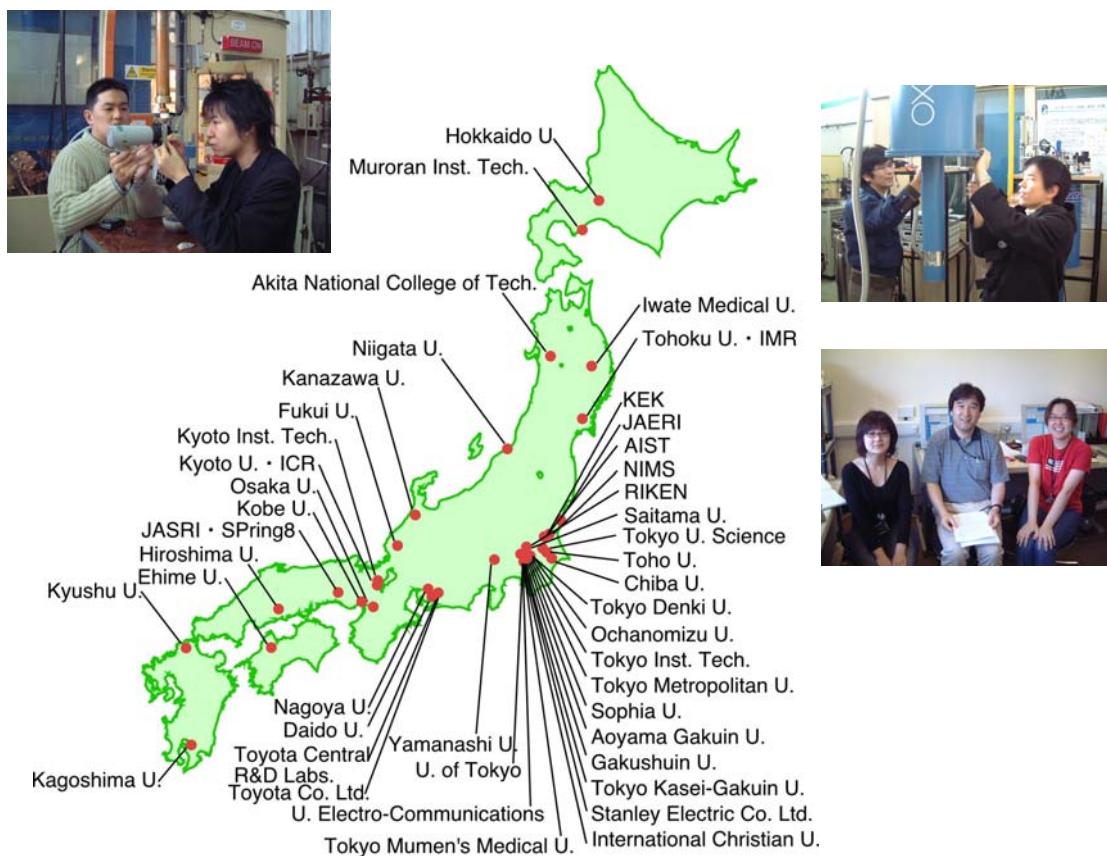


Figure: Domestic collaboration map with more than 50 groups for the material science.

### Publications

See the publication list at the end of this report.

## International Collaborations for $\mu$ SR on the Material Science and Creation of New Activities of $\mu$ SR in Asian Areas

### Abstract

We have collaborated with 8 foreign research groups concerning the material science using  $\mu$ SR. In addition, we have made a Memory of Understandings (MOU) with three Indonesian universities on the basis of the material sciences using  $\mu$ SR at the RIKEN-RAL Muon Facility. .

### Achievement

We have collaborated with 8 foreign research groups using  $\mu$ SR at the RIKEN-RAL Muon Facility for the material sciences. Recently, we are keen to create new  $\mu$ SR activities in the Asian area, because there has been no  $\mu$ SR core user in the Asian area except for Japan. On the basis of this direction of expansion of the  $\mu$ SR activity at the RIKEN-RAL Muon Facility, we have made MOU in 2008 with three Indonesian national universities, which are Institute Technology of Bandung (ITB), Padjadjaran University (UNPAD) and Institute Technology of Surabaya (ITS). Left-hand photo shows the sign-up ceremony with the former director of the Nishina Center, Dr. Yano, and the former rector of ITB, Prof. Santoso, shaking hands to celebrate MOU. Through this collaboration, researchers are being exchanged and symposiums are being organized. We are planning to expand these activities to other Asian countries aiming to organize Asian  $\mu$ SR activities on the basis of the RIKEN-RAL Muon Facility and J-PARC in the near future.



Left-hand photo: MOU signed by Dr. Yano (left) and Prof. Santoso (right).

Right-hand photo: Prof. Nugroho (right) from ITB enjoying  $\mu$ SR at the RIKEN-RAL Muon Facility with his student (left).

### Publications

- 1) A.A. Nugroho *et al.*, Physica B **404**, 785-788 (2009).

## Development of a Gas-Pressurized High-Pressure Setup for $\mu$ SR Experiments at the RIKEN-RAL Muon Facility

### Abstract

A gas-pressurized high-pressure  $\mu$ SR setup for RIKEN-RAL has been developed under the collaboration between the RIKEN  $\mu$ SR group and the ISIS high-pressure group. The system is pressurized up to 6.4 kbar and cooled down to 2 K using the existing cryostat. The developed system is the world-wide **FIRST** high-pressure  $\mu$ SR system for the pulsed muon beam.

### Achievement

A high-pressure  $\mu$ SR setup which is pressurized by gas He up to 6.4 kbar has been developed for the RIKEN-RAL Muon Facility. This work has been done under a close collaboration between RIKEN  $\mu$ SR group and the ISIS high-pressure group. The system can be cooled down to 2 K using the existing cryostat. The advantage of this system is that a homogeneous pressure can be applied to a sample from the room temperature down to 2 K and that the pressure can be changed smoothly from zero to the maximum without removing the sample from the high-pressure cell. Figure (a) shows the developed high-pressure cell and Figure (b) shows a high-pressure rotary pump and a gas intensifier. The gas He is pressurized up to 2.7 kbar using the high-pressure rotary pump and finally pressurized up to 6.4 kbar using the intensifier. This system is the **FIRST** high-pressure  $\mu$ SR system in the world for the pulsed muon beam. Thus, it is expected to open a new range of applications of  $\mu$ SR for material sciences.

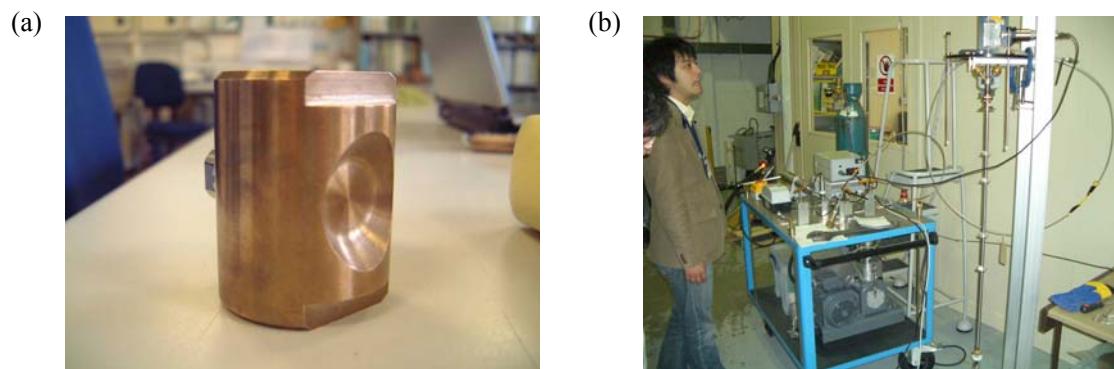


Figure: (a) Developed pressure cell made by CuBe. (b) The high-pressure  $\mu$ SR setup. The high-pressure cell is connected to a high-pressure stick to introduce pressurized gas He.

### Publications

- 1) I. Watanabe *et al.*, Physica B **404**, 993 (2009).
- 2) R. Done *et al.*, Proceedings of 12th International Conference on Pressure Vessel Technology (ICPVT-12), 329 (2009).

## Development of New Laser System for Laser-Irradiated Pump-Probe Type μSR Experiments

### Abstract

Use of laser pulses synchronized with muon implantation provides an extremely useful tool that extends the scope of μSR. A laser system for μSR setup has been developed at RIKEN-RAL.

### Achievement

The laser irradiation will cause coherent electron excitation or ionization in the studied sample or even affect the muonium state directly. Additionally, circularly polarized light can induce spin polarization of the excited electrons in the sample. The effects of the pump laser pulse on the sample can then be sensitively probed using μSR technique. To allow routine use of such pump-probe technique we have built a dedicated laser room and laser beam delivery system next to the ARGUS μSR spectrometer (bottom left figure). The laser system was built by pooling equipment together in close collaboration between RIKEN, KEK, Yamanashi University, University of California and ISIS and is based on an optical parametric oscillator (OPO) pumped by a Nd:YAG laser. The output wavelength from OPO is tunable between 400 nm and 2500 nm. Six groups has already conducted μSR involving laser on topics such as spintronics in semiconductors and reaction dynamics of muonium with excited H<sub>2</sub> molecules (bottom right figure).

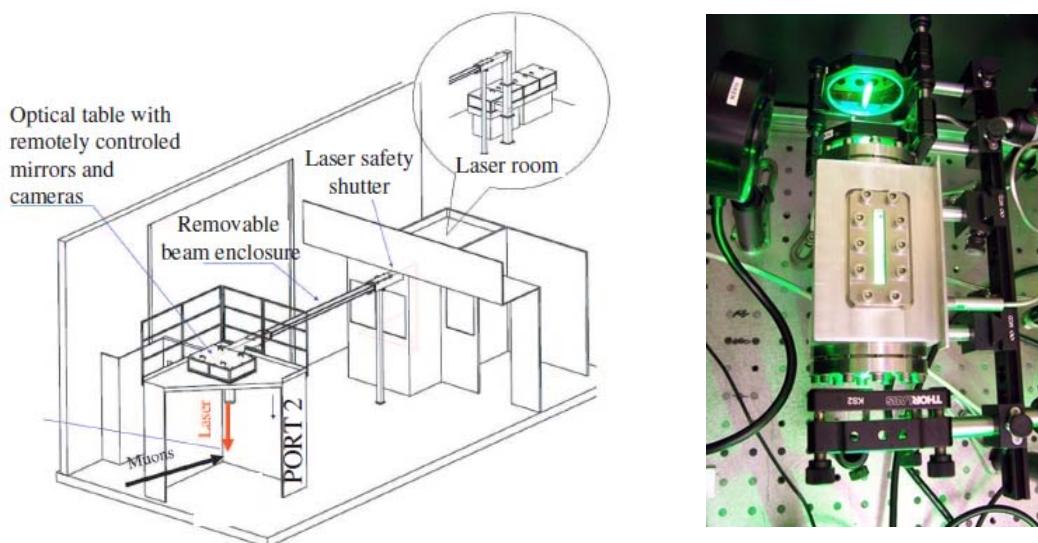


Figure: (Left) Layout of the laser enclosure and beam delivery to μSR sample, (Right) Gas H<sub>2</sub> chamber being pumped by 532nm laser light.

### Publications

- 1) "Installation of a new laser system for laser-irradiated pump-probe type μSR experiments at RIKEN-RAL muon facility", P. Bakule et al., RIKEN Accel. Prog. Rep. **42**, 233 (2009).

## Development of the New Multi-Channel Spectrometer (CHRONUS) for $\mu$ SR Studies at Port-4, RIKEN-RAL

### Abstract

A new multi-channel  $\mu$ SR spectrometer is newly developed at Port-4. The 606 fine segmented detectors and homogeneous magnetic field enable us to perform  $\mu$ SR experiments under various extreme conditions using the intense pulsed beam. The installation and the beam commissioning were almost completed and practical  $\mu$ SR experiments commence to operate.

### Achievement

A research project to explore various matters under multiple extreme conditions combined with  $\mu$ SR technique are going on. In order to promote the project more intensively and efficiently, a high performance and multi-channel spectrometer named **CHRONUS** is newly installed at Port-4, which enables us to perform experiments in parallel at Port-2 and Port-4.

We employed new technologies for the spectrometer: magnets with good field homogeneity ( $<0.1\%$ ) (Fig.1) and with a large gap of 150mm, the 606 highly segmented and direction-sensitive  $\mu$ -e counters (Fig.2), the counter mold (Fig.2 top) and frame, fiber light guides (Fig.2 bottom), multi-anode photomultiplier tubes covered with magnetic field shields (Fig.1) and a new data acquisition system (Fig.3). At present, Beam commissioning is in progress to prove availability for practical  $\mu$ SR measurements. We are going to start operation in 2010.

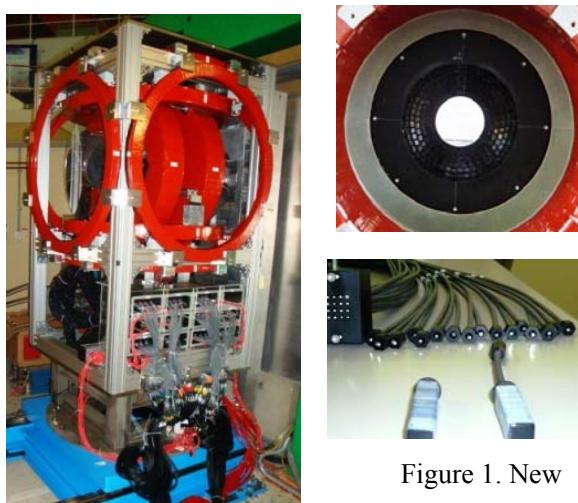


Figure 1. New spectrometer.

Figures 2. Counters, fiber light guides and counter mold.

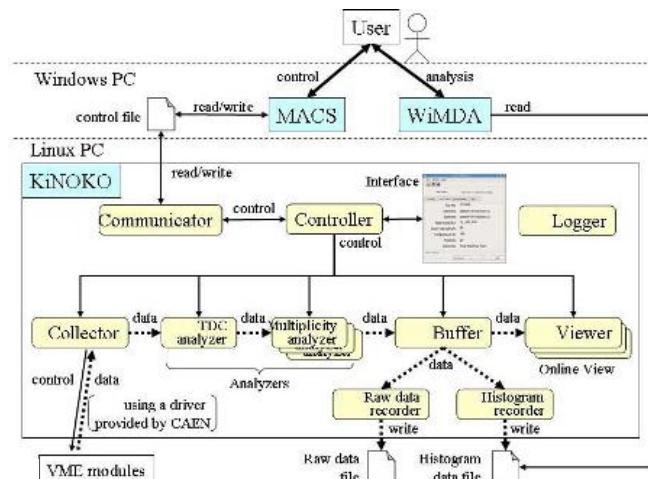


Figure 3 Block diagram of the data acquisition system.

### Publications

- 1) "Development of new  $\mu$ -e decay counters in a new multi-channel  $\mu$ SR spectrometer for an intense Pulsed muon beam", D. Tomono, et al., Nucl. Inst. Meth A **600**, 44 (2009).

## Study of Beam Density Enhancement Effect with Tapered Tubes for Application to the Muon Beam-Line at RIKEN-RAL

### Abstract

The beam density of 54 MeV/c muons can be increased almost by a factor of two when a tapered glass tube is inserted coaxially along with the muon beam. This opens a new technique to increase the muon intensity effectively.

### Achievement

More effective and efficient use of the muon beam is currently one of the important issues for  $\mu$ SR studies and high density muonium generation at RIKEN-RAL. We are planning to employ the capillary method to focus the intense muon beam. In the first experiment at RIKEN-RAL, we observed beam density enhancement effect at 54 MeV/c when a tapered glass tube was inserted coaxially along with the muon beam as shown Fig.1. The muon beam was scattered on their inner wall surface toward its outlets, resulting in increase of the number of available muons almost by a factor of two compared with using a normal beam collimator. The tapered angle dependence, outlet diameter dependence and particle charge dependence are shown in Fig. 2 and Fig. 3 comparing with the Monte Carlo simulation with a Coulomb scattering model. These results suggest that enhancement is explained by this model reasonably and that its material dependence is expected when tapered tubes are made of heavy materials rather than glass tube due to the large coefficient of the Coulomb scattering. This observation shows the possibility to increase the number of available muons further by optimizing the inner shape of collimators. For practical application, it is necessary to know the details of the energy and angular distribution of outgoing muons. A dedicated measurement was done at TRIUMF and analysis is in progress. These experiments were performed by the collaboration between atomic physics laboratory and advanced meson science laboratory.



Figure 1 One of the glass capillaries

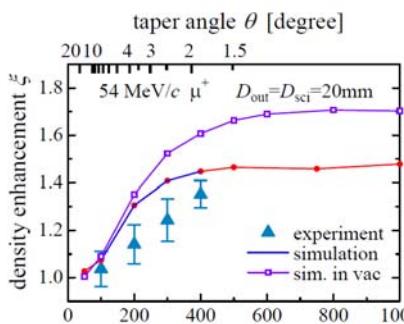


Figure 2 The beam density enhancement as a function of the tube length.

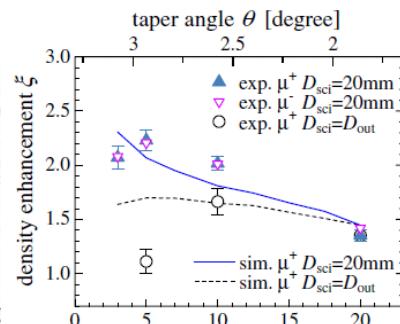


Figure 3 The beam density enhancement as a function of the outlet diameter  $D_{\text{out}}$ .

### Publications

as a function of the tube length. as a function of the outlet diameter  $D_{\text{out}}$ .

- 1) "Density Enhancement of Muon Beams with Tapered Glass Tubes", T M. Kojima, D. Tomono, T. Ikeda, K. Ishida, Y. Iwai, M. Iwasaki, Y. Matsuda, T. Matsuzaki, Y. Yamazaki, J. Phys. Soc. Jpn., **76**, 093501(2007).

## Development of an Ultra-Cold Muon Beam for the New Muon g-2 Measurement

### Abstract

Measurement of the muon's anomalous magnetic moment gives one of the most rigorous tests of the Standard Model of particle physics. We started a project to improve its precision using an ultra-cold muon beam. We first plan an increase of the muon intensity by two orders of magnitude with developments of muonium production targets, muonium ionizing laser and muon acceleration.

### Achievement

We have started, in collaboration with KEK and other laboratories, a project to measure the muon g-2 to 0.1 ppm precision. An ultra cold muon beam is a key for the success of the measurement. We had been developing an ultra slow muon beam as a tool for material science. To fulfill severer requirements for its application to muon g-2, several new developments are planned this year.

1) In order to produce an ultra-cold muon beam with small enough transverse momentum spread so it can be stored in the muon g-2 storage ring, we will utilize a muonium production target at room temperature instead of the 2100 K tungsten target we had been using. Silica-powders are known to produce similar yield of thermal muonium. We will study silica aerogels and nano-channeled targets as well. We will select the best target utilizing the muonium position measurement by positron tracking with the d.c. muon beam at TRIUMF.

2) We are collaborating with the RIKEN laser group to increase the Lyman- $\alpha$  laser intensity from 1  $\mu\text{J}/\text{pulse}$  to 100  $\mu\text{J}/\text{pulse}$ . Thus we will increase the ionization efficiency by two orders of magnitude. Manufacturing of the laser is in progress at RIKEN.

3) We are to adopt a new muon acceleration scheme based on the electron microscopes design so as not to heat the ultra-cold muon beam during acceleration. In addition, we plan a pulsed acceleration in the first stage so the field gradient does not increase the muon energy spread much.

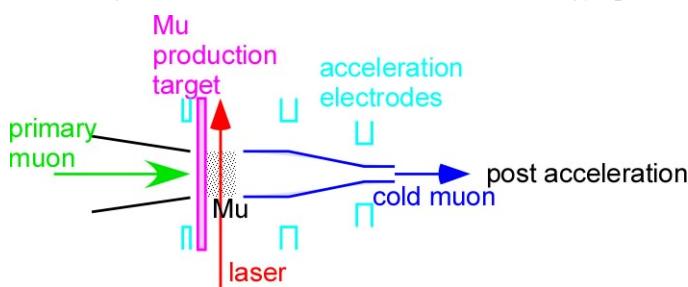


Figure: Schematic of the ultra-cold muon source

### Publications

- 1) "An Experimental Proposal on a New Measurement of the Muon Anomalous Magnetic Moment g-2 and Electric Dipole Moment at J-PARC", N. Saito (Contact person), M. Iwasaki (Co-contact person) et al., Proposal P34 submitted to the 9th J-PARC NPP PAC meeting, Tsukuba, Jan 2010.

## **Mössbauer spectroscopy**

The condensed matter physics using Mössbauer spectroscopy is rather new subject in the Advanced Meson Science Laboratory. Previously, Mössbauer spectroscopy was one of the subjects of Applied Nuclear Physics Laboratory initiated by a chief scientist Koichiro Asahi. In 2007, He resigned from RIKEN to concentrate Professorship in Tokyo Institute of Technology.

Because the Mössbauer spectroscopy is unique tool to study the property of condensed matter, giving complementary to  $\mu$ SR, so that it is one of the important subjects in Nishina Center. In addition, our group has developed on-line Mössbauer spectroscopy using a RI-beam produced in RIKEN-RIBF. It was the first application of the RI-beam to Mössbauer spectroscopy. It can be possible to observe the dynamic behaviors and the meta-stable states of located atoms in any condensed matter by in-beam Mössbauer spectroscopy. To maintain the activity, the Mössbauer group merged to the Advanced Meson Science Laboratory.

### 3.2.3.1 Individual topics of the Mössbauer spectroscopy

#### In-beam Mössbauer Spectroscopy of $^{57}\text{Mn}$ Implanted into $\text{Al}_2\text{O}_3$ and $\text{MgO}$

#### Abstract

Room-temperature ferromagnetism has been reported in some non-magnetic and semiconductive oxides, like  $\text{Al}_2\text{O}_3$ ,  $\text{MgO}$  and  $\text{ZnO}$ , including the trace magnetic transition-metals. It is suggested that the vacancies in non-magnetic oxides play an important role in the occurrence of ferromagnetism, but the origin of magnetism has not been understood yet. To this end, the  $^{57}\text{Mn}$  implantation Mössbauer spectroscopy was applied to clarify the occupation sites of the trace Fe atoms in  $\text{Al}_2\text{O}_3$  and  $\text{MgO}$ .

#### Achievement

The  $^{57}\text{Mn}$  beam was produced by the nuclear projectile fragmentation of  $^{58}\text{Fe}$  accelerated at 500 MeV/nucleon bombarding to a Be target, and was simultaneously implanted into a sample after passing through degraders. The 14.4 keV Mössbauer  $\gamma$ -rays emitted from  $^{57}\text{Mn}$  were detected by a parallel-plate avalanche counter (PPAC) combined with a plastic scintillation  $\beta$ -veto counter.

Figure 1 shows In-beam  $^{57}\text{Fe}$  Mössbauer spectrum of implanted  $^{57}\text{Mn}$  into  $\text{MgO}$  at room temperature. The spectrum could be analyzed with *Singlet-1* ( $\delta = -0.91 \text{ mm/s}$ ), *Doublet-1* ( $\delta = -0.57 \text{ mm/s}$ ,  $\Delta E_Q = 0.77 \text{ mm/s}$ ) and *Doublet-2* ( $\delta = -0.61 \text{ mm/s}$ ,  $\Delta E_Q = 2.15 \text{ mm/s}$ ), from the results of *ab initio* electron density calculations (see Fig.2). *Singlet-1* was assigned to high-spin  $\text{Fe}^{2+}$  at substitutional position of Mg atom. *Doublet-1* and *Doublet-2* are assigned to be Fe atoms with the neighboring O vacancies and those with the neighboring Mg vacancies, respectively.

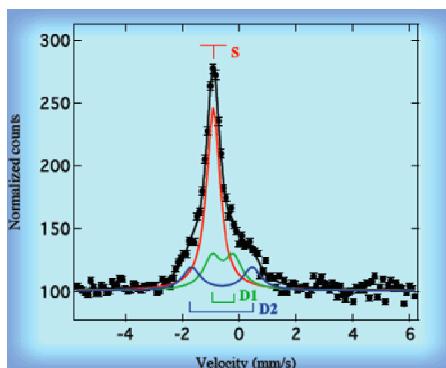


Figure 1: In-beam  $^{57}\text{Fe}$  Mössbauer spectrum of implanted  $^{57}\text{Mn}$  into  $\text{MgO}$  at R.T.

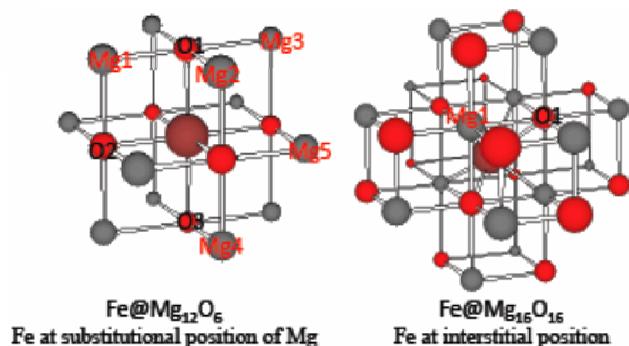


Figure 2: Cluster models of  $\text{MgO}(\text{Fe})$  for electron density calculations.

#### References

- 1) T. Nagatomo *et al.*, “In-beam Mössbauer Spectroscopy of  $^{57}\text{Mn}$  Implanted into Aluminum Oxide” and “In-beam Mössbauer Spectroscopy of  $^{57}\text{Mn}$  Implanted into Magnesium Oxide”, in *Asia-Pacific Symposium of Radiochemistry*, PO-2-124/125, Napa (USA), Nov. (2009).

## Development of Detector System for $^{57}\text{Mn}$ Implantation Mössbauer Studies

### Abstract

The detection system for  $^{57}\text{Mn}$  implantation Mössbauer studies was improved by using an anticoincidence method where a thin plastic scintillation counter was set between the detector and a sample in order to reject the  $\beta$  rays from  $^{57}\text{Mn}$ .  $^{57}\text{Mn}$  implantation Mössbauer spectrum with sufficient signal-to-noise (S/N) ratio that is about 20 times higher than that in previous measurements was successfully obtained.

### Achievement

The on-line Mössbauer spectroscopy using  $^{57}\text{Mn}$  as an radioisotope (RI) beam is considered to be one of the most powerful techniques to obtain *atomistic* information concerning the final lattice positions, valence states, and dynamic behavior of localized Fe atoms in materials. However, there still remained a significant problem that the  $\beta$  rays emitted from  $^{57}\text{Mn}$  nuclei much degraded a S/N ratio of the spectra, because the Mössbauer detector system has poor energy resolution to discriminate the conversion electrons emitted by Mössbauer effect from the incoming  $\beta$  rays.

The detector system was improved by using an anticoincidence method where a thin plastic scintillation counter was set between the Mössbauer  $\gamma$ -ray detector and a sample in order to reject the  $\beta$  rays from  $^{57}\text{Mn}$  (see Fig.3). The implantation Mössbauer spectrum with sufficient resonance that is about 20 times higher than that in previous measurements was successfully obtained (see Fig.4).

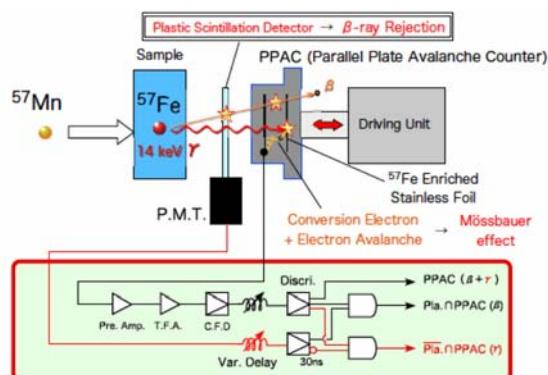


Figure 3: Schematic layout of the anticoincidence method for the  $^{57}\text{Mn}$  implantation Mössbauer measurement

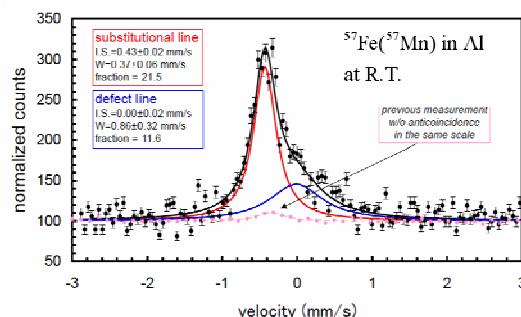


Figure 4: In-beam  $^{57}\text{Fe}$  Mössbauer spectrum of implanted  $^{57}\text{Mn}$  into Al at R.T. Pink shows the previous result without anticoincidence method.

### Reference

- 1) T. Nagatomo *et al.*, "Improvement of Signal-to-Noise Ratios in  $^{57}\text{Mn}$  Implantation Mössbauer Spectroscopy", in *Asia-Pacific Symposium of Radiochemistry*, 14-S12-6, Napa (USA), Nov. (2009).

## <sup>99</sup>Ru and <sup>61</sup>Ni Mössbauer Spectroscopic Studies Using RIKEN AVF Accelerator

### Abstract

Ru and Ni are important constituent elements of a wide variety of compounds, including oxides, catalysts, and functional materials in solar batteries and shape-memory alloys. <sup>99</sup>Ru and <sup>61</sup>Ni Mössbauer spectroscopy can reveal important information about the physical and chemical properties of these compounds. We have developed irradiation equipment for these Mössbauer sources at RIKEN AVF accelerator and an off-line measurement system at low temperatures.

### Achievement

The Mössbauer source nuclides, <sup>99</sup>Rh ( $T_{1/2}=15.0$  d) for <sup>99</sup>Ru Mössbauer spectroscopy and <sup>61</sup>Cu ( $T_{1/2}=3.4$  h) for <sup>61</sup>Ni Mössbauer studies, were produced via <sup>99</sup>Ru( $p,n$ )<sup>99</sup>Rh and <sup>58</sup>Ni( $\alpha,p$ )<sup>61</sup>Cu reactions at RIKEN AVF cyclotron in RNC. <sup>99</sup>Ru Mössbauer spectroscopy were applied to the studies of CaRuO<sub>3</sub>, Hg<sub>2</sub>Ru<sub>2</sub>O<sub>7</sub> and Ru-containing skutterudites, in addition of <sup>61</sup>Ni Mössbauer studies of NiO catalysts.

<sup>61</sup>Ni Mössbauer spectra have been obtained by a new technique using synchrotron radiation at SPring-8 and ESRF. Moreover, short-lived RI beam implantation methods will be available for studying of these elements. Conventional methods that use Mössbauer sources produced by nuclear reactions play a complementary role to these sophisticated Mössbauer techniques, but they remain the important Mössbauer measurement technique.

### Publications

- 1) Y. Kobayashi, *J. Phys.*, “<sup>99</sup>Ru and <sup>61</sup>Ni Mössbauer Spectroscopic Studies Using the Accelerator at RIKEN”, *in press*.
- 2) A. Koriyama M. Ishizaki, T. C. Ozawa, T. Taniguchi, Y. Nagata, H. Samata, Y. Kobayashi, and Y. Noro: “Magnetism of CaRuO<sub>3</sub> crystal”, *J. Alloys Compounds*, **372** (2004) 58-64.
- 3) S. Tsutsui, J. Umemura, H. Kobayashi, T. Tazaki, S. Nasu, Y. Kobayashi, Y. Yoda, H. Onodera, H. Sugawara, T. D. Matsuda, D. Kikuchi, H. Sato, C. Sekine and I. Shirotani: “Elastic Properties of Filled-Skutterudite Compounds Probed by Mössbauer Nuclei”, *Hyperfine Interact.*, **168** (2006) 1073-1077.
- 4) S. Tsutsui, Y. Kobayashi, T. Okada, H. Haba, H. Onodera, Y. Yoda, M. Mizumaki, H. Tanida, T. Uruga, C. Sekine, I. Shirotani, D. Kikuchi, H. Sugawara, and H. Sato: “A Possible Novel Magnetic Ordering in SmRu<sub>4</sub>P<sub>12</sub>”, *J. Phys. Soc. Jpn.*, **75** (2006) 0937031-4.

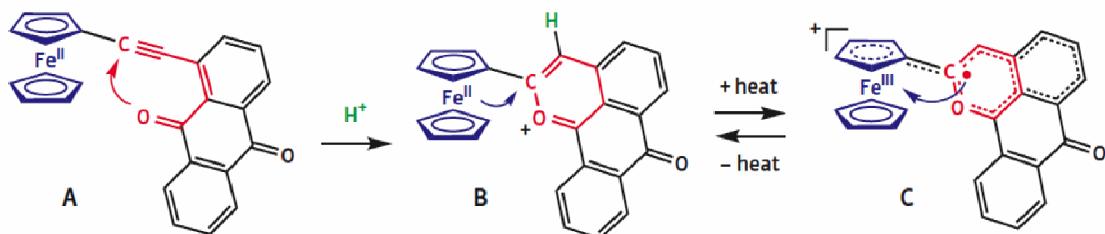
## Applications of $^{57}\text{Fe}$ Mössbauer Spectroscopy to Inorganic Chemistry

### Abstract

Chemical compounds that can be converted between two states with different physical properties are promising building blocks for molecular devices. If this process can be triggered with an external stimulus such as heat or light, it can be controlled remotely and integrated into a functional system. One method of introducing bistability into a material is to exploit valence tautomerization. We developed a new class of compounds exhibiting this type of behavior, and devised a novel chemical reaction to make cyclic structures that can switch between different valence isomers..

### Achievement

As the figure below indicates, when the molecule is treated with a strong organic acid ( $\text{H}^+$ ), a cyclization reaction occurs in which a bond forms between the carbon and oxygen atoms (Figure A). As a result of this process, a larger region of alternating single and double bonds, that is a  $\pi$ -conjugated system, is created. The Fe ion in the ferrocene group is able to transfer one of its electrons to the ring and is oxidized from Fe(II) to Fe(III) (Figure B and C). The donated electron can hop back to the ferrocene group to reform the compound B. The reorganization of the electronic structure depends on temperature; the structure in compound C is favored as it increases from 12 to 290 K.



### Publications

- 1) M. Kondo, M. Uchikawa, W. W. Zhang, K. Namiki, S. Kume, M. Murata, Y. Kobayashi, and H. Nishihara: "Protonation-Induced Cyclocondensation of 1-Aryl Ethynylanthraquinones: Expanding the  $\pi$  Conjugation", *Angewandte Chemie International Edition*, **46** (2007) 6271-6274.
- 2) M. Kondo, M. Uchikawa, K. Namiki, W. Zhang, S. Kume, E. Nishibori, H. Suwa, S. Aoyagi, M. Sakata, M. Murata, Y. Kobayashi, and H. Nishihara: "Counterion-Dependent Valence Tautomerization of Ferrocenyl-Conjugated Pyrylium Salts", *J. Am. Soc. Chem.*, **131** (2009) 12112-12124.

### 3.3. Record of Members

Current members (as of January 1, 2010).

|   | Number | Name  |
|---|--------|---|
| Research Staff<br>( permanent position )              | 9      | IWASAKI Masahiko<br>MATSUZAKI Teiichiro<br>ISHIDA Katsuhiko<br>KOBAYASHI Yoshio<br>OUTA Haruhiko<br>WATANABE Isao<br>ITAHASHI Kenta<br>OHNISHI Hiroaki<br>SAKUMA Fuminori |
| Special/Foreign<br>Postdoctoral Researchers           | 3      | OHISHI Kazuki<br>IIO Masami<br>RISDIANA   |
| Research Staff<br>( contract )                        | 8      | SUZUKI Takao<br>KOIKE Takahisa<br>KAWAMATA Takayuki<br>ISHII Yasuyuki<br>TSUKADA Kyo<br>MIZUNO Katsuya<br>TOMONO Dai<br>YOKOYAMA Koji                                     |
| Research Collaborative<br>Advisors                    | 4      | ITO Atsuko<br>AKAISHI Yoshinori<br>YAMAZAKI Toshimitsu<br>KAMIMURA Masayasu   |
| Visiting Researchers                                  | 2      | YAGI Eiichi<br>NAGATOMO Takashi   |
| Visiting<br>Researchers/Technicians<br>(Lab. Outside) | 128    |   |
| Junior Research Associates                            | 3      | ITO Satoshi<br>HIRAIWA Toshihiko<br>FUJIWARA Yuya   |
| Student Trainees                                      | 3      | TOKUDA Makoto<br>KOU Hiroshi<br>SATAKE Manami   |
| Student Trainees<br>(Lab. Outside)                    | 35     |   |
| Assistants  | 2      | SATO Junko<br>FUJITA Yoko   |

## History

### Research staff ( permanent position )

| Name                   | 2002 | 2003 | 2004 | 2005 | 2006 | 2007 | 2008 | 2009 | move to                               |
|------------------------|------|------|------|------|------|------|------|------|---------------------------------------|
| IWASAKI<br>Masahiko    |      |      |      |      |      |      |      | →    |                                       |
| MATSUZAKI<br>Teiichiro |      |      |      |      |      |      |      | →    |                                       |
| ISHIDA<br>Katsuhiko    |      |      |      |      |      |      |      | →    |                                       |
| KOBAYASHI<br>Yoshio    |      |      |      |      |      |      | →    |      |                                       |
| OUTA<br>Haruhiko       |      |      | →    |      |      |      |      | →    |                                       |
| WATANABE<br>Isao       |      |      |      |      |      |      |      | →    |                                       |
| ITAHASHI<br>Kenta      | →    |      |      |      |      |      |      | →    |                                       |
| OHNISHI<br>Hiroaki     |      |      |      | →    |      |      |      | →    |                                       |
| SAKUMA<br>Fuminori     |      |      |      |      |      |      | →    |      |                                       |
| MATSUDA<br>Yasuyuki    |      |      |      |      |      | →    |      |      | Associate Professor<br>Univ. of Tokyo |
| Sum                    | 6    | 7    | 7    | 8    | 8    | 9    | 8    | 9    |                                       |

### Special/Foreign Postdoctoral Researchers

| Name                | 2002 | 2003 | 2004 | 2005 | 2006 | 2007 | 2008 | 2009 | move to                               |
|---------------------|------|------|------|------|------|------|------|------|---------------------------------------|
| OHISHI<br>Kazuki    |      |      |      |      |      |      | →    |      |                                       |
| IIO<br>Masami       |      |      |      |      |      |      | →    |      |                                       |
| RISDIANA            |      |      |      |      |      |      | →    | →    |                                       |
| OKAMOTO<br>Satoshi  | →    |      |      |      |      |      |      |      | PD, Univ. of Columbia<br>(USA)        |
| YAMASE<br>Hiroyuki  | →    |      |      |      |      |      |      |      | Special PD Researcher<br>Takagi Lab.  |
| OHIRA<br>Seiko      | →    |      |      |      |      |      |      |      | Contract Researcher<br>Iwasaki Lab.   |
| KOIKE<br>Takahisa   |      | →    |      |      | →    |      |      |      | Contract Researcher<br>Iwasaki Lab.   |
| OKADA<br>Shinji     |      |      | →    |      |      |      |      |      | Contract Researcher<br>Iwasaki Lab.   |
| NEMURA<br>Hidekatsu |      |      |      |      | →    |      |      |      | Contract Researcher<br>Hiyama Lab.    |
| SUZUKI<br>Takatoshi |      |      |      |      | →    |      |      |      | Assistant Professor<br>Univ. of Tokyo |
| TOMONO<br>Dai       |      |      |      |      | →    |      | →    |      | Contract Researcher<br>Iwasaki Lab.   |
| SAKUMA<br>Fuminori  |      |      |      |      |      | →    | →    |      | Researcher<br>Iwasaki Lab.            |
| FUJIOKA<br>Hiroyuki |      |      |      |      |      |      | →    |      | Assistant Professor<br>Kyoto Univ.    |
| Sum                 | 3    | 1    | 2    | 5    | 4    | 4    | 6    | 3    |                                       |

**Contract Researchers**

| Name                 | 2002 | 2003 | 2004 | 2005 | 2006 | 2007 | 2008 | 2009 | move to                               |
|----------------------|------|------|------|------|------|------|------|------|---------------------------------------|
| SUZUKI<br>Takao      |      |      |      |      |      |      |      | →    |                                       |
| KOIKE<br>Takahisa    |      |      |      |      |      | →    |      |      |                                       |
| KAWAMATA<br>Takayuki |      |      |      |      | →    |      |      |      |                                       |
| ISHII<br>Yasuyuki    |      |      |      |      |      | →    | →    |      |                                       |
| TSUKADA<br>Kyo       |      |      |      |      |      | →    |      |      |                                       |
| MIZUNO<br>Katsuya    |      |      |      |      |      | →    | →    |      |                                       |
| TOMONO<br>Dai        |      |      |      | →    |      |      | →    | →    |                                       |
| YOKOYAMA<br>Koji     |      |      |      |      |      |      |      | →    |                                       |
| ITAHASHI<br>Kenta    | →    |      |      |      |      |      |      |      | Researcher<br>Iwasaki Lab.            |
| STRASSER<br>Patrick  | →    |      |      |      |      |      |      |      | Assistant Professor,<br>KEK           |
| OHIRA<br>Seiko       |      |      | →    |      |      |      |      |      | PD<br>Ochanomizu Univ.                |
| BAKULE<br>Pavel      |      |      | →    |      |      |      | →    |      | Researcher Scientist<br>ISIS, RAL(UK) |
| IIO<br>Masami        |      |      |      | →    |      |      | →    |      | Special PD Researcher<br>Iwasaki Lab. |
| OKADA<br>Shinji      |      |      |      |      | →    |      | →    |      | PD<br>LNF-INFN(Italy)                 |
| SATO<br>Masaharu     |      |      |      |      | →    |      | →    |      | PD<br>Univ. of Tokyo                  |
| Sum                  | 2    | 3    | 2    | 5    | 6    | 7    | 9    | 8    |                                       |

**Junior Research Associates**

| Name                 | 2002 | 2003 | 2004 | 2005 | 2006 | 2007 | 2008 | 2009 | move to                                 |
|----------------------|------|------|------|------|------|------|------|------|---|
| ITO<br>Satoshi       |      |      |      |      |      |      | →    |      |   |
| HIRAIWA<br>Toshihiko |      |      |      |      |      |      | →    |      |   |
| FUJIWARA<br>Yuya     |      |      |      |      |      |      | →    |      |   |
| TOMONO<br>Dai        | →    |      |      |      |      |      |      |      | Collaborated<br>Researcher, KEK         |
| TANAKA<br>Hiroyuki   | →    |      |      |      |      |      |      |      | PD<br>Univ. of California,<br>Riverside |
| SATO<br>Masaharu     |      |      |      |      | →    |      |      |      | Contract Researcher<br>Iwasaki Lab.     |
| HACHITANI<br>Kenichi |      |      |      |      |      | →    |      |      | Private cooperation                     |
| Sum                  | 2    | 1    | 1    | 2    | 2    | 1    | 2    | 3    |   |

### Student Trainees

| Name                | 2002 | 2003 | 2004 | 2005 | 2006 | 2007 | 2008 | 2009 | move to                               |
|---------------------|------|------|------|------|------|------|------|------|---------------------------------------|
| TOKUDA<br>Makoto    |      |      |      |      |      |      | →    |      |                                       |
| KOU<br>Hiroshi      |      |      |      |      |      |      | →    |      |                                       |
| SATAKE<br>Manami    |      |      |      |      |      |      | →    |      |                                       |
| OKADA<br>Shinji     | →    |      |      |      |      |      |      |      | Special PD Researcher<br>Iwasaki Lab. |
| YONEYAMA<br>Tetsu   | →    |      |      |      |      |      |      |      | Private corporation                   |
| SUZUKI<br>Takatoshi | →    |      |      |      |      |      |      |      | Special PD Researcher<br>Iwasaki Lab. |
| SATO<br>Masaharu    | →    |      |      |      |      |      |      |      | JRA<br>Iwasaki Lab.                   |
| KATAYAMA<br>Takeshi | →    |      |      |      |      |      |      |      | Private corporation                   |
| FUKUDA<br>Yoshiyuki |      |      |      | →    |      |      |      |      | Private corporation                   |
| HIRAYAMA<br>Yuzo    |      |      |      |      | →    |      |      |      | Private corporation                   |
| HANAKI<br>Toshio    |      |      |      |      | →    |      |      |      | Private corporation                   |
| TATSUNO<br>Hideyuki |      |      |      |      | →    |      |      |      | JSPS<br>DC2                           |
| FUJIWARA<br>Yuya    |      |      |      |      |      | →    |      |      | JRA<br>Iwasaki Lab.                   |
| Sum                 | 5    | 4    | 1    | 1    | 4    | 2    | 2    | 3    |                                       |

Half of the researchers of Advanced Meson Science Laboratory are concurrently appointed in RIKEN-RAL branch (RRB) to perform muon science researches and the facility operation. The budget of this activity is mainly supported by operational budget from RIKEN, executed by the chief scientist and inspected also by RRB director Teiichiro Matsuzaki. Another half of the researchers are focusing on the hadron physics. The budget of this activity is mainly supported by the external competitive fund from MEXT / JSPS, executed by the chief scientist.

### 3.4. Record of Fundings

| Fiscal year                          |  | 2002    | 2003    | 2004    | 2005    | 2006    | 2007    | 2008    | 2009    |
|--------------------------------------|--|---------|---------|---------|---------|---------|---------|---------|---------|
| Goverment<br>funding<br>(thru RIKEN) | Laboratory funding                       | 15,800  | 12,440  | 82,130  | 50,180  | 110,950 | 67,498  | 47,150  | 55,150  |
|                                      | RIKEN -RAL funding<br>(annual operating) | 280,522 | 252,135 | 205,224 | 205,224 | 205,224 | 205,224 | 194,977 | 175,479 |
| Research grants                      |  | 4,650   | 18,800  | 31,600  | 31,000  | 75,840  | 108,660 | 97,610  | 23,500  |
| Commissioned<br>research funding     |  | 420     | 0       | 0       | 0       | 0       | 0       | 0       | 300     |
| Others                               |  | 0       | 0       | 0       | 0       | 0       | 4,545   | 0       | 0       |
|                                      |  | 303,394 | 285,378 | 320,958 | 288,409 | 394,020 | 387,934 | 341,745 | 256,438 |

## 3.5. Research Outputs

| Year              |               | 2002 | 2003 | 2004 | 2005 | 2006 | 2007 | 2008 | 2009 |
|-------------------|---------------|------|------|------|------|------|------|------|------|
| Publication       | International | 26   | 37   | 40   | 40   | 56   | 47   | 55   | 64   |
|                   | Domestic      | 0    | 0    | 1    | 3    | 0    | 0    | 2    | 0    |
| Oral presentation | International | 12   | 13   | 18   | 47   | 32   | 35   | 35   | 26   |
|                   | Domestic      | 9    | 12   | 14   | 23   | 37   | 39   | 31   | 24   |
| Patent            | International | 0    | 0    | 0    | 0    | 0    | 0    | 0    | 0    |
| Application       | Domestic      | 0    | 0    | 0    | 0    | 0    | 0    | 0    | 0    |

### 3.5.1. Original papers

#### 2010

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**“DD 系ミュオン触媒核融合の凝縮系オルソパラ効果”**,  
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Awards given to the work done in Advanced Meson Science Laboratory

- T. Yamazaki(Research Consultant) received  
“The Person of Cultural Merit Award (Agency for Cultural Affairs)” in 2009
- S. Okada (former Special Postdoctoral Researcher) received  
“3<sup>rd</sup> ” Young Scientist Award of the Physical Society of Japan” in 2009
- I. Watanabe (Senior Researcher) received  
“44<sup>th</sup> Toray Science and Technology Grant” in 2003
- K. Itahashi (Senior Researcher) received  
“The GSI Exotic Nuclei Community Membership Award” in 2003

### 3.6. Future Plan

Feasibility studies of the future projects themselves are part of our research activity, and some of them are described in research topics of this paper. Actually, many experimental proposals / letter of intents have been submitted to J-PARC PAC from RNC researchers including ourselves. Hadron physics side, it is very difficult to realize those experimental projects in short term with the present hadron hall at J-PARC. Thus, we are proposing “institutional-based contribution to J-PARC” to internal RNC conceptual-planning committee, focusing on research related with “origin of the hadron (matter) mass” together with Radiation Laboratory of RNC. As a part of discussions of the future project of RNC, our proposal was discussed as one of the interesting scenarios, although it is not very simple because RNC has primary mission to operate RIBF.

In the muon science side, we are facing at extreme difficulty for the future plan. because quite similar facility MUSE is starting up in J-PARC MLF hall, although there is still a long way to go for the full commissioning. It is quite important to make our facility as competitive as possible, and to establish work sharing with KEK muon group. One of the key issues is to make substantial progress in ultra-cold muon production and establish the technology to accelerate the ultra-cold muon without heating up. If it is realized, one can realize “muon magnetic microscope” which can open new research field in muon science. This is also a key to realize “new precise muon anomalous-magnetic-moment ( $g-2$ ) measurement” to confirm the physics beyond standard model. The high temperature solid  $\mu$ CF is also an important research program to be done at RIKEN-RAL, because of the difficulty of radiation safety at J-PARC.

The Advanced Meson Science Laboratory’s future plan has a deep relationship with J-PARC. We wish to try our best to realize “RIKEN J-PARC cooperation center” under the leadership of the RNC director.

### 3.7. Curriculum Vitae (Chief Scientist)

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#### Education and Degrees

|            |   |
|------------|---|
| March 1982 | Graduate from Department of Physics,<br>Faculty of Science, University of Tokyo         |
| March 1987 | Completion of Department of Physics,<br>Graduate School of Science, University of Tokyo |
| June 1987  | Doctor of Science, University of Tokyo  |

#### Appointment:

|                      |  |
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| April 1987           | Assistant Professor, Meson Science Laboratory,<br>Faculty of Science, University of Tokyo  |
| January 1997         | Associate Professor, Department of Physics,<br>Graduate School of Science and Engineering,<br>Tokyo Institute of Technology  |
| April 2002           | Chief Scientist, Muon Science Laboratory,<br>Discovery Research Institute,<br>RIKEN (Institute for Physical and Chemical Research)<br>Visiting Professor, Department of Physics,<br>Graduate School of Science and Engineering,<br>Tokyo Institute of Technology |
| April 2004           | Chief Scientist, Advanced Meson Science Laboratory,<br>Discovery Research Institute,<br>RIKEN (Institute for Physical and Chemical Research)   |
| April 2005 – Present | Professor, Department of Physics,<br>Graduate School of Science and Engineering,<br>Tokyo Institute of Technology  |
| April 2006 - Present | Chief Scientist, Advanced Meson Science Laboratory,<br>RIKEN Nishina Center for Accelerator-Based Science,<br>RIKEN (Institute for Physical and Chemical Research)   |

#### Prize:

|            |   |
|------------|---|
| March 2005 | 22 <sup>nd</sup> Inoue Prize for Science /<br>Inoue Foundation for Promotion of Science |
|------------|---|

**Academic Activities:**

- (present)      Chair of Chief Scientist Assembly, RIKEN  
                  Member of Japan Physics Society  
                  Member of Japan Meson/Muon Science Society  
                  Member of Program Advisory Committee of Muon Science of  
                  J-PARC