

Light kaonic nuclei at J-PARC



F. Sakuma, RIKEN



on behalf of the J-PARC E15/E80
collaborations

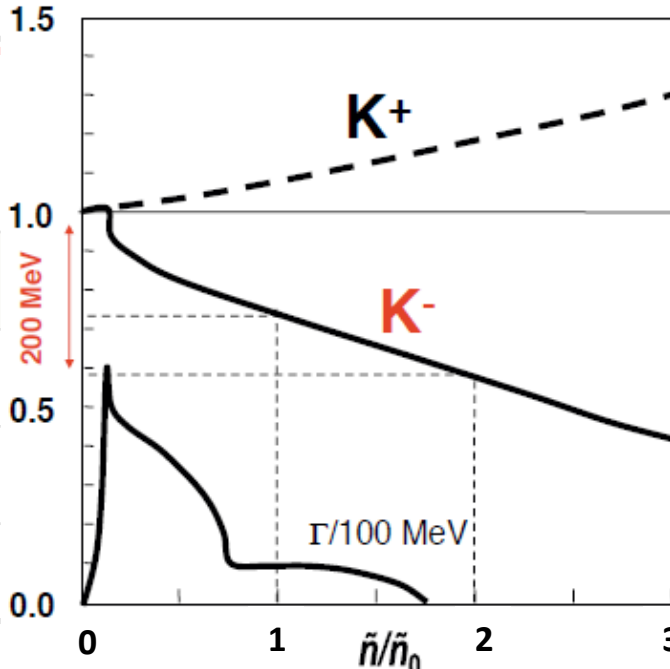
Strangeness Nuclear Physics Workshop, 2021
18-19 December 2021, Online

Kaonic Nuclei

Predicted from

Kaonic Nuclei	Binding Energy [MeV]
$\Lambda(1405) = K^-p$	27
K^-pp	48
K^-ppp	97
K^-ppn	118

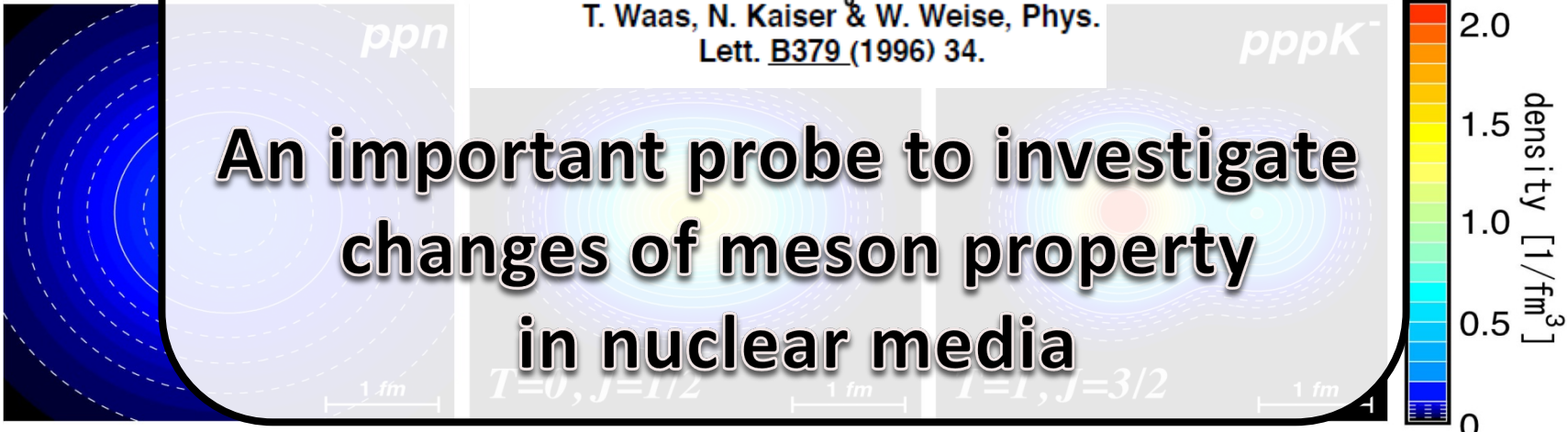
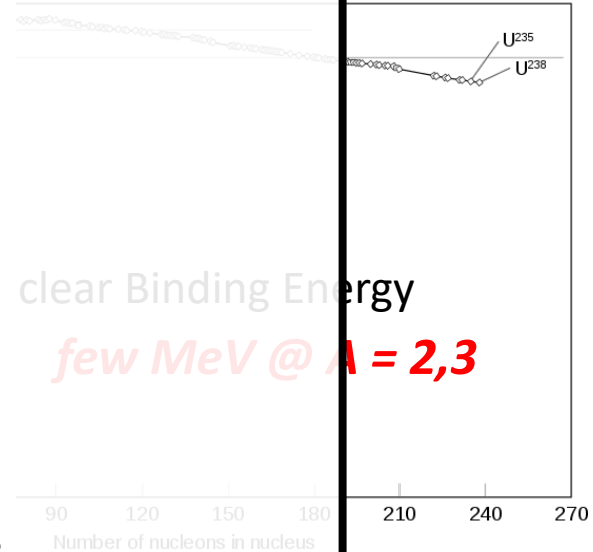
m_K^*/m_K in nuclear matter



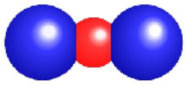
T. Waas, N. Kaiser & W. Weise, Phys. Lett. **B379** (1996) 34.

action in $I=0$

88., PRC65(2002)044005., etc.



An important probe to investigate changes of meson property in nuclear media



K⁻pp

Theoretical Calculations on “K⁻pp”

The simplest kaonic nuclei, $\bar{K}NN$ ($J^P=0^-, I=1/2$)

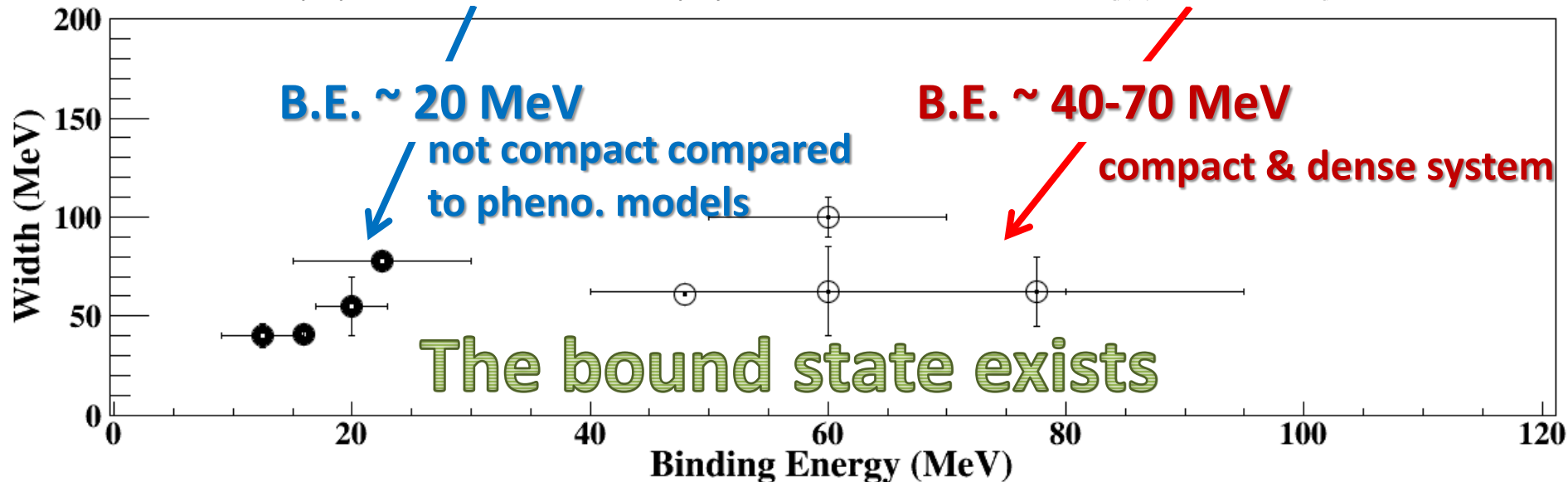
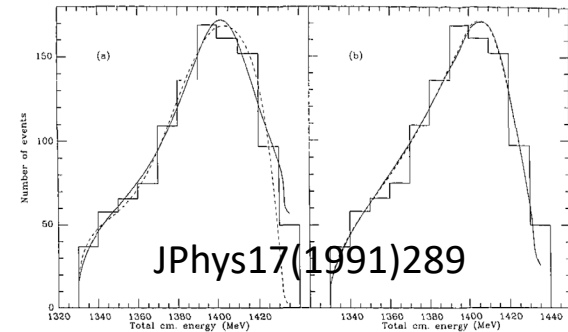
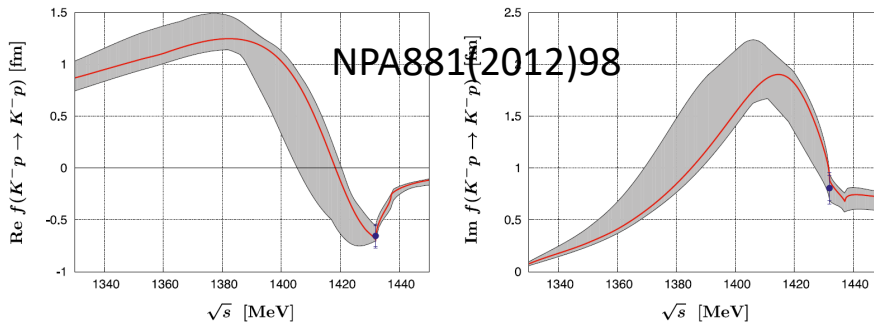
$\bar{K}N$ int.

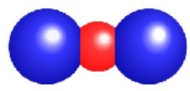
Chiral SU(3)
(energy dependent)

Phenomenological
(energy independent)

$M_{\Lambda(1405)} \sim 1420$, double pole

$M_{\Lambda(1405)} \sim 1405$, single pole







Experimental Situation "before E15"

Negative results

AMADEUS@DAΦNE 
 $^{12}\text{C}(K^-_{\text{stopped}}, \Lambda p)$ EPJ**C79**(2019)190

HADES@GSI 
 $p + p \rightarrow (\Lambda + p) + K^+$ @ 3.5 GeV
PLB**742**(2015)242

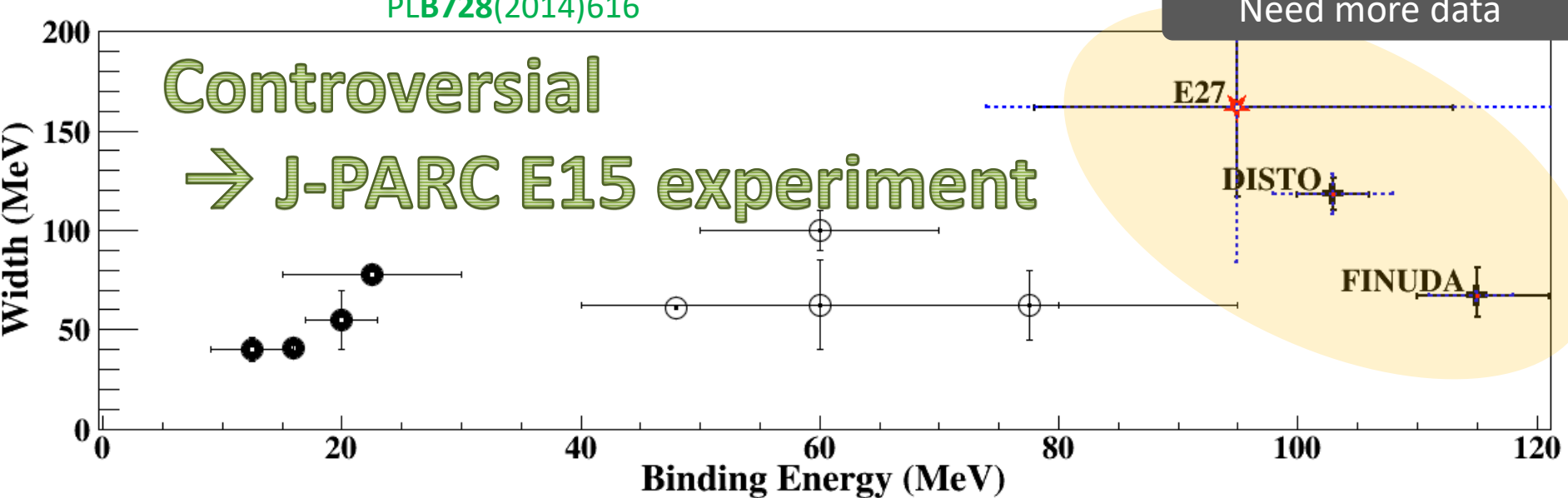
LEPS@SPring-8 
 $d(\gamma, \pi^- K^+)X$ @ 1.5-2.4 GeV
PLB**728**(2014)616

Positive results

FINUDA@DAΦNE PRL**94**(2005)212303
 $^6\text{Li}/^7\text{Li}/^{12}\text{C}(K^-_{\text{stopped}}, \Lambda p)$ 

DISTO@SATURNE PRL**104**(2010)132502
 $p + p \rightarrow (\Lambda + p) + K^+$ @ 2.85 GeV

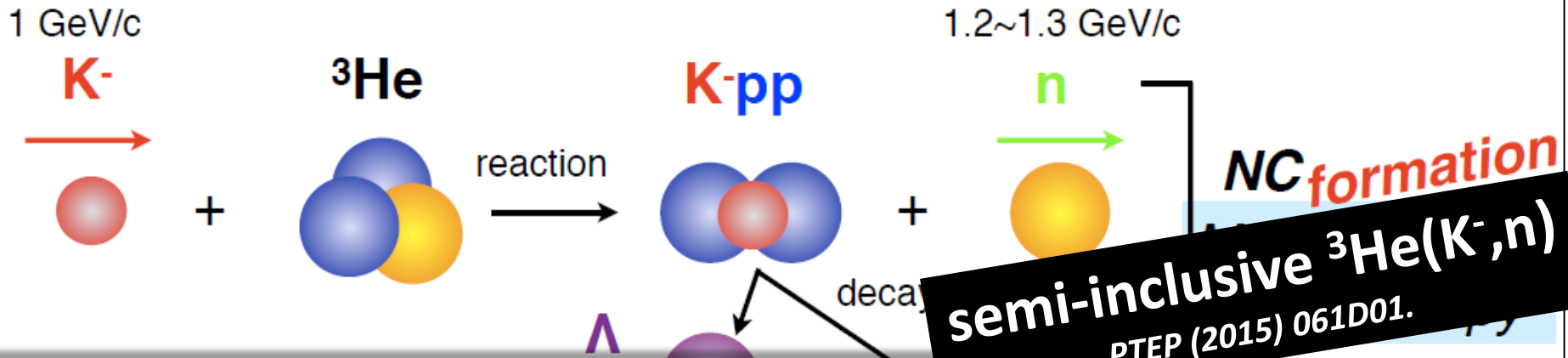

E27@J-PARC PTEP(2015)021D01.
 $d(\pi^+, K^+)\Sigma^0 p$ @ 1.69 GeV/c



J-PARC E15 Experiment

- ${}^3\text{He}(\text{in-flight } K^-, n)$ reaction @ 1.0 GeV/c

😊 **2NA and Λ decays can be discriminated kinematically**



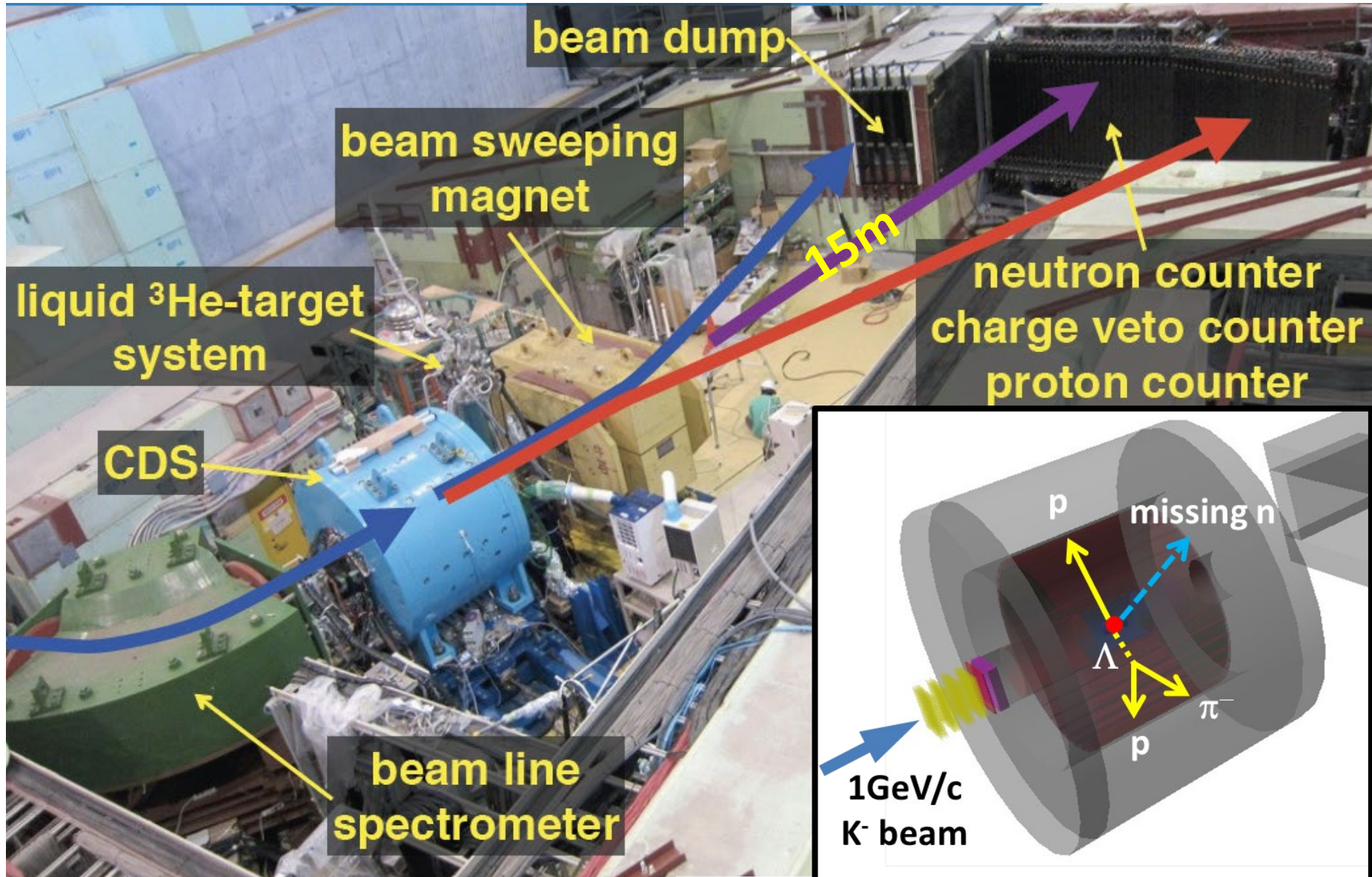
PHYSICAL REVIEW C **102**, 044002 (2020)

Observation of a $\bar{K}NN$ bound state in the ${}^3\text{He}(K^-, \Lambda p)n$ reaction

T. Yamaga,^{1,*} S. Ajimura,² H. Asano,¹ G. Beer,³ H. Bhang,⁴ M. Bragadireanu,⁵ P. Buehler,⁶ L. Busso,^{7,8} M. Cargnelli,⁶ S. Choi,⁴ C. Curceanu,⁹ S. Enomoto,¹⁴ H. Fujioka,¹⁵ Y. Fujiwara,¹² T. Fukuda,¹³ C. Guaraldo,⁹ T. Hashimoto,²⁰ R. S. Hayano,¹² T. Hiraiwa,² M. Iio,¹⁴ M. Iliescu,⁹ K. Inoue,² Y. Ishiguro,¹¹ T. Ishikawa,¹² S. Ishimoto,¹⁴ K. Itahashi,¹ M. Iwai,¹⁴ M. Iwasaki,^{1,†} K. Kanno,¹² K. Kato,¹¹ Y. Kato,¹ S. Kawasaki,¹⁰ P. Kienle,^{16,‡} H. Kou,¹⁵ Y. Ma,¹ J. M. Y. Matsuda,¹⁷ Y. Mizoi,¹³ O. Morra,⁷ T. Nagae,¹¹ H. Noumi,^{2,14} H. Ohnishi,²² S. Okada,²³ H. Ota,¹ K. Piscicelli,¹⁸ Y. Sada,²² A. Sakaguchi,¹⁰ F. Sakuma,¹ M. Sato,¹⁴ A. Scordo,⁹ M. Sekimoto,¹⁴ H. Shi,⁶ K. Shirotori,² D. Sirghi,^{9,5} F. S. Suzuki,¹⁴ T. Suzuki,¹² K. Tanida,²⁰ H. Tatsuno,²¹ M. Tokuda,¹⁵ D. Tomono,² A. Toyoda,¹⁴ K. Tsukada,¹⁸ O. Vazquez Doce,^{9,16} E. Widmann,⁶ T. Yamazaki,^{12,1} H. Yim,¹⁹ Q. Zhang,¹ and J. Zmeskal⁶
 (J-PARC E15 Collaboration)

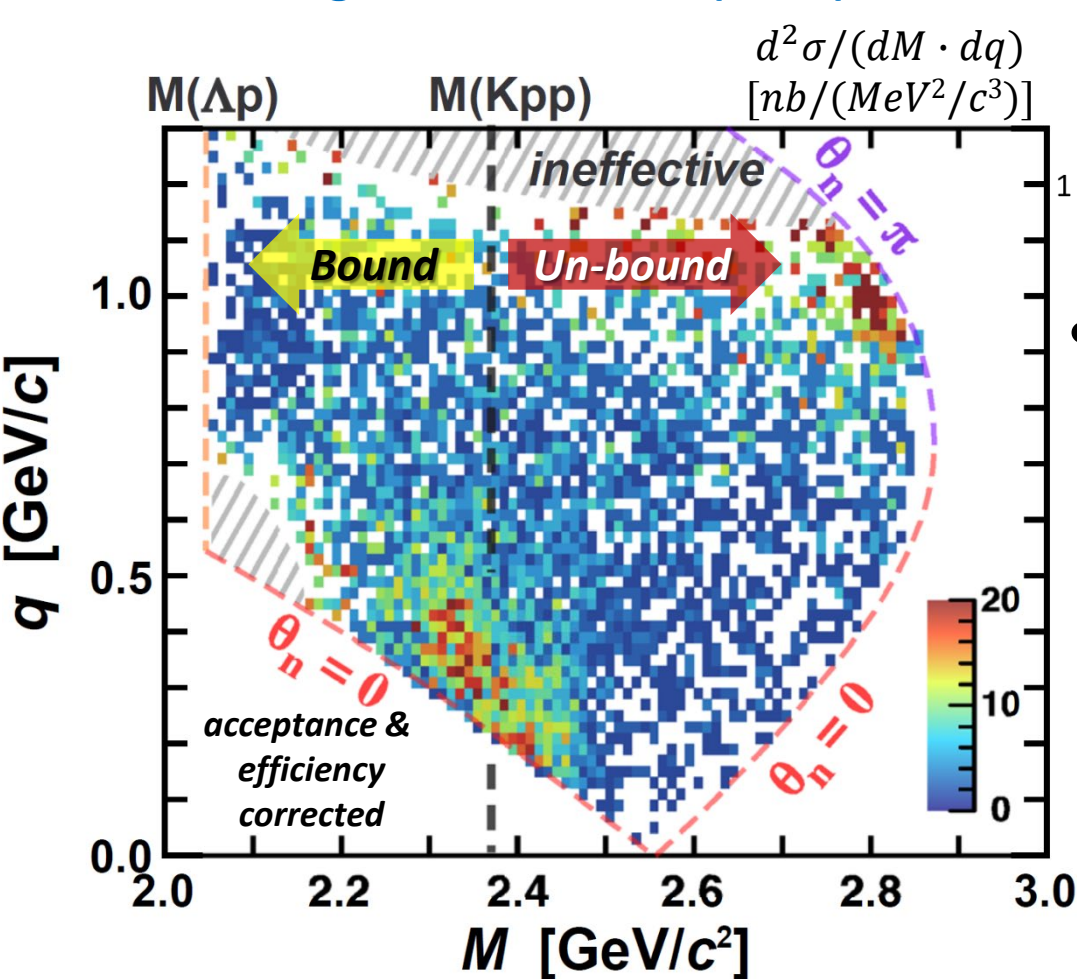
exclusive $(K^-, \Lambda p)n$ missing
 PTEP (2016) 051D01.,
 PLB789(2019)620., PRC102(2020)044002.

Experimental Setup @ K1.8BR



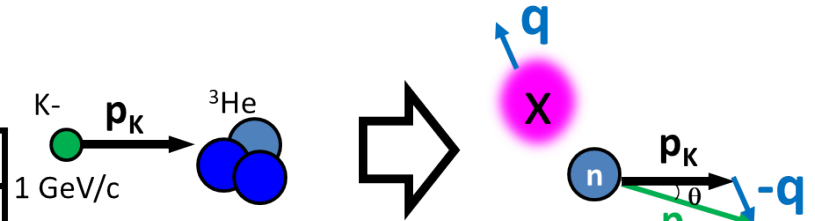
Exclusive ${}^3\text{He}(\text{K}^-, \Lambda p)n$

T.Yamaga et. al., PRC102(2020)044002.



q : momentum transfer of (K⁻,n)

M : invariant mass of Λp



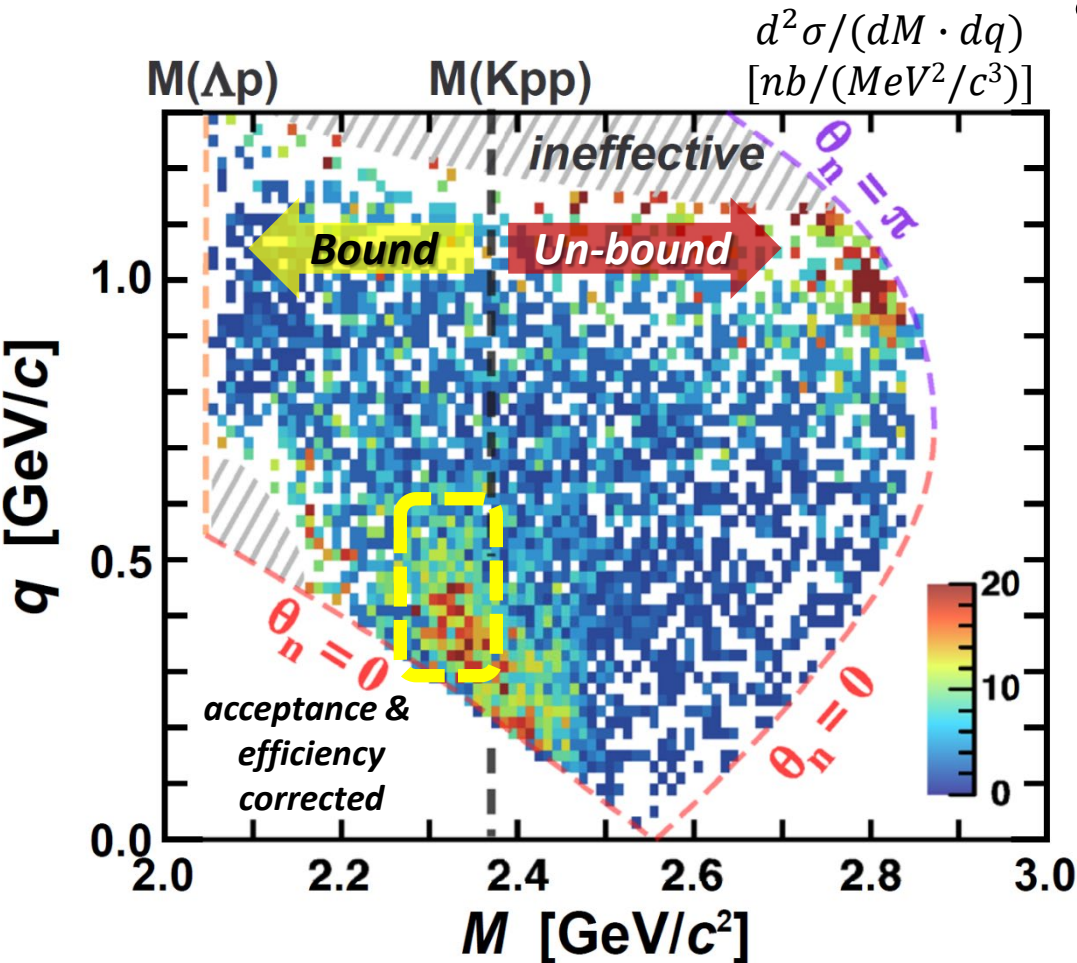
- Momentum transfer analysis using the (K⁻,n) reaction

✓ $M(\Lambda p)$ vs. q

✓ give a clear information on reaction processes

Exclusive ${}^3\text{He}(K^-, \Lambda p)n$

T.Yamaga et. al., PRC102(2020)044002.

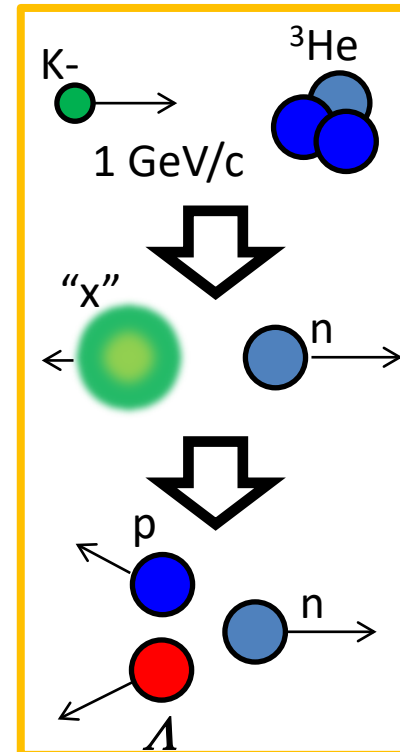


q : momentum transfer of (K^-, n)

M : invariant mass of Λp

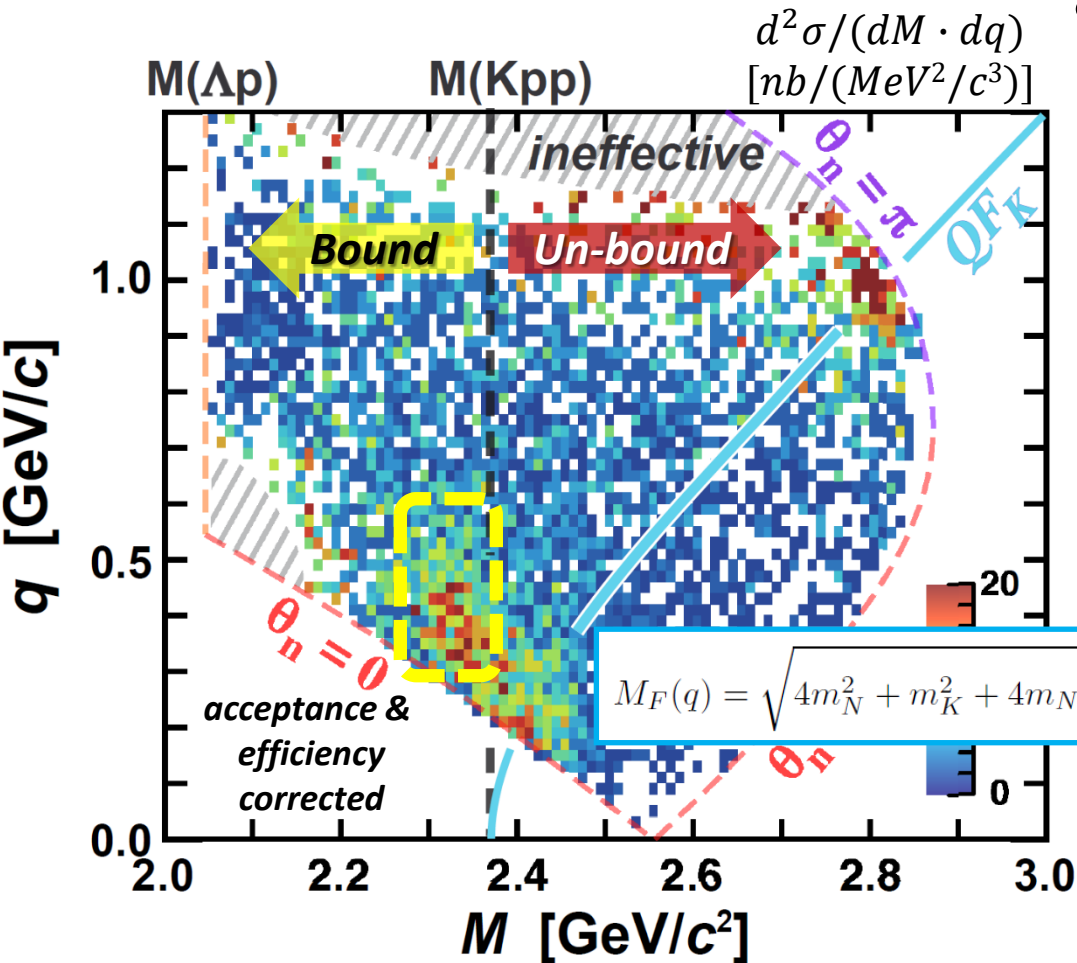
- “K⁻pp” bound state exists

✓ q -independent



Exclusive ${}^3\text{He}(K^-, \Lambda p)n$

T.Yamaga et. al., PRC102(2020)044002.



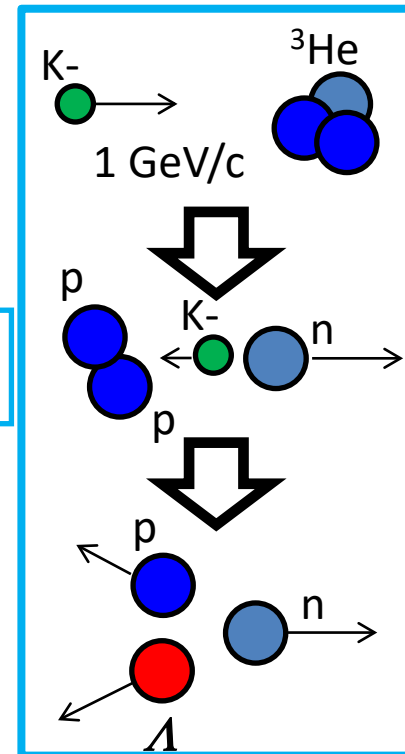
q : momentum transfer of (K^-, n)

M : invariant mass of Λp

- **QF followed by 2NA**

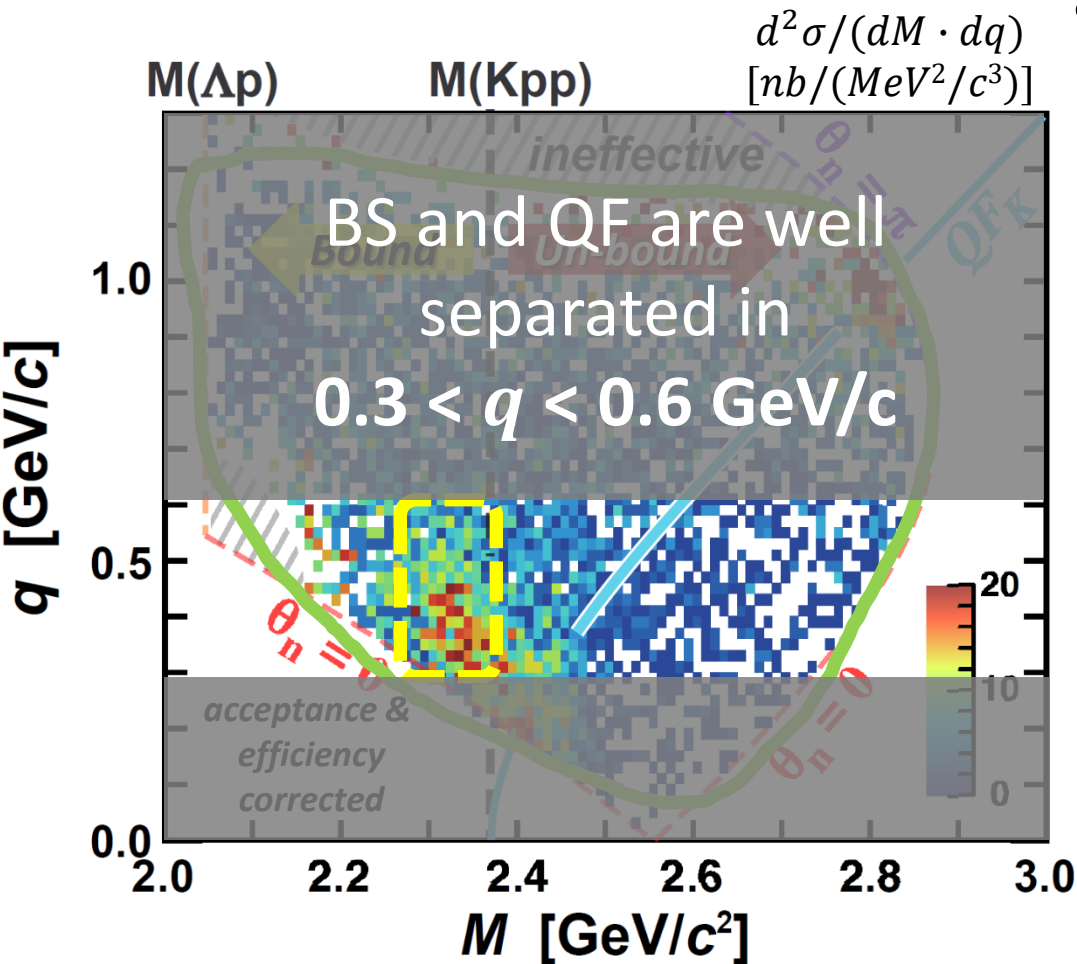
- ✓ via “on-shell K^- ”

- ✓ q -dependent



Exclusive ${}^3\text{He}(K^-, \Lambda p)n$

T.Yamaga et. al., PRC102(2020)044002.



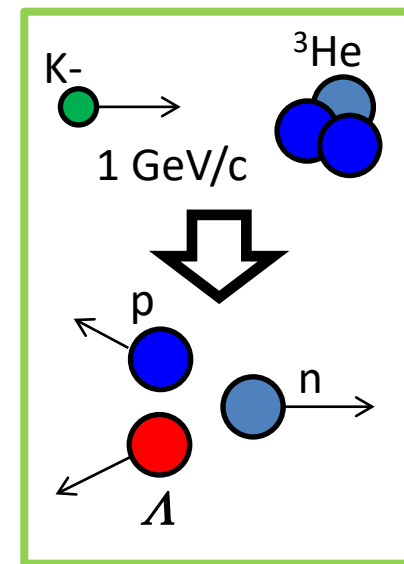
q : momentum transfer of (K^-, n)

M : invariant mass of Λp

- **Broad Component**

- ✓ 3NA reaction?

- ✓ Further investigations are ongoing

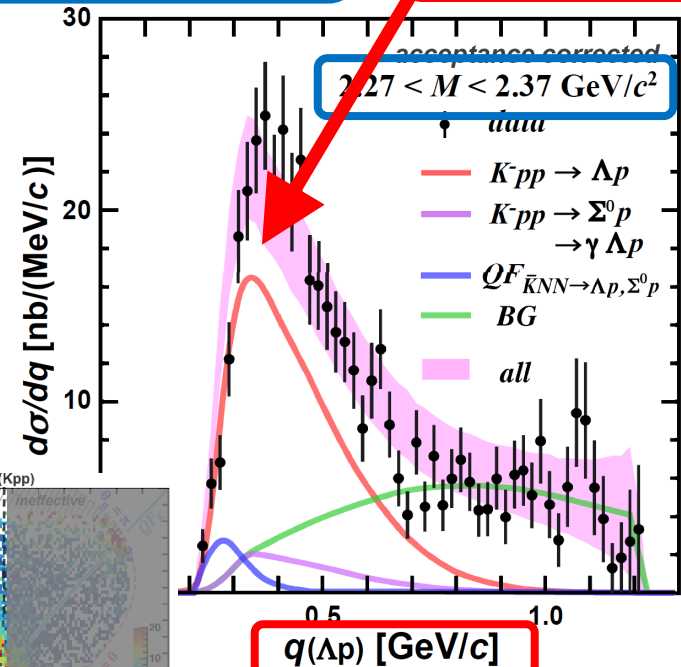
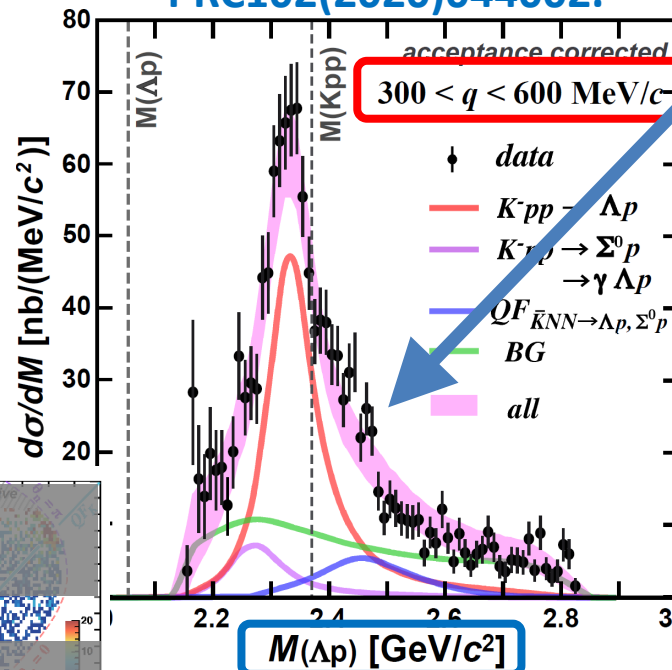


“K-pp” Bound State

Fit with PWIA

$$\sigma(M, q) \propto \rho(M, q) \times \frac{(\Gamma_{Kpp}/2)^2}{(M - M_{Kpp})^2 + (\Gamma_{Kpp}/2)^2} \times \exp\left(-\frac{q^2}{Q_{Kpp}^2}\right)$$

PRC102(2020)044002.



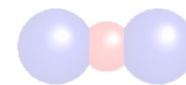
$B_{Kpp} \sim 40$ MeV, $\Gamma_{Kpp} \sim 100$ MeV
 \rightarrow large binding energy

$Q_{kpp} \sim 400$ MeV (c.f. $Q_{QF} \sim 200$ MeV)
 \rightarrow wide momentum transfer

\rightarrow suggest the “K-pp” is quite compact

Many Questions to be Answered

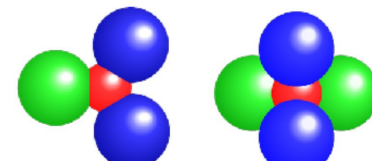
- Further details of the $\bar{K}NN$
 - Spin and parity of the “ K^-pp ”?
 - Really compact and dense system?
- $\Lambda(1405)$ state
 - $\bar{K}N$ quasi-bound state as considered?
 - Size?
 - Relation between $\bar{K}N$ and $\bar{K}NN$?
- **More heavier kaonic nuclei?**
 - Mass number dependence?
- **Double kaonic nuclei?**
 - Much compact and dense system?



K^-pp

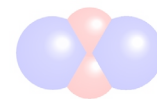


K^-p

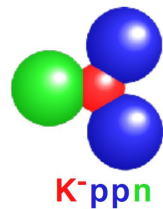


K^-ppn

K^-ppnn



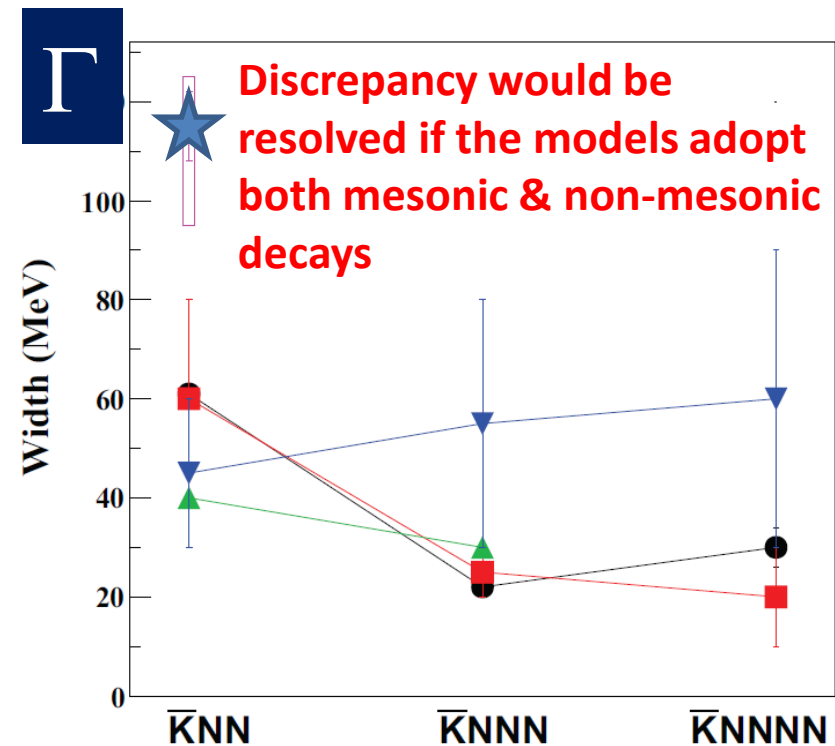
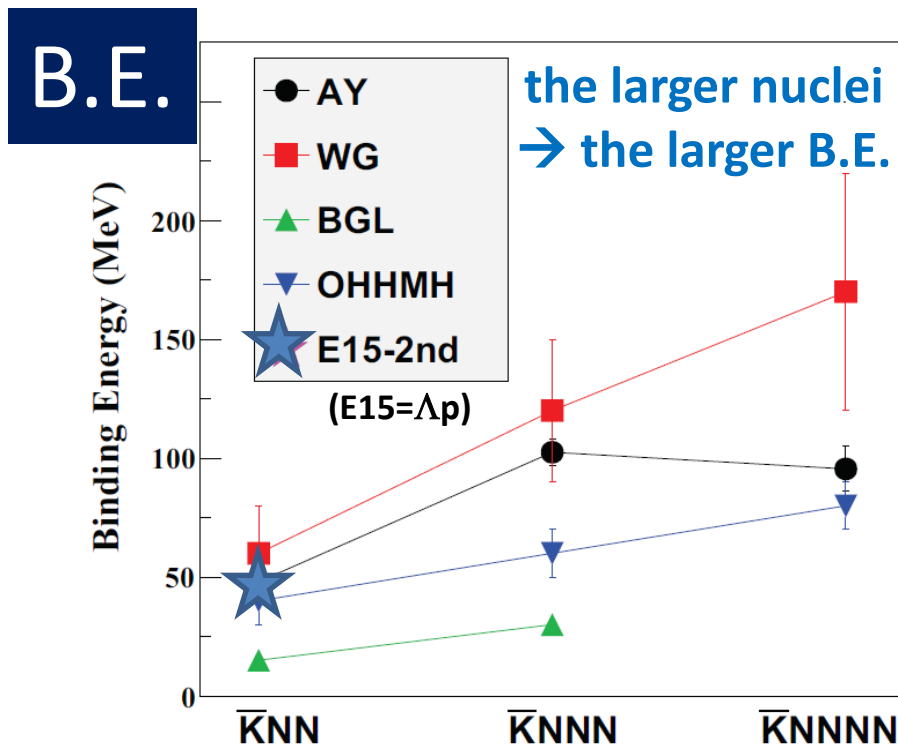
K^-K^-pp



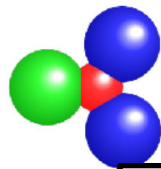
$\bar{K}NNN$ and $\bar{K}NNNN$



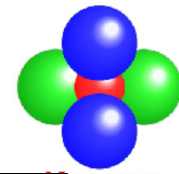
- More heavier system must be explored to provide more conclusive evidence of the kaonic nuclei



- The $\bar{K}NNNN$ system is expected to be the most compact system due to an α particle configuration



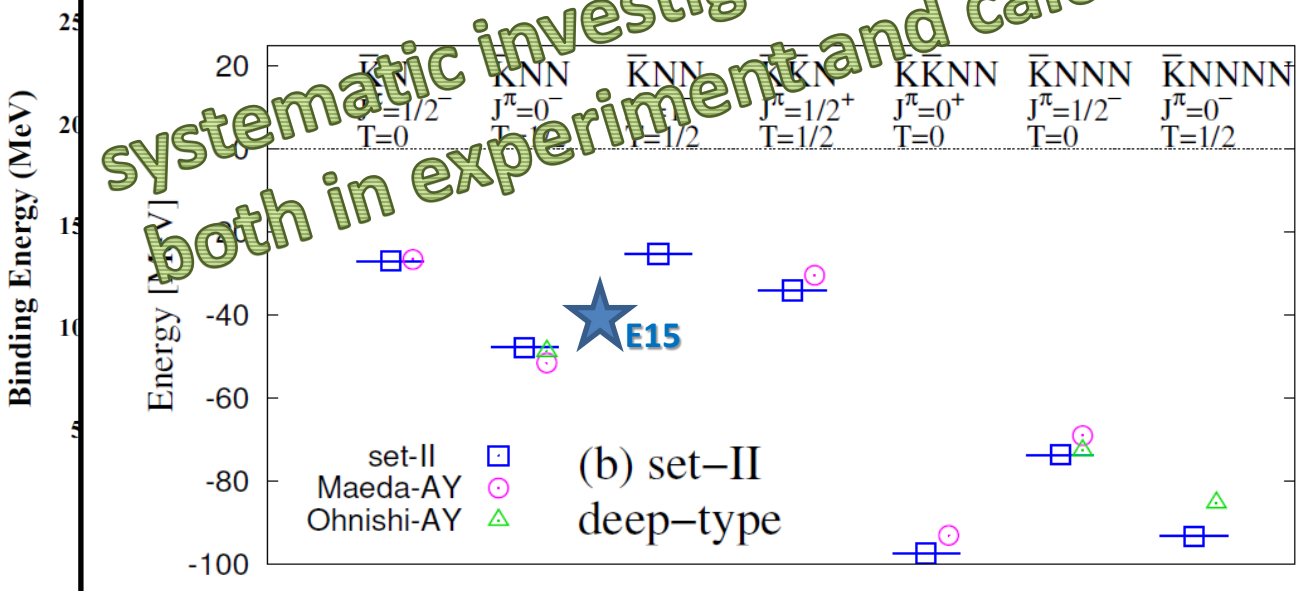
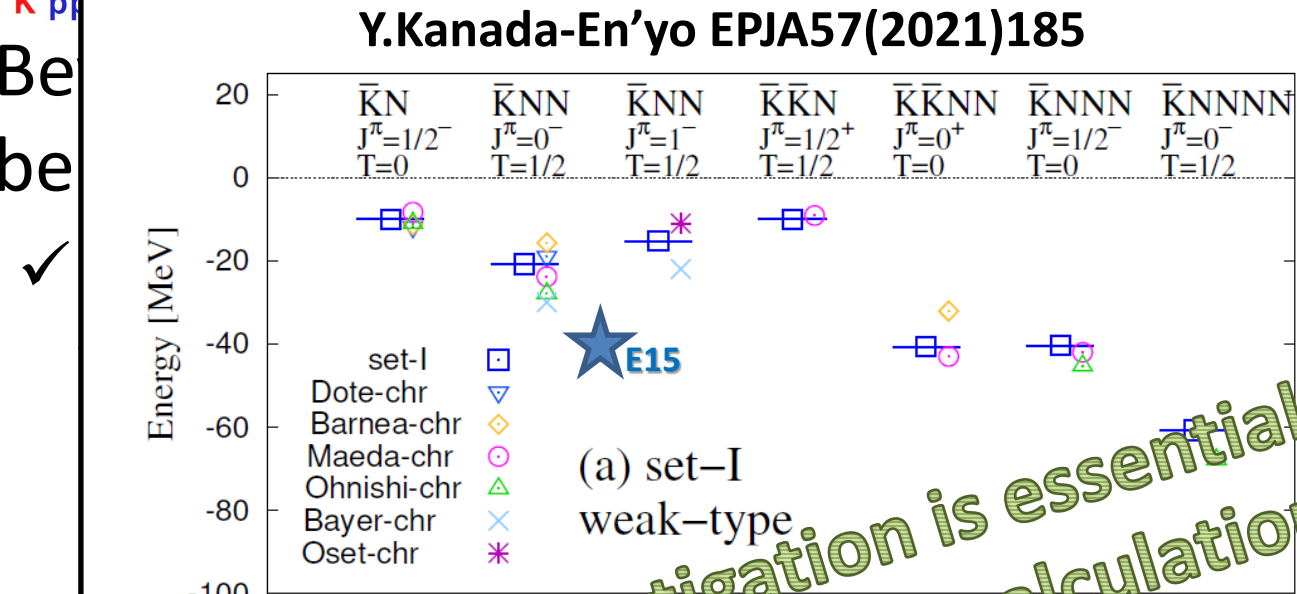
$\bar{K}NNN$ and $\bar{K}NNNN$



ppnn

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systematic investigation is essential
both in experiment and calculation



NNNN

New Project @ J-PARC

-- systematic investigation of the light kaonic nuclei --

Strategy of the New Project

- for systematic study from the $\bar{K}N$ to $\bar{K}NNNN$ systems -

	Reaction	Decays	Key
$\bar{K}N$	$d(K^-,n)$	$\pi^{\pm 0}\Sigma^{\mp 0}$	F-factor \rightarrow n/ γ identification
$\bar{K}NN$	${}^3\text{He}(K^-,N)$	$\Lambda p/\Lambda n$	$J^P \rightarrow$ polarimeter
$\bar{K}NNN$	${}^4\text{He}(K^-,N)$	$\Lambda d/\Lambda pn$	large acceptance \leftarrow A first step
$\bar{K}NNNN$	${}^6\text{Li}(K^-,d)$	$\Lambda t/\Lambda dn/\Lambda pnn$	<i>many body decay</i>
$\bar{K}\bar{K}NN$	$\bar{p} + {}^3\text{He}$	$\Lambda\Lambda$	<i>\bar{p} beam yield</i>

- To realize the systematic measurements, we need

- \square a large acceptance spectrometer **\leftarrow new CDS****

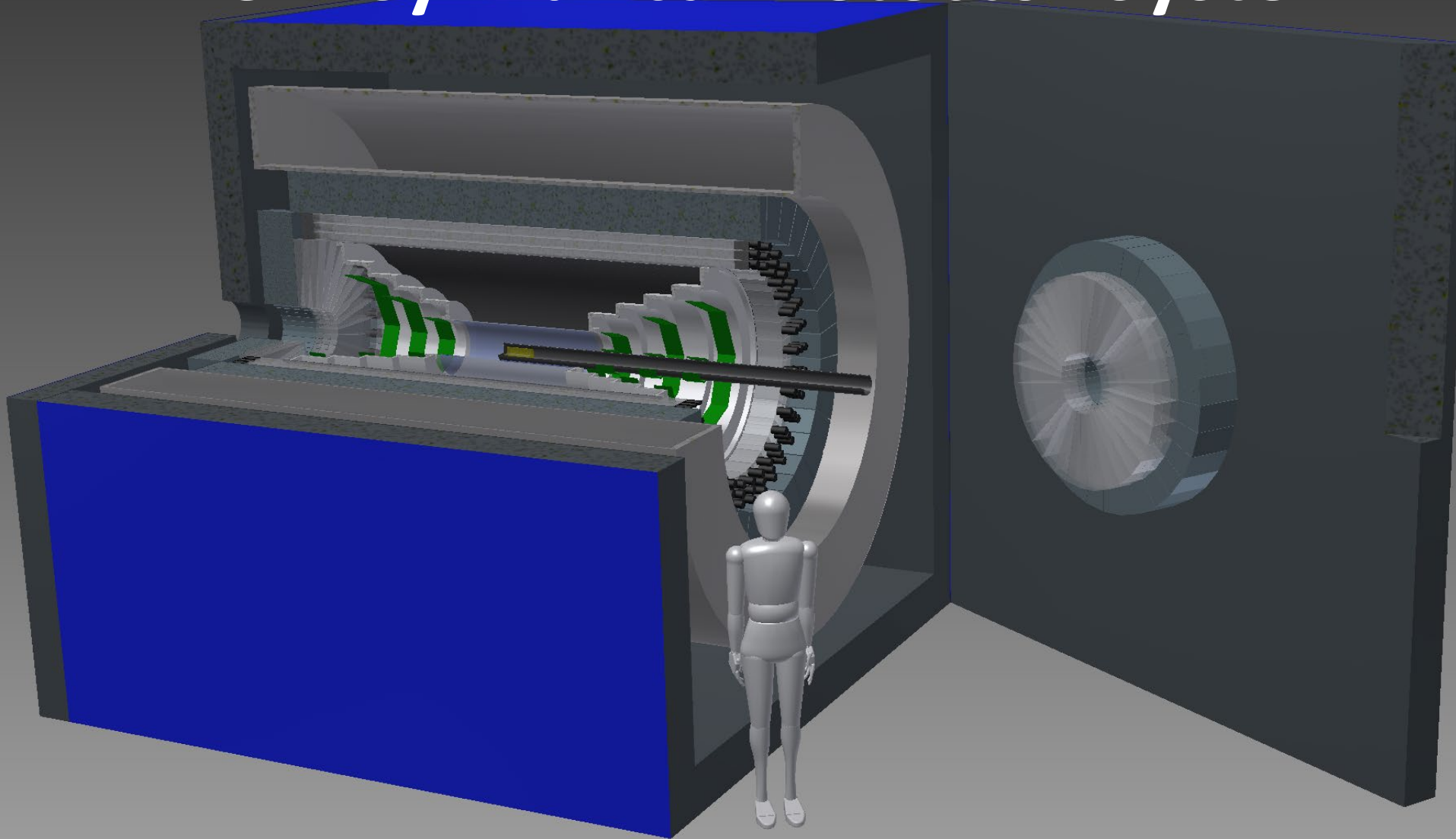
- detect/identify all particles to specify the reaction

- \square high-intensity kaon beam **\leftarrow modified K1.8BR****

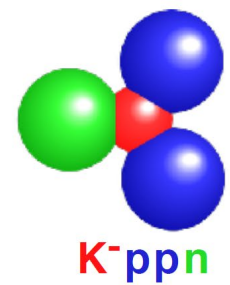
- more K^- yield than the existing beamline

- We take a **step-by-step** approach*

A New Cylindrical Detector System



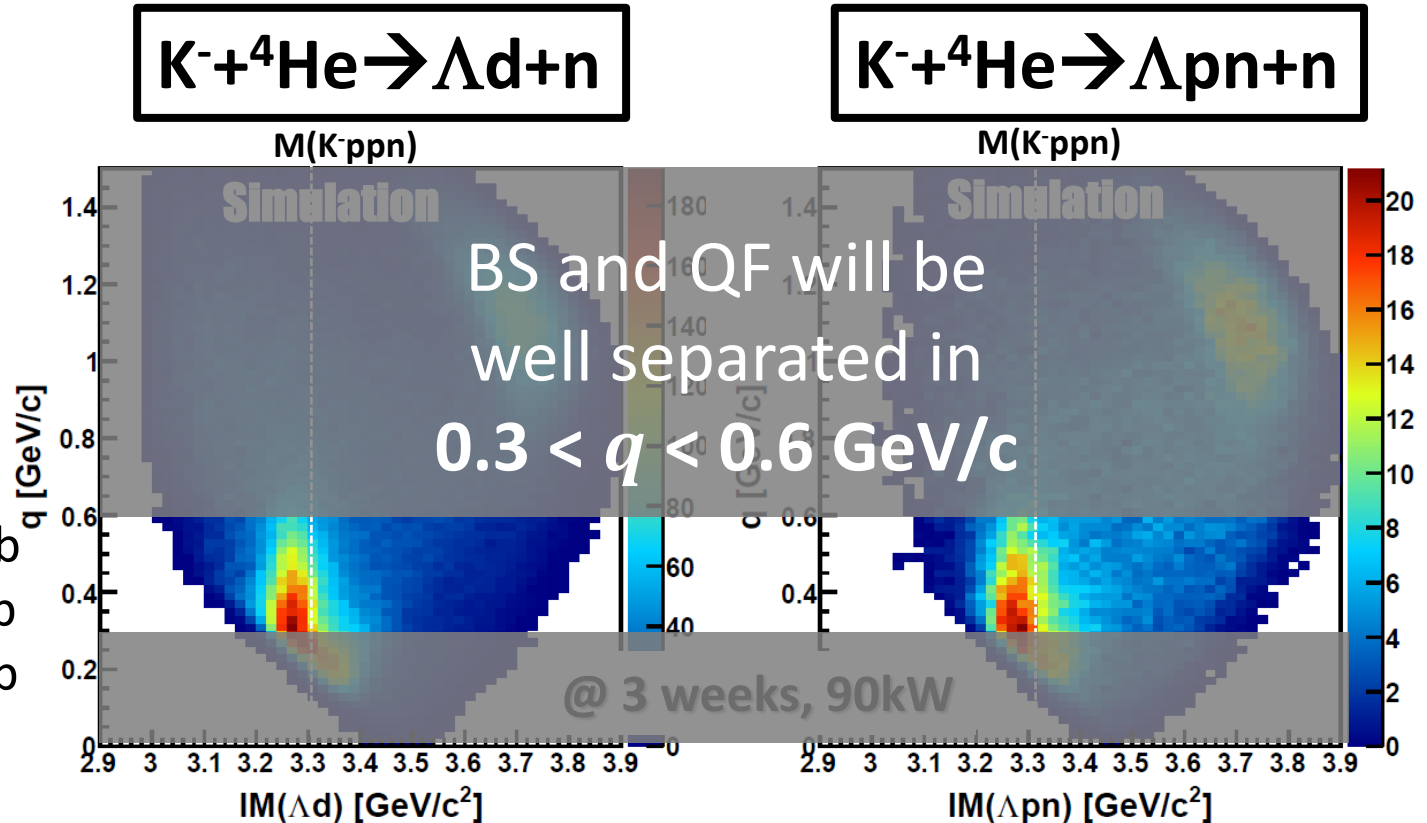
A new 4π spectrometer with n/ γ detection capability



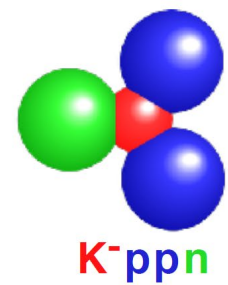
Expected $\bar{K}NNN$ in $K^-+{}^4\text{He}$

$B_{Kppn} \sim 40 \text{ MeV}$
 $\Gamma_{Kppn} \sim 100 \text{ MeV}$
 $Q_{Kppn} \sim 400 \text{ MeV}/c$

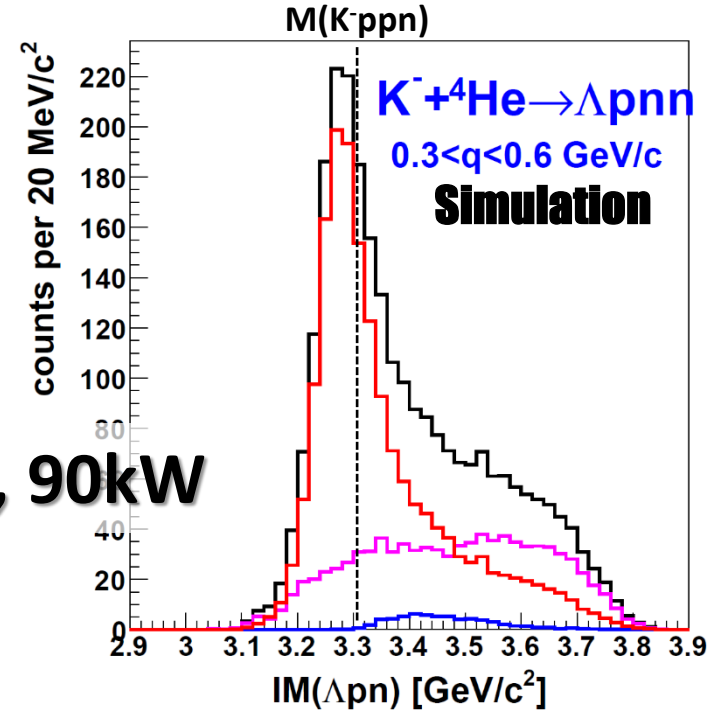
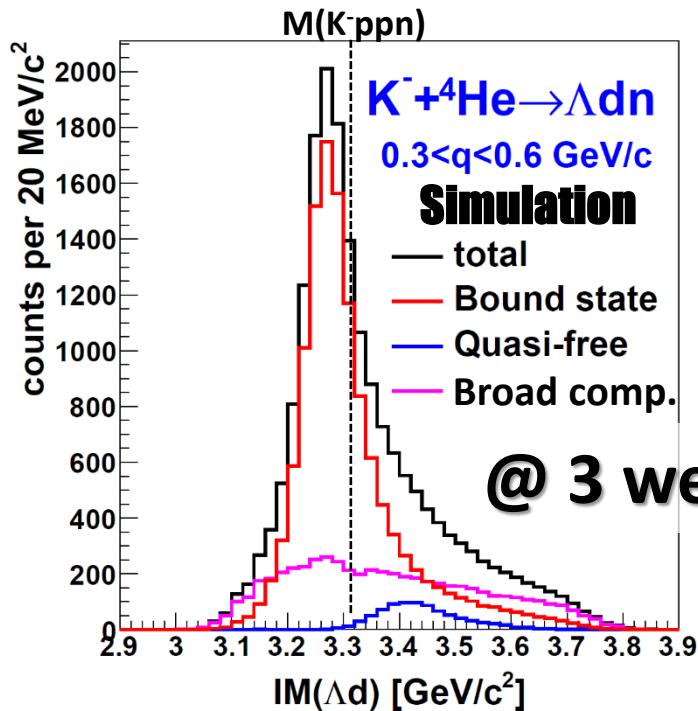
$\sigma(K^-ppn) * Br \sim 10 \mu\text{b}$
 $\sigma(\text{QF}) \sim 10 \mu\text{b}$
 $\sigma(\text{BG}) \sim 20 \mu\text{b}$



- Similar parameters obtained with the $K^-+{}^3\text{He} \rightarrow \Lambda pn$ (PRC102(2020)044002.) are adopted to $K^-ppn/\text{QF}/\text{BG}$ shapes
- “ K^-ppn ” signal [q -independent] will be seen clearly



Expected $\bar{K}NNN$ in $K^-+^4\text{He}$



- Assumption: similar parameters obtained at E15
- **Mass-number dependence of the kaonic nuclei will be provided for the first time.**

Summary

- We observed the “K⁻pp” bound state in ${}^3\text{He}(\text{K}^-, \Lambda\text{p})\text{n}$

J-PARC E15

✓ PLB789(2019)620., PRC102(2020)044002.

- As the next step, the new project has been launched to reveal the properties of the light kaonic nuclei from the $\bar{K}N$ to $\bar{K}NNNN$

- a powerful probe to understand low energy QCD
- the best approach to cold & high-density nuclear matter

- We take a step-by-step approach:

J-PARC E80

– a $\bar{K}NNN$ search via ${}^4\text{He}(\text{K}^-, \text{N})$ reactions as the first step
























J-PARC P89

– a spin/parity measurement of the $\bar{K}NN$ as the second step

– experimental challenges of $\bar{K}N$, $\bar{K}NNNN$, and $\bar{K}\bar{K}NN$ will also be followed

J-PARC E15 Collaboration

S. Ajimura^a, H. Asanoⁿ, G. Beer^b, C. Berucci^f, H. Bhang^c, M. Bragadireanu^e, P. Buehler^f, L. Busso^{g,h}, M. Cargnelli^f, S. Choi^c, C. Curceanu^d, S. Enomoto^o, H. Fujioka^m, Y. Fujiwara^k, T. Fukuda^l, C. Guaraldo^d, T. Hashimoto^u, R. S. Hayano^k, T. Hiraiwa^a, M. Iio^o, M. Iliescu^d, K. Inoue^a, Y. Ishiguro^j, T. Ishikawa^k, S. Ishimoto^o, K. Itahashiⁿ, M. Iwai^o, M. Iwasaki^{m,n*}, K. Kanno^k, K. Kato^j, Y. Katoⁿ, S. Kawasakiⁱ, P. Kienle^{+p}, H. Kou^m, Y. Maⁿ, J. Marton^f, Y. Matsuda^q, Y. Mizoi^l, O. Morra^g, T. Nagae^{j,s}, H. Noumi^a, H. Ohnishi^w, S. Okadaⁿ, H. Outaⁿ, K. Piscicchia^d, Y. Sada^a, A. Sakaguchiⁱ, F. Sakumaⁿ, M. Sato^o, A. Scordo^d, M. Sekimoto^o, H. Shi^d, K. Shirotori^a, D. Sirghi^{d,e}, F. Sirghi^{d,e}, K. Suzuki^f, S. Suzuki^o, T. Suzuki^k, K. Tanida^u, H. Tatsuno^v, M. Tokuda^m, D. Tomono^a, A. Toyoda^o, K. Tsukada^r, O. Vazquez Doce^{d,p}, E. Widmann^f, T. Yamagaⁿ, T. Yamazaki^{k,n}, H. Yim^t, Q. Zhangⁿ, and J. Zmeskal^f

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J-PARC E80 Collaboration



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**We're looking for
new collaborators!**



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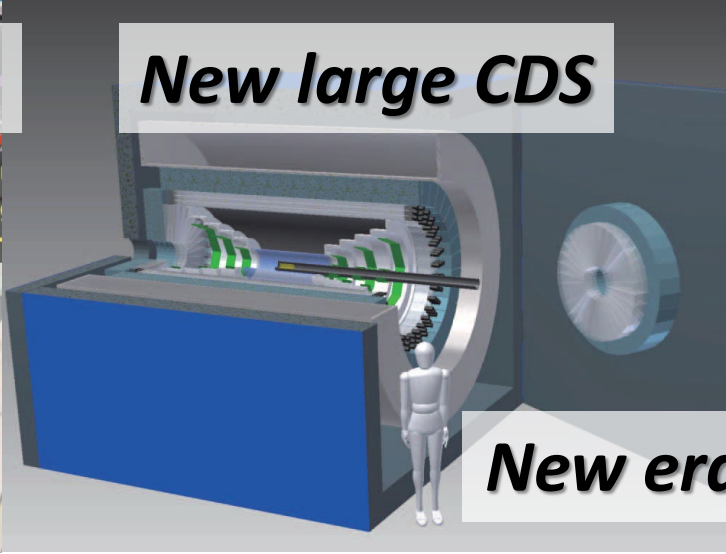
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D. Sirghi, F. Sirghi
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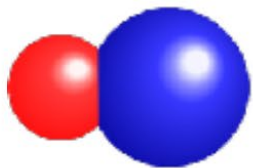


P. Buehler, M. Simon, E. Widmann, J. Zmeskal
Stefan-Meyer-Institut für subatomare Physik, A-1090 Vienna, Austria

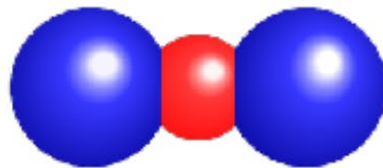


Thank you for your attention!

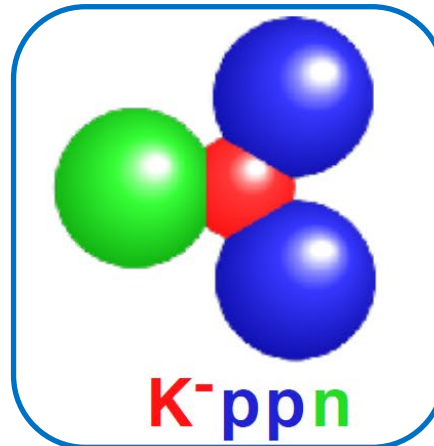
A first step of the new project



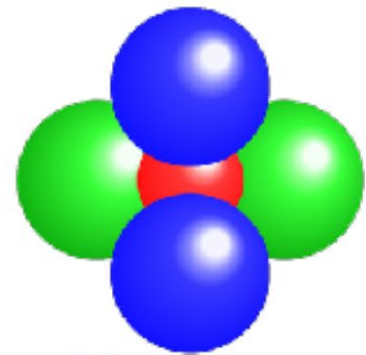
K^-p



K^-pp



K^-ppn



K^-ppnn

via in-flight $^4\text{He}(K^-,N)$

Investigation of fundamental properties of the $\bar{K}NN$ state

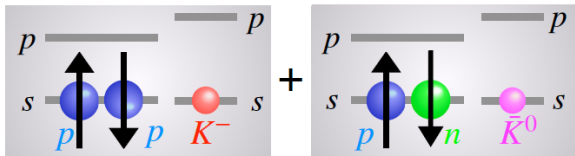
J-PARC P89

Internal structure & spin-parity

There are two possible J^P as for the $\bar{K}NN$ ground state.

" $(NN)_{(I.sym \times S.asym)} \otimes \bar{K}$ "

$$J^P = 0^-$$

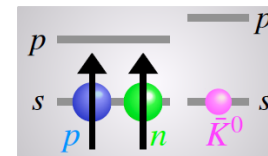


$$\frac{|I_{\bar{K}N} = 0|}{|I_{\bar{K}N} = 1|} = \frac{3}{1}$$

Deeper bound expected

" $(NN)_{(I.asym \times S.sym)} \otimes \bar{K}$ "

$$J^P = 1^-$$



$$\frac{|I_{\bar{K}N} = 0|}{|I_{\bar{K}N} = 1|} = \frac{1}{3}$$

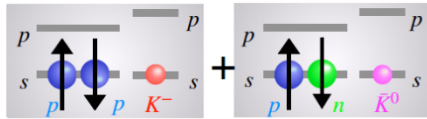
Shallower bound expected

* Positive parity state should be higher excited state if exist.

How to determine J^P

– Λp spin-spin correlation ($\alpha_{\Lambda p}$) in $K^-pp \rightarrow \Lambda p$ decay –

$$J^P = 0^-$$



To make negative parity from Λp

$$L_{\Lambda p} = 1$$

To be $J = 0$

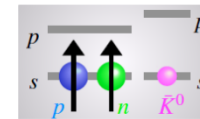
$$S_{\Lambda p} = 1$$



$$\alpha_{\Lambda p} = +1$$

Spin parallel

$$J^P = 1^-$$



To make negative parity from Λp

$$L_{\Lambda p} = 1$$

To be $J = 1$

$$BR = 1/3 \quad S_{\Lambda p} = 0 \quad + \quad S_{\Lambda p} = 1 \quad BR = 2/3$$



$$\alpha_{\Lambda p} = -1$$

Spin anti-parallel



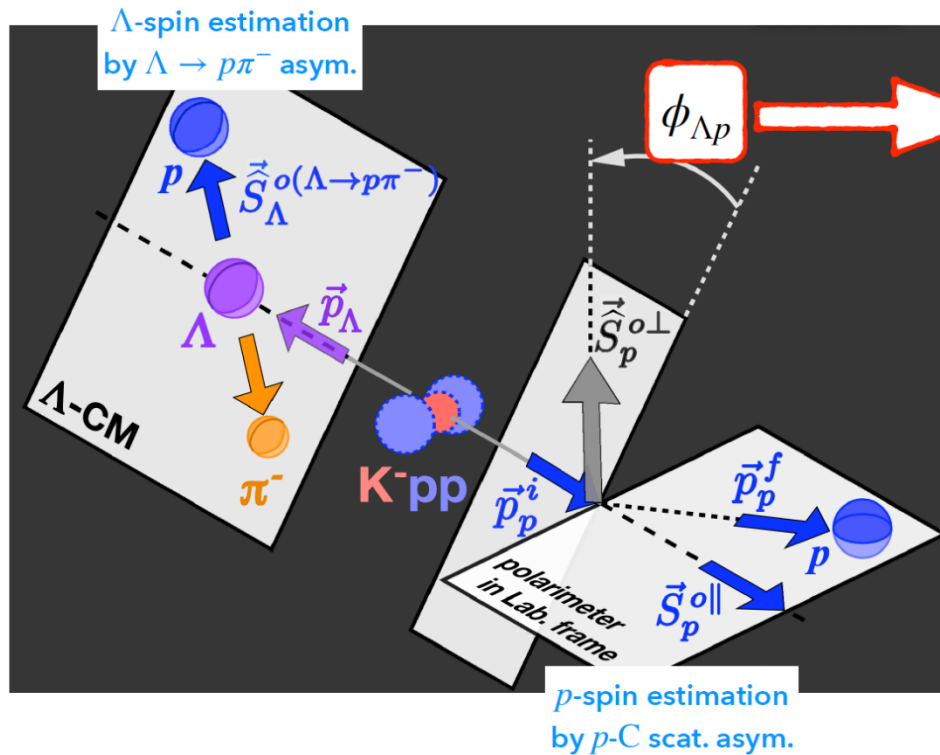
$$\alpha_{\Lambda p} = +1$$

Spin parallel

We can deduce J^P from $\alpha_{\Lambda p}$ measurement.

How to measure spin-spin correlation

– Spin alignment measurement by $\Lambda \rightarrow p\pi^-$ & p -C scattering –



Spin-spin correlation on ϕ -asymmetry

$$N(\phi_{\Lambda p}) = N_0 \cdot (1 + r^{(J^P)} \cdot \alpha_{\Lambda p} \cos \phi_{\Lambda p})$$

$r^{(J^P)}$: asymmetry reduction factor defined by;

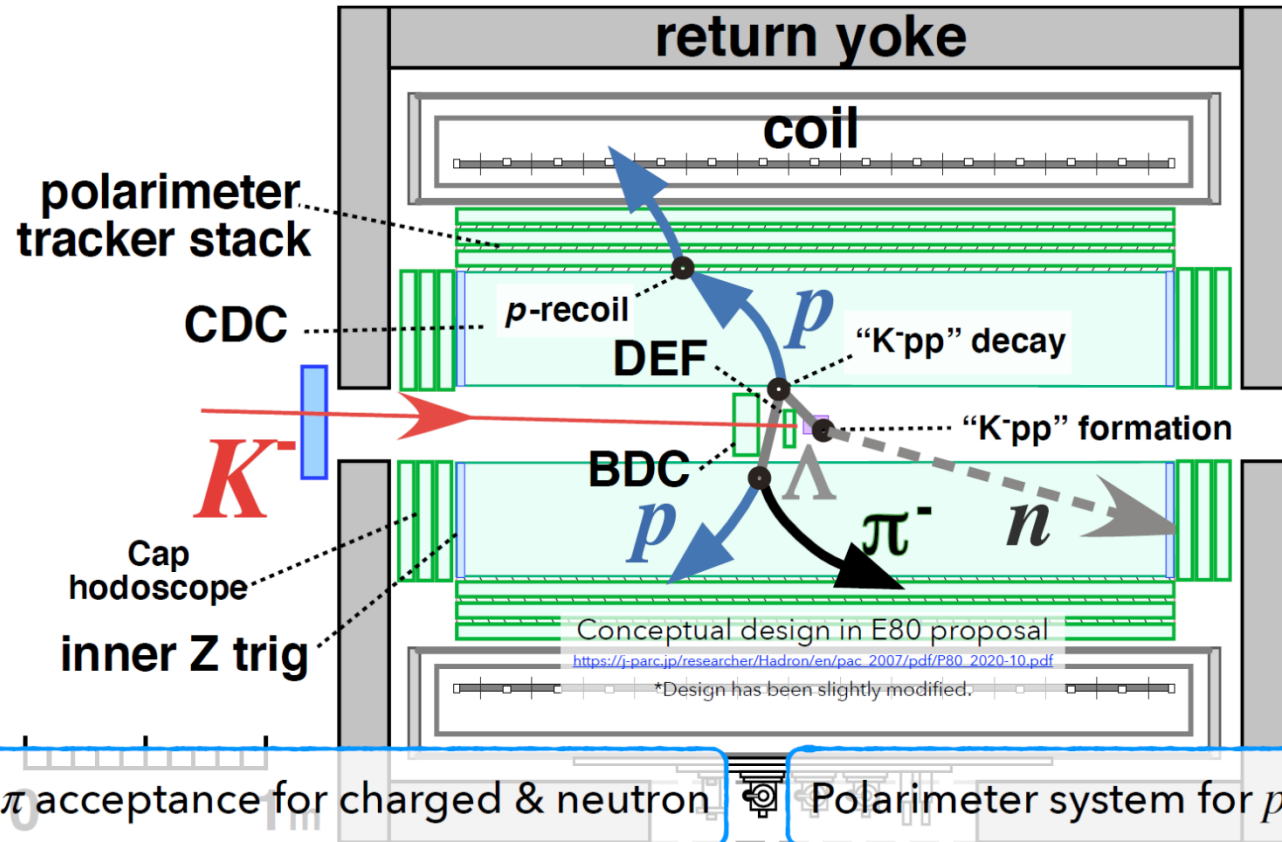
α_- : Λ asym. parameter B : Magnetic field

A_{pC} : Analyzing power $B_{\bar{K}}$: Binding energy

$f_{\bar{S}_\Lambda}$: Spin distribution q : Momentum transfer

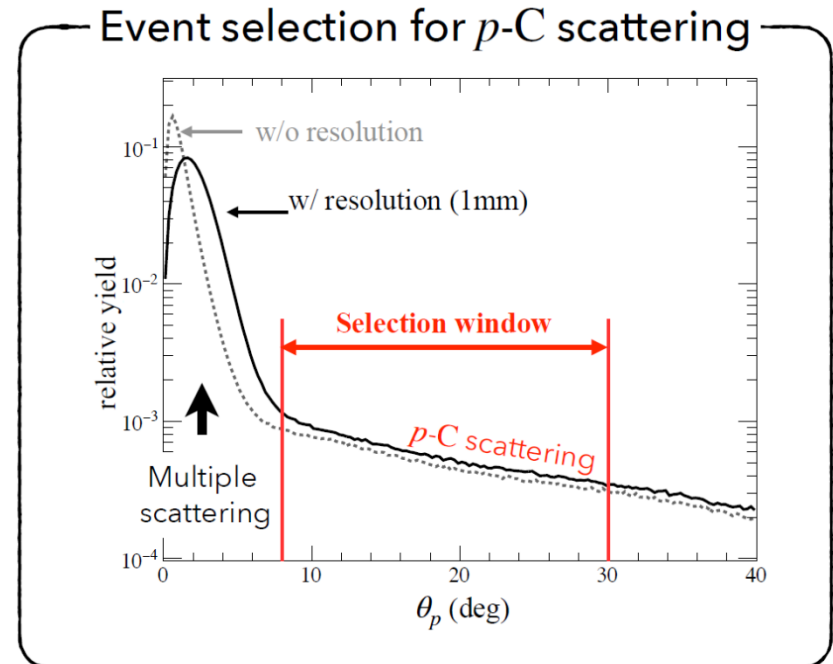
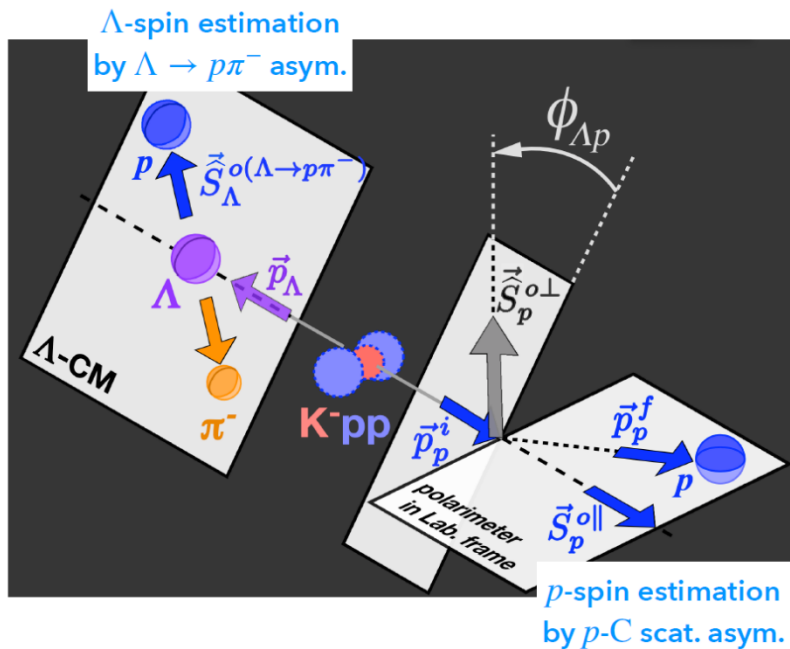
We can observe $\alpha_{\Lambda p}$ from $\phi_{\Lambda p}$ -distribution.

Cylindrical detector system



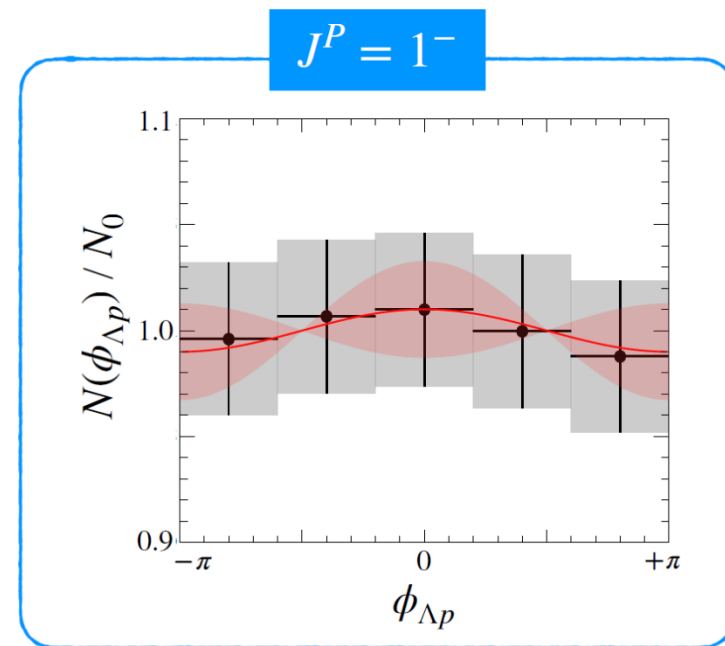
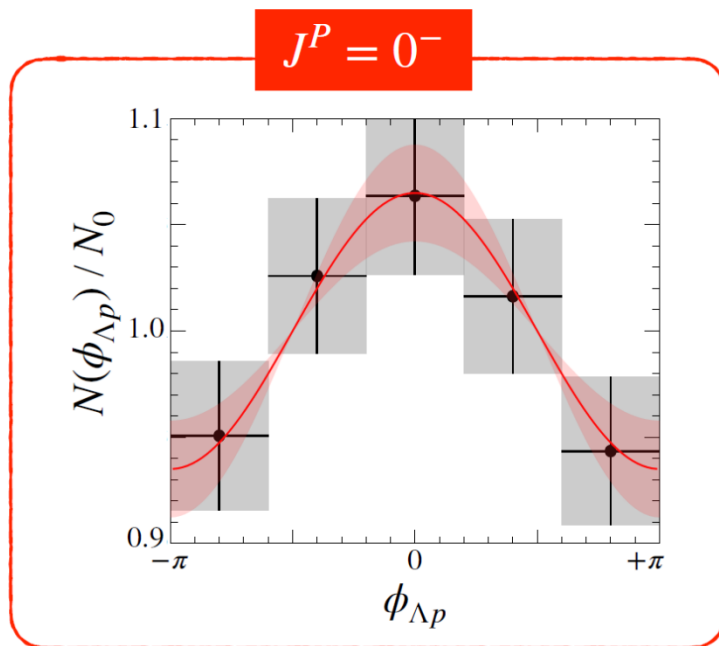
Expected spectra

– To measure $\phi_{\Delta p}$ -asymmetry for J^P determination –



Expected spectra

– To measure $\phi_{\Lambda p}$ -asymmetry for J^P determination –



We would exclude $J^P = 1^-$ with **95%** confidence level from **only $\phi_{\Lambda p}$ -asym.**

Size of “K⁻pp”?

Fit with PWIA $\sigma(M, q) \propto \rho(M, q) \times \frac{(\Gamma_{Kpp}/2)^2}{(M - M_{Kpp})^2 + (\Gamma_{Kpp}/2)^2} \times \exp\left(-\frac{q^2}{Q_{Kpp}^2}\right)$

$B_{Kpp} \sim 40$ MeV

→ large binding energy

$Q_{kpp} \sim 400$ MeV

→ wide momentum transfer

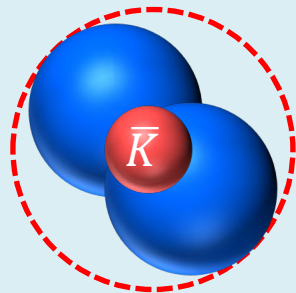
suggest the “K⁻pp” is quite compact ($R_{Kpp} \sim 0.6$ fm)

→ *Need more realistic theoretical calculations*

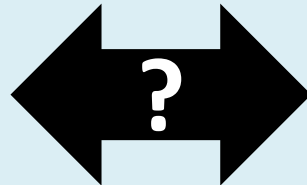
**a whole new probe to reveal
the structure of a nucleon**

Radius of “K⁻pp”

Interaction range of “K⁻”



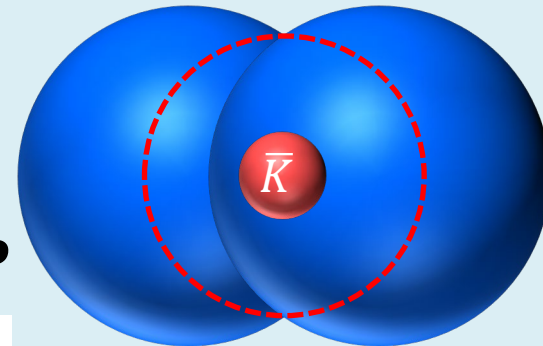
~ 0.6 fm



& effect of K size?

Charge radius

$$\langle r \rangle = 0.560 \pm 0.031 \text{ fm}$$



~ 0.8 fm

~ 0.6 fm

A Theoretical Interpretation

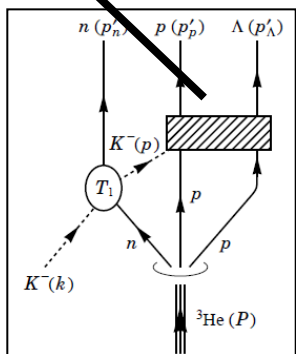
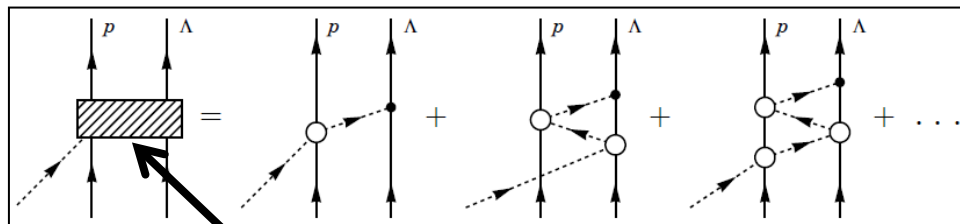
- A calculation with chiral unitary approach reproduces the mass spectrum with the $\bar{K}NN$

PTEP

Prog. Theor. Exp. Phys. **2016**, 123D03 (27 pages)
DOI: 10.1093/ptep/ptw166

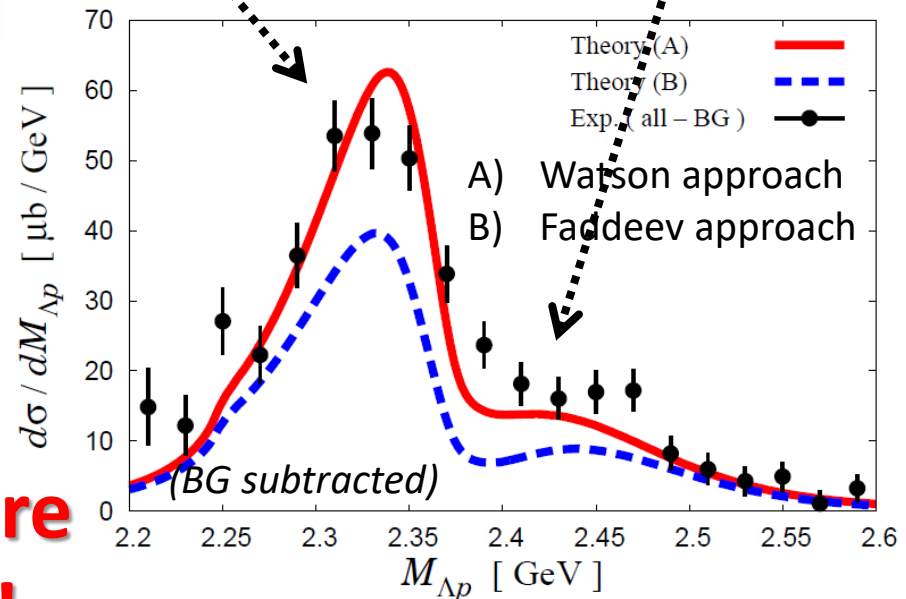
On the structure observed in the in-flight ${}^3\text{He}(K^-, \Lambda p)n$ reaction at J-PARC

Takayasu Sekihara^{1,*}, Eulogio Oset², and Angels Ramos³



Theoretical investigations are indispensable!

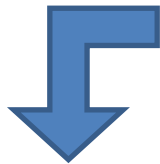
$\bar{K}^{\text{bar}}NN$ bound-state quasi-elastic kaon scattering



JPS Conf. Proc.26(2019)023009.

Many Questions to be Answered

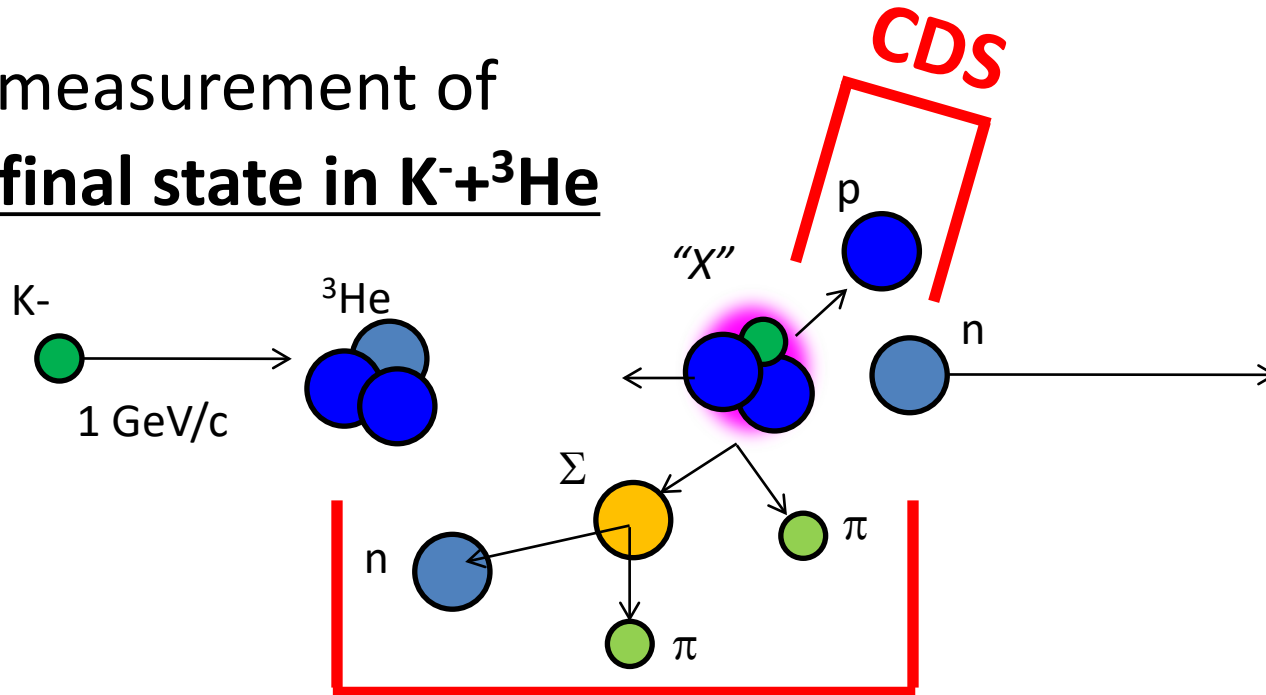
- **Further details of the $\bar{K}NN$**
 - Spin and parity of the “K⁻pp”?
 - Really compact and dense system?
 - Other decay modes?



- **$\pi\Sigma N$ mesonic decay**
 - expected to be the dominant channel
 - only YN non-mesonic decays were reported
- **Reaction mechanism**
 - relation between $\Lambda(1405)$ =“K⁻p” & “K⁻pp”
 - “K⁻pp” is expected to be produced via $\Lambda(1405)+p \rightarrow$ “K⁻pp” door-way process

$K^- \ ^3\text{He} \rightarrow \pi \Sigma \text{pn} @ \text{E15}$

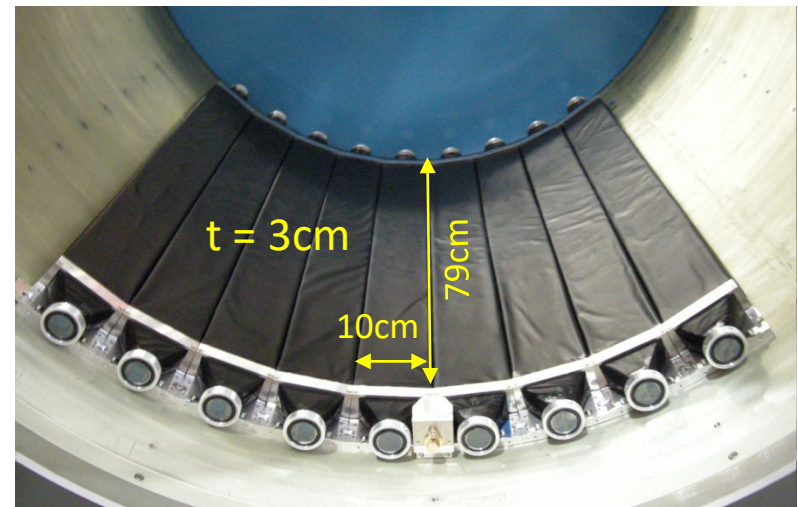
- Exclusive measurement of $\pi^\pm \Sigma^\mp \text{pn}$ final state in $K^- + ^3\text{He}$



CDS

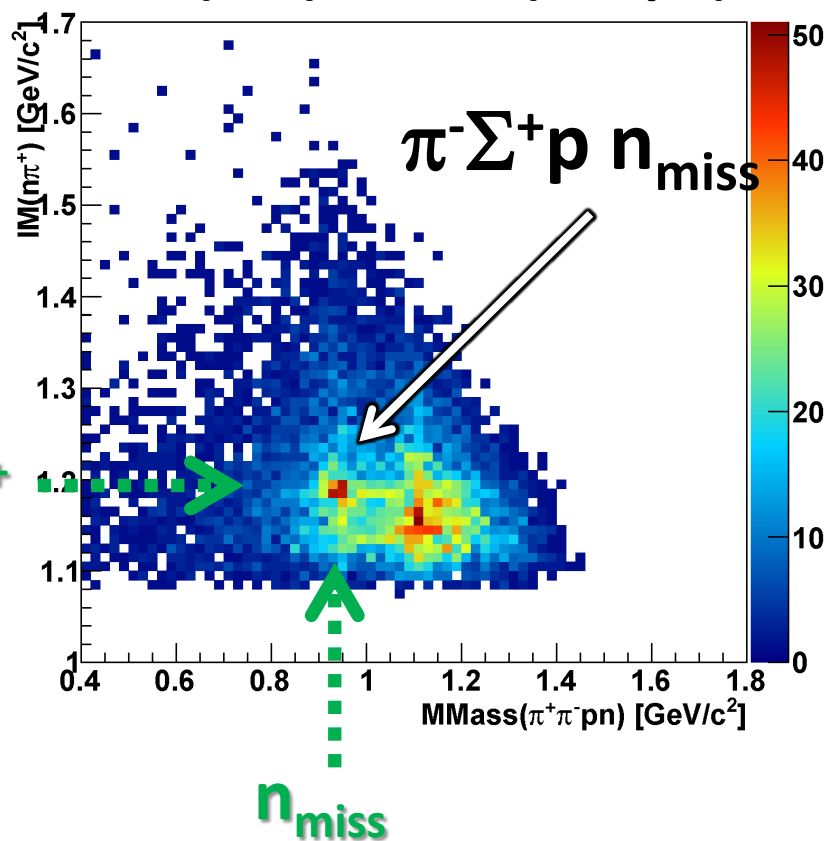
- Experimental challenge of neutron detection with thin scintillation counter ($t=3\text{cm}$)

n detection efficiency $\sim 3\text{-}10\%$

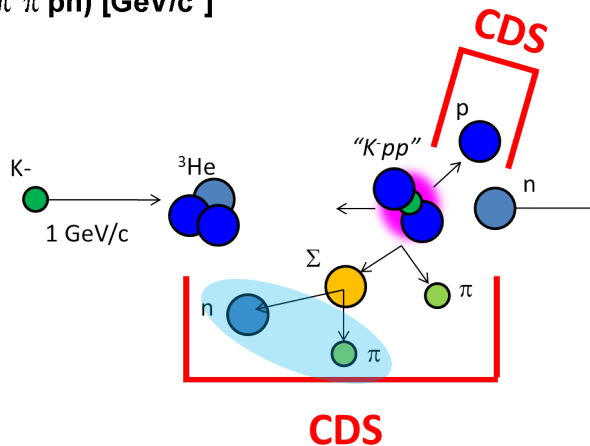
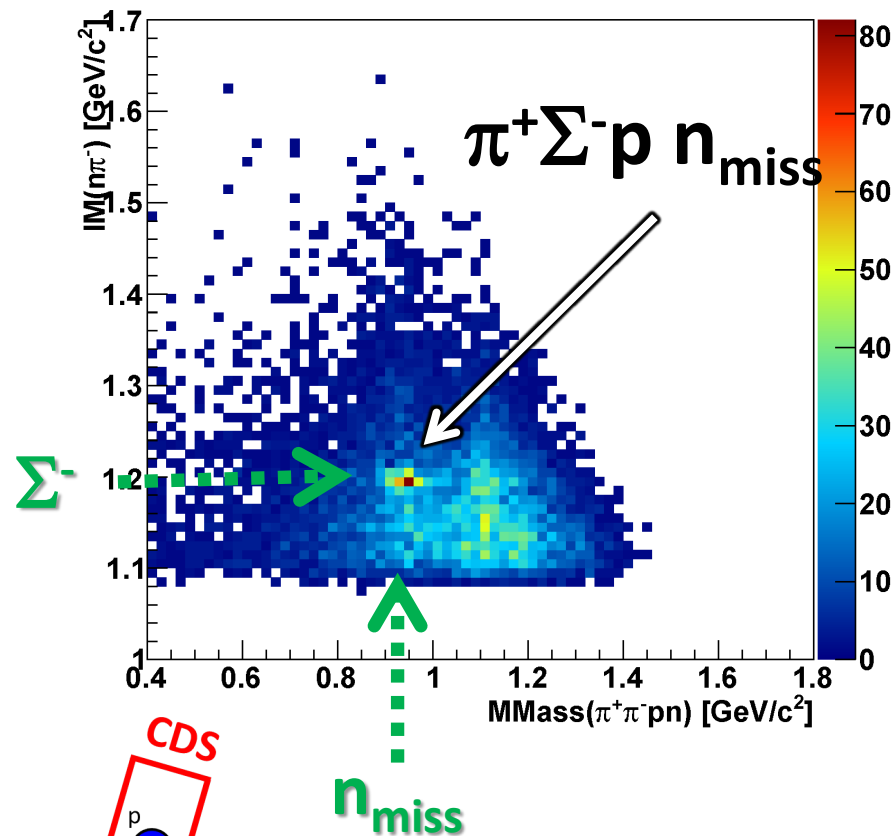


$\pi\Sigma\rho n$ Events

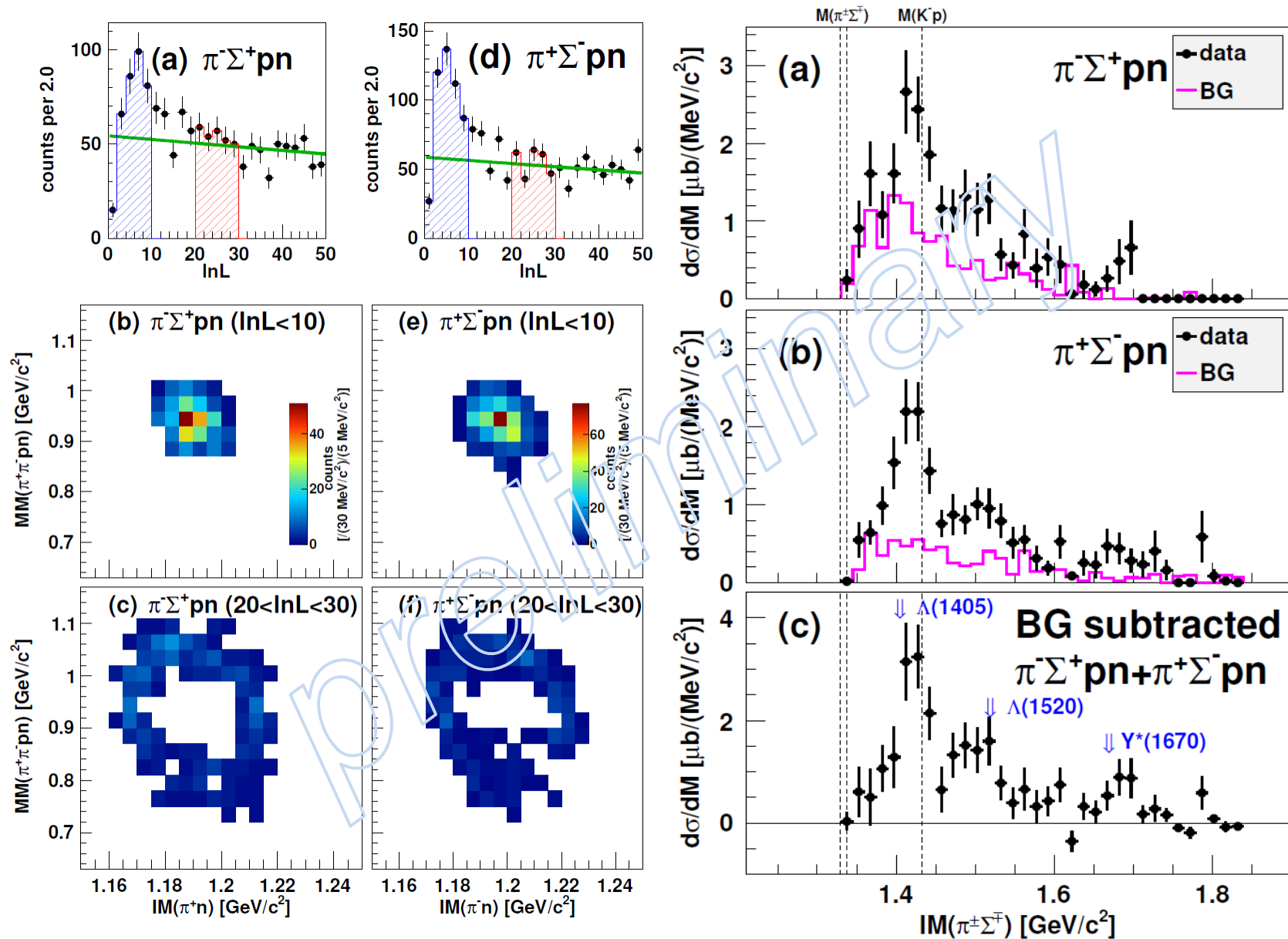
IM($n\pi^+$) vs MM($\pi^+\pi^-pn$)



IM($n\pi^-$) vs MM($\pi^+\pi^-pn$)



BG Subtracted IM($\pi^\pm\Sigma^\mp$) in $\pi^\pm\Sigma^\mp pn$



Y^* pn Final State

$\Lambda(1405)$

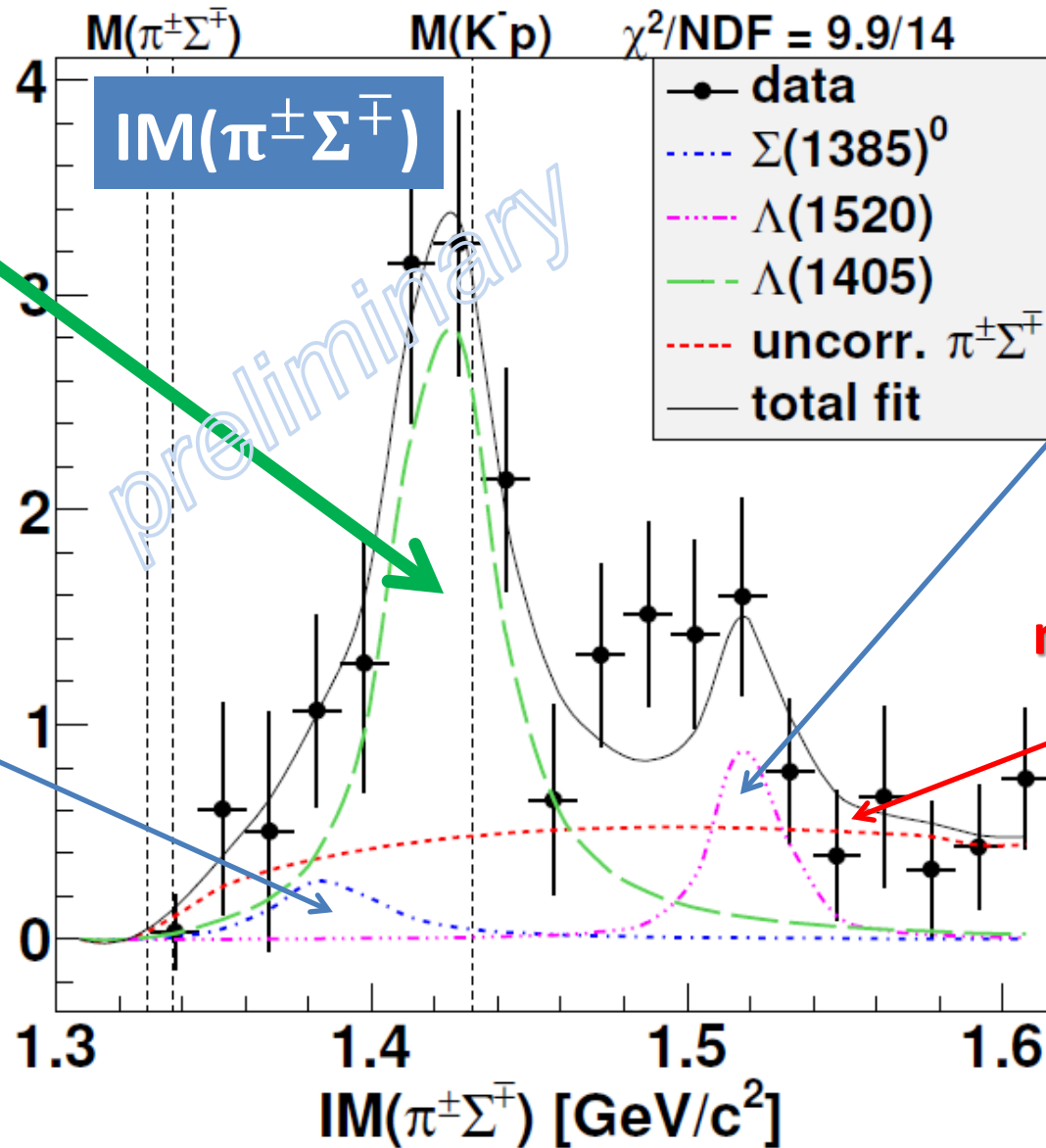
$\sim 250 \mu\text{b}$

Fit with BW:
 $M \sim 1422 \text{ MeV}$
 $\Gamma \sim 40 \text{ MeV}$

$\Sigma(1385)^0$

$\sim 100\text{-}150 \mu\text{b}$
*[evaluated from
 $\Sigma(1385)^{\pm} \rightarrow \pi^{\pm} \Lambda$
 measurement]*

$(\Sigma(1385) \rightarrow \pi\Lambda/\pi\Sigma : 87.0/11.7\%)$



$\Lambda(1520)$

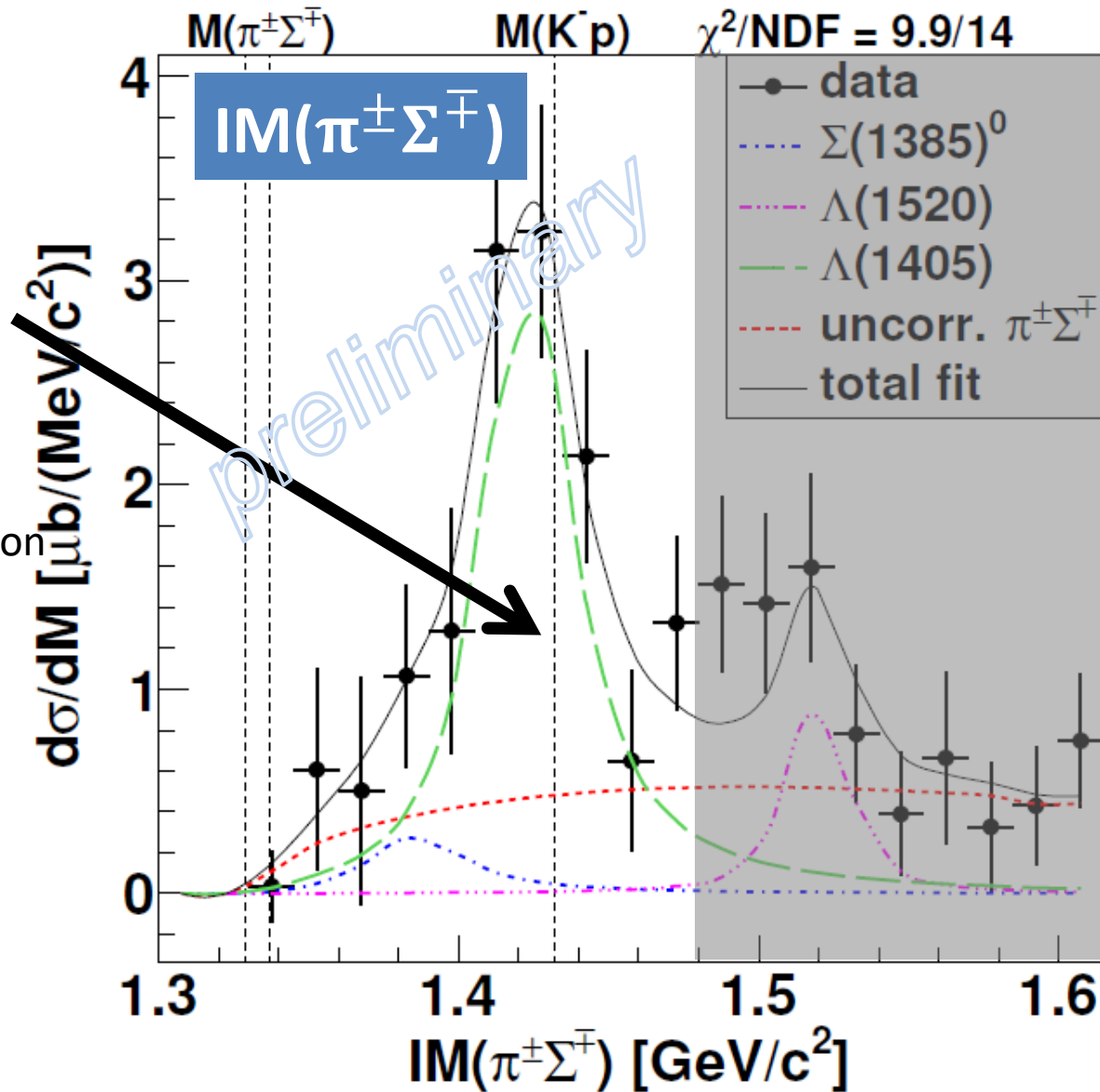
$\sim 100 \mu\text{b}$

$\pi\Sigma pn$
 non-resonant

$\Lambda(1405)pn$ Final State Selection

Select
 $\Lambda(1405)$
region

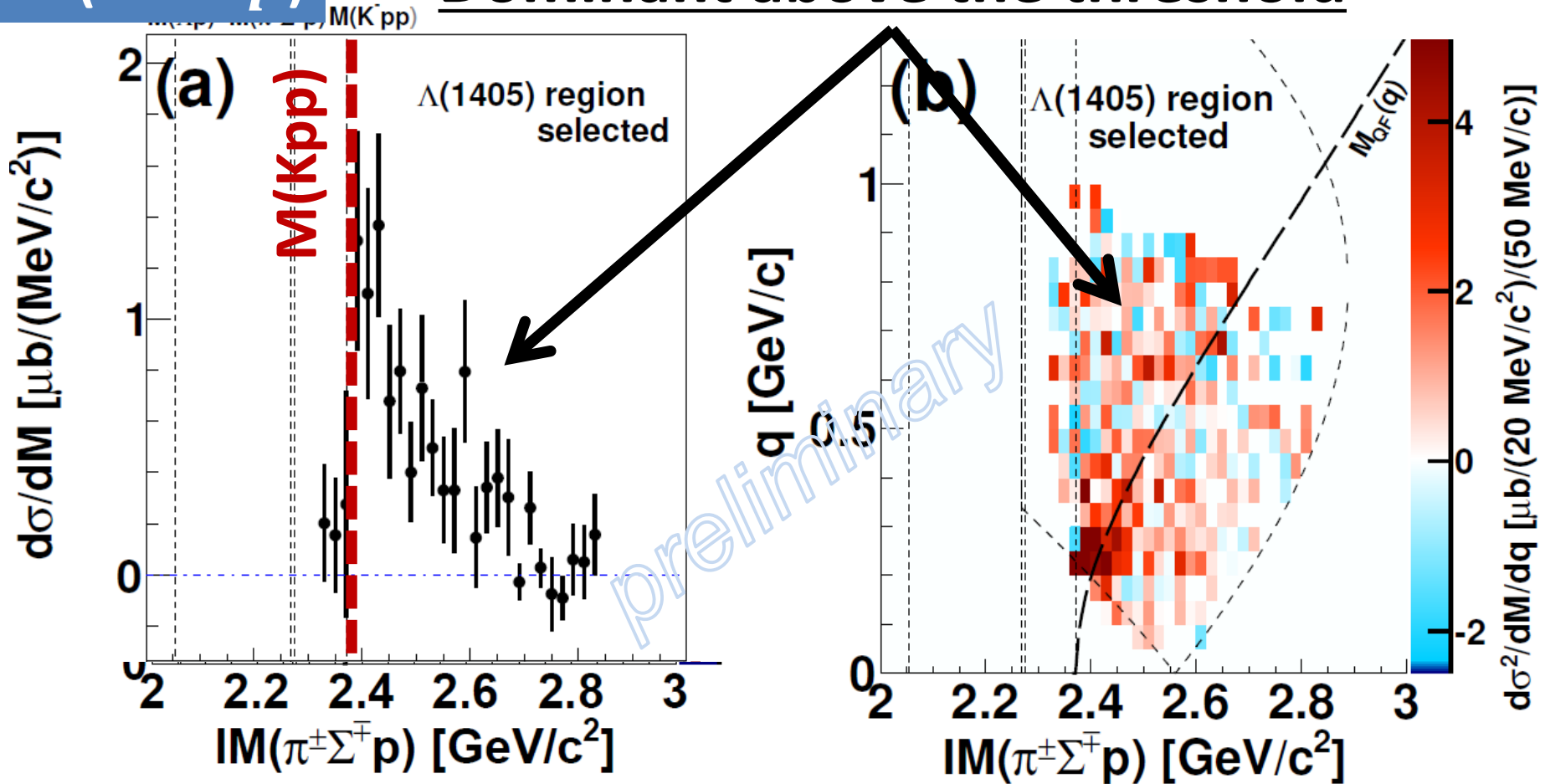
- Below $\Lambda(1520)$
- Small contribution from $\Sigma(1385)$



IM($\pi\Sigma p$) in $\Lambda(1405)pn$ Final State

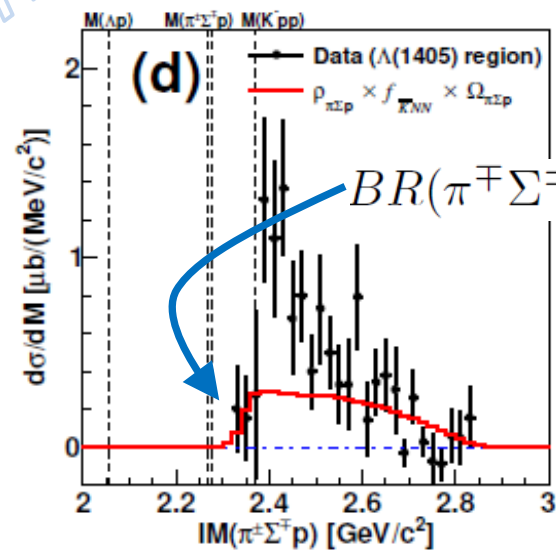
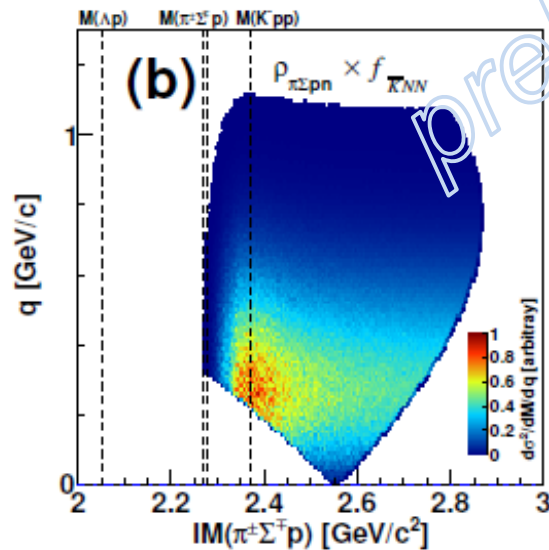
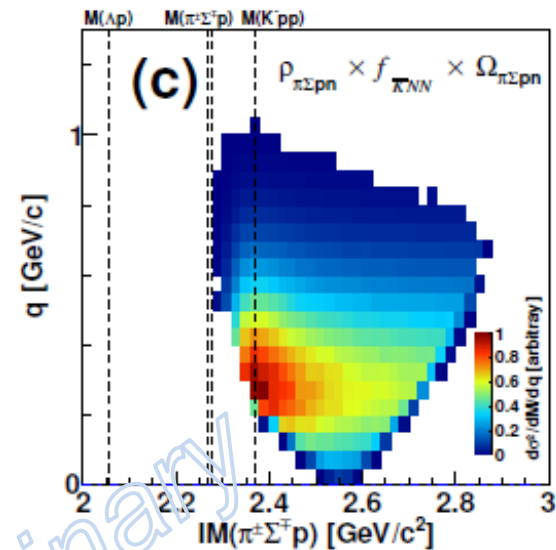
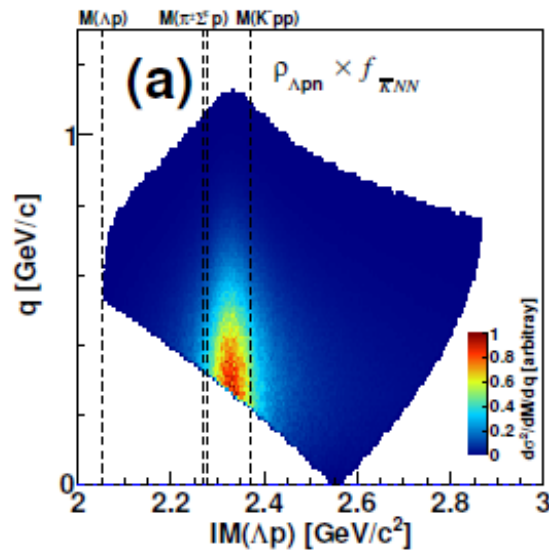
IM($\pi^\pm\Sigma^\mp p$)

Dominant above the threshold



- IM($\Lambda(1405)p$) distributes above the $M(Kpp)$
- QF $K^-N \rightarrow K^{\text{bar}}n$ followed by $K^{\text{bar}}NN \rightarrow \Lambda(1405)p$

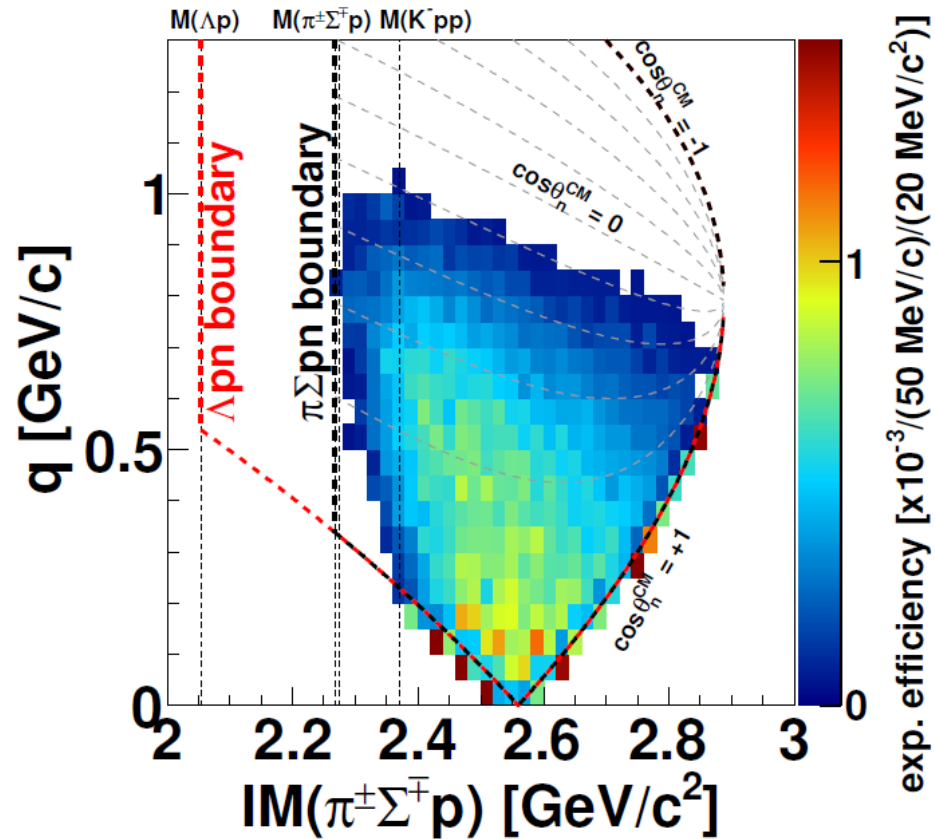
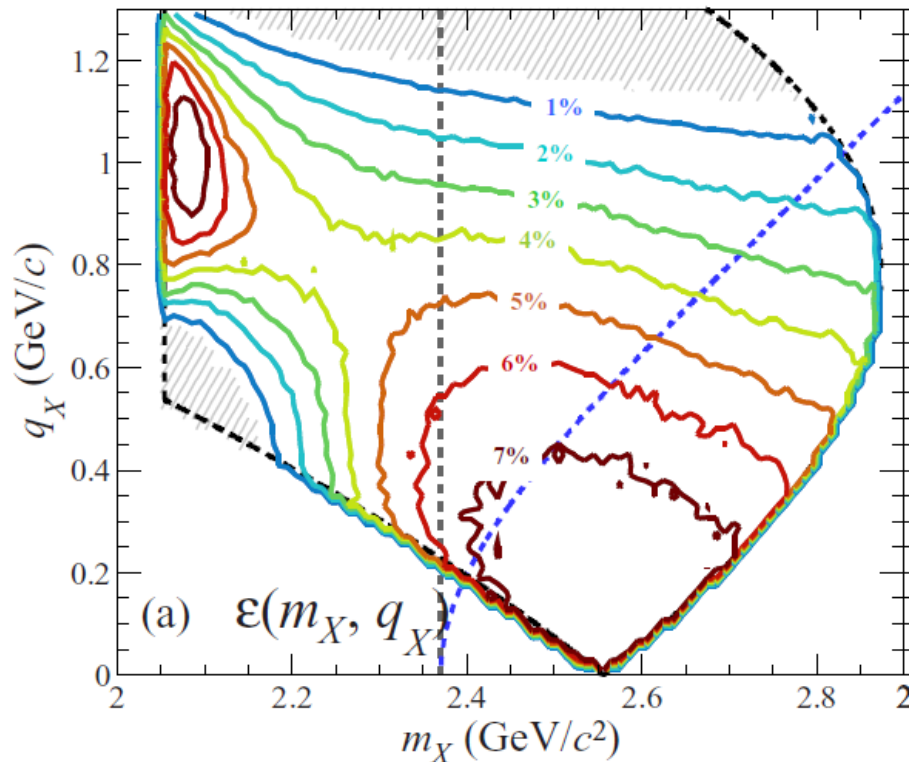
Investigation of $K^-pp \rightarrow \pi\Sigma p$ decay



Detector acceptances of Λp and $\pi\Sigma p$

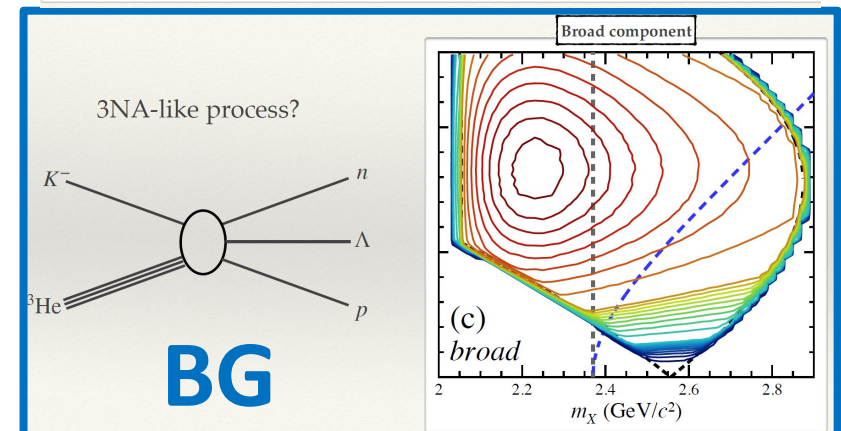
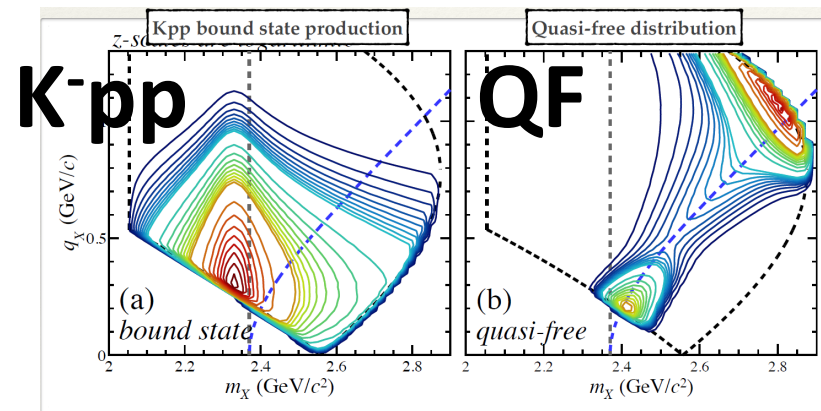
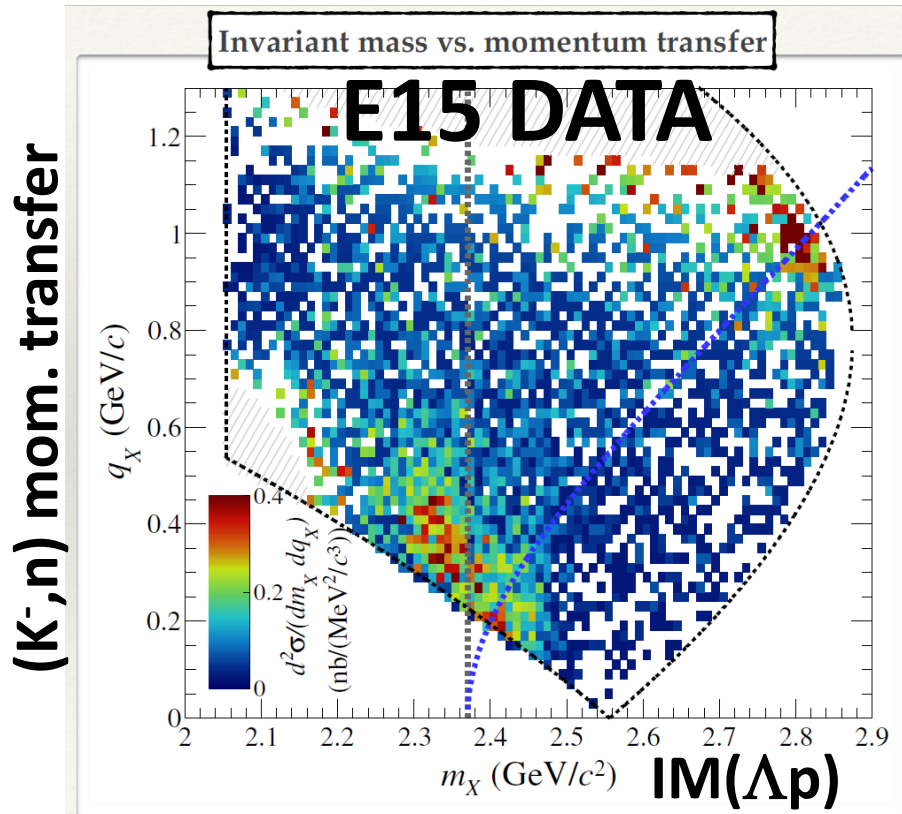
Λp

$\pi\Sigma p$



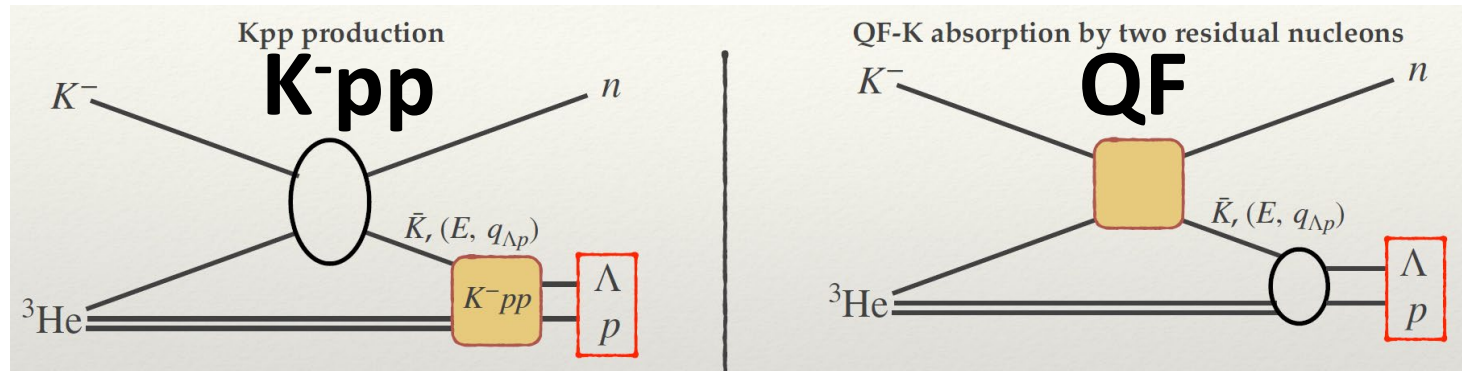
Recent progress of the E15 analysis

- In the E15 analysis so far, we assumed a point-like 3NA process for the background to explain the $\text{IM}(\Lambda p)$ spectrum, by parametrizing a fitting function

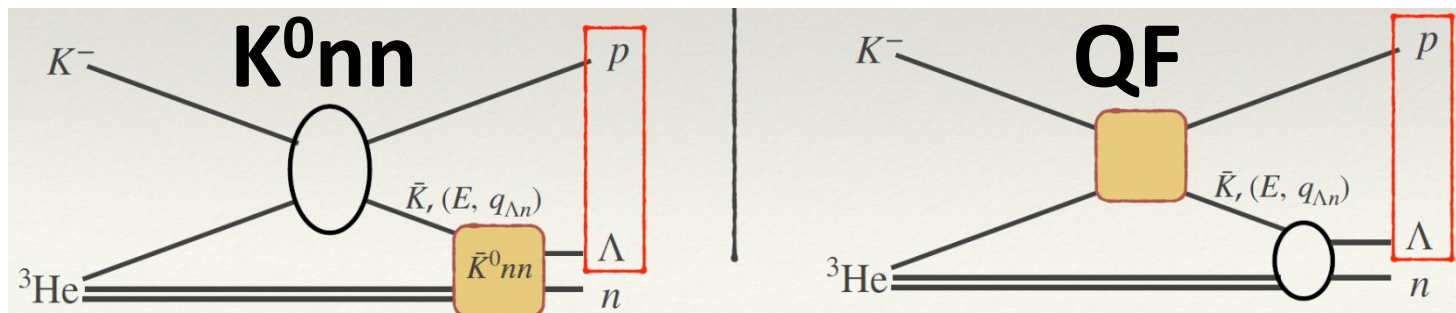


Recent progress of the E15 analysis

- For the K^-pp and QF, we assume the following processes



- On the other hand, “p” and “n” can be swapped in the reactions when the isospin partner of the K^-pp ($=K^0nn$) is also generated

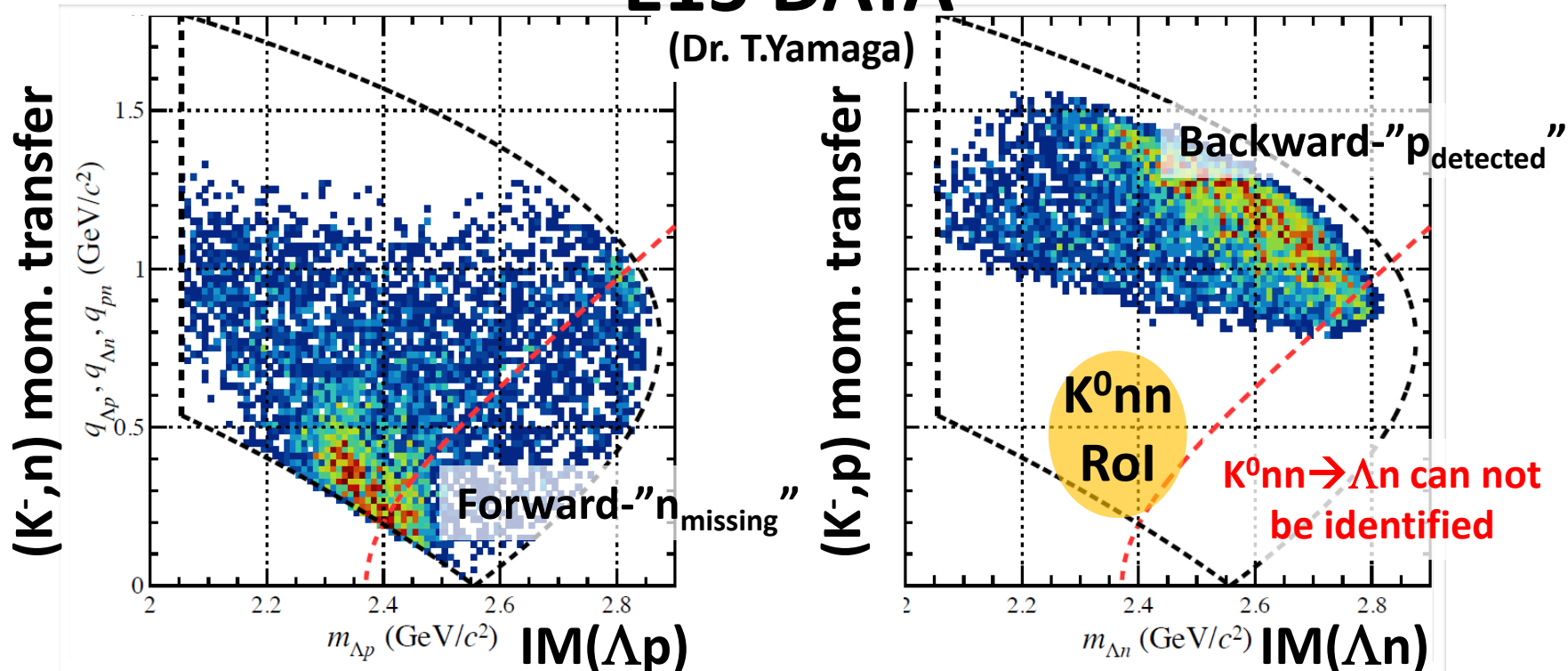


Recent progress of the E15 analysis

- $IM(\Lambda p)$ and $IM(\Lambda n)$
 - Acceptances are quite different between the “ Λp ” and “ Λn ”
 - In $IM(\Lambda n)$, a forward going proton is out of the acceptance

E15 DATA

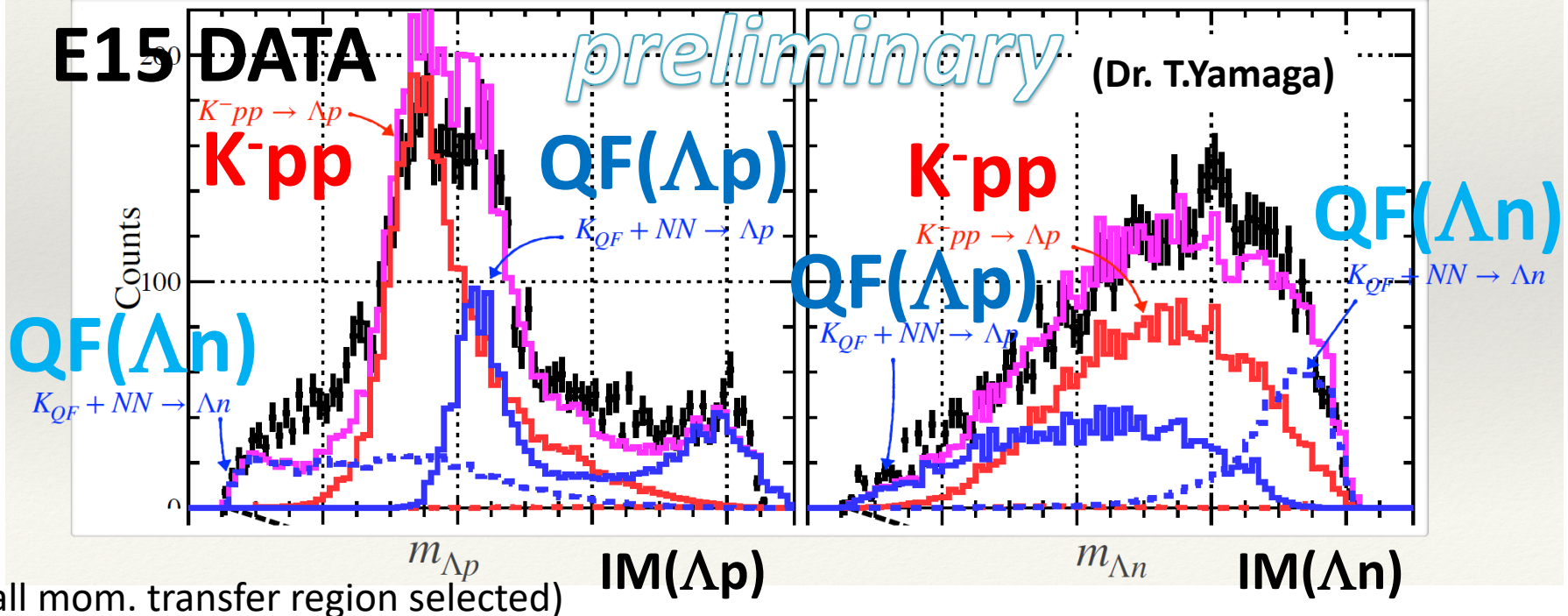
(Dr. T.Yamaga)



Recent progress of the E15 analysis

- Both of $IM(\Lambda p)$ and $IM(\Lambda n)$ can be reproduced by the “signal” and “QFs”
 - Eye fit results \rightarrow further analysis is on going

need further data with the new 4π spectrometer



Σ^*N bound state? Other possibilities?

$\Sigma(1385) 3/2^+$

$$I(J^P) = 1(\frac{3}{2}^+)$$

$\Sigma(1385)$ DECAY MODES

Fraction (Γ_i/Γ)

$\Lambda\pi$

(87.0 \pm 1.5) %

$\Sigma\pi$

(11.7 \pm 1.5) %

- Σ^* coupling through $K^{\text{bar}}-N$ channel (P-wave) would be weak
 - ✓ A.Cieply et al., PRC84(2011)045206, etc.
- Naively, Σ^*N system with $1^+/2^+$ state (S-wave) could not be bound, because corresponding ΔN system (non-strangeness sector) is considered to be no-bound or quite-weakly bound
 - ✓ R. D. Mota et al., PRC59(1999)46, etc.

need J^P determination with a polarimeter

- The $K^{\text{bar}}NN$ state ($I=1/2, J^P=0^-$) is calculated with a $K^{\text{bar}}NN-\pi\Sigma N-\pi\Lambda N$ coupled channel system, where the $\pi\Lambda N$ coupling is expected to be small
- The $K^{\text{bar}}NN$ state with $J^P=1^-$ ($K^{\text{bar}}-d$ like configuration) is expected to not be bound, or have small B.E.
 - ✓ S. Ohnishi et al., PRC95(2017)065202, etc.

Σ^*N bound state? Other possibilities?

- One theoretical possibility is a “ $\pi\Lambda N$ - $\pi\Sigma N$ dibaryon”

Nuclear Physics A 897 (2013) 167–178

Relativistic three-body calculations of a $Y = 1, I = \frac{3}{2},$
 $J^P = 2^+ \pi \Lambda N - \pi \Sigma N$ dibaryon

H. Garcilazo^a, A. Gal^{b,*}

- Calculated $\pi\Lambda N$ resonance with Σ^*N - $\Delta\Sigma$ configuration is:
 - $l=1/2, J^P=2^+ : \mathbf{E = -10-i52 MeV}$
 - $l=3/2, J^P=2^+ : \mathbf{E = -120-i2.6 MeV}$ with respect to $M(K^{\text{bar}}NN)$
- The obtained K^-pp parameter at E15 is **$\mathbf{E=-40-i50 MeV}$**
- Therefore, the “observed K^-pp structure” would be different from the “ $\pi\Lambda N$ - $\pi\Sigma N$ dibaryon”

“K-ppn” Candidates so far

- A few candidates have been reported in *inclusive* measurements

◆ **E471/E549@KEK** ← **NULL results**

– ${}^4\text{He}(\text{stopped-}K^-, p/n)X$

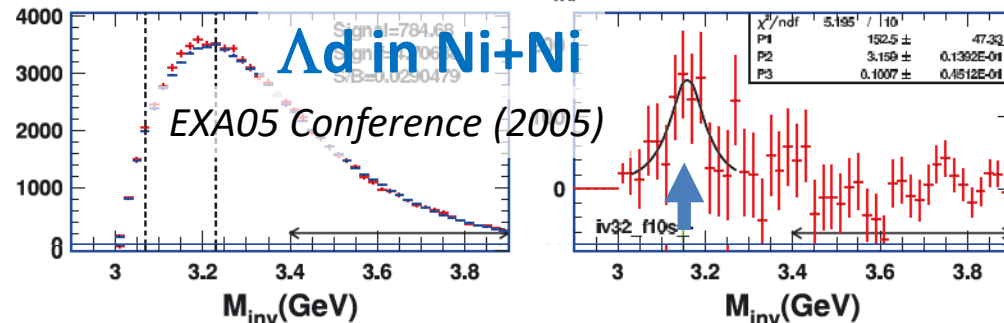
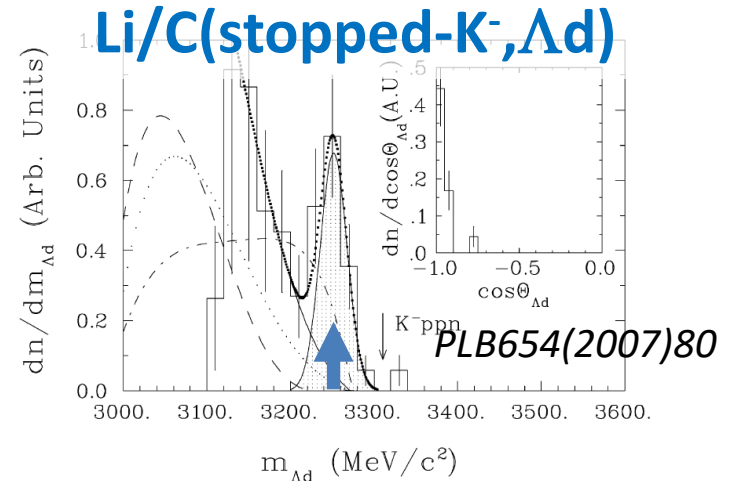
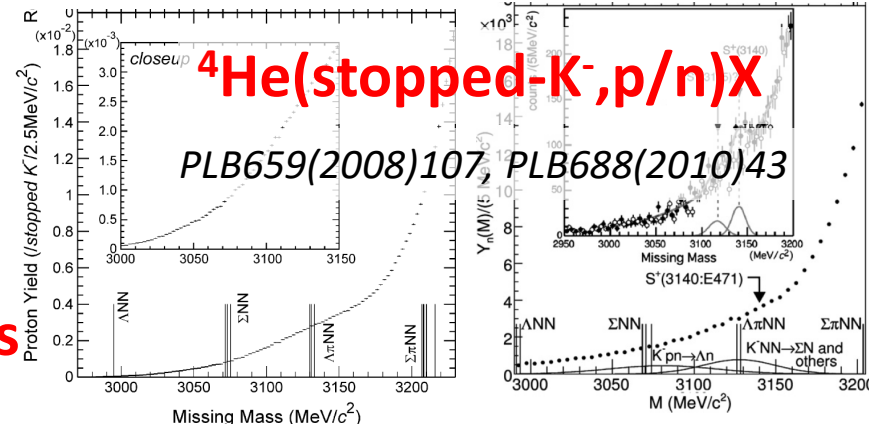
◆ **FINUDA@DAΦNE** ← **Observed?**

– $\text{Li}/\text{C}(\text{stopped-}K^-, \Delta d)$

◆ **FOPI@GSI** ← **Observed?**

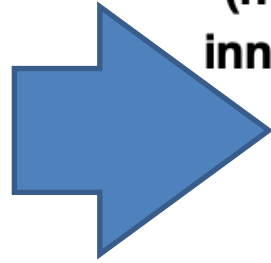
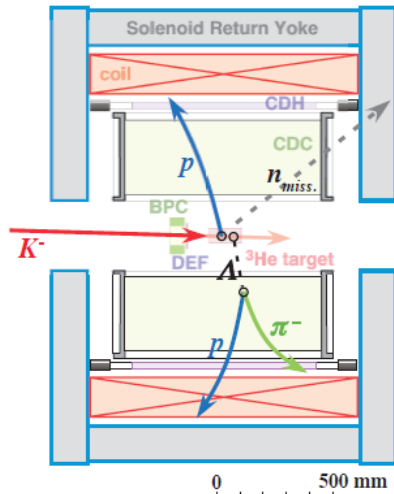
– Δd in $\text{Ni}+\text{Ni}$

• **Exclusive measurement using a simple reaction (in-flight & light nuclei) is crucial**



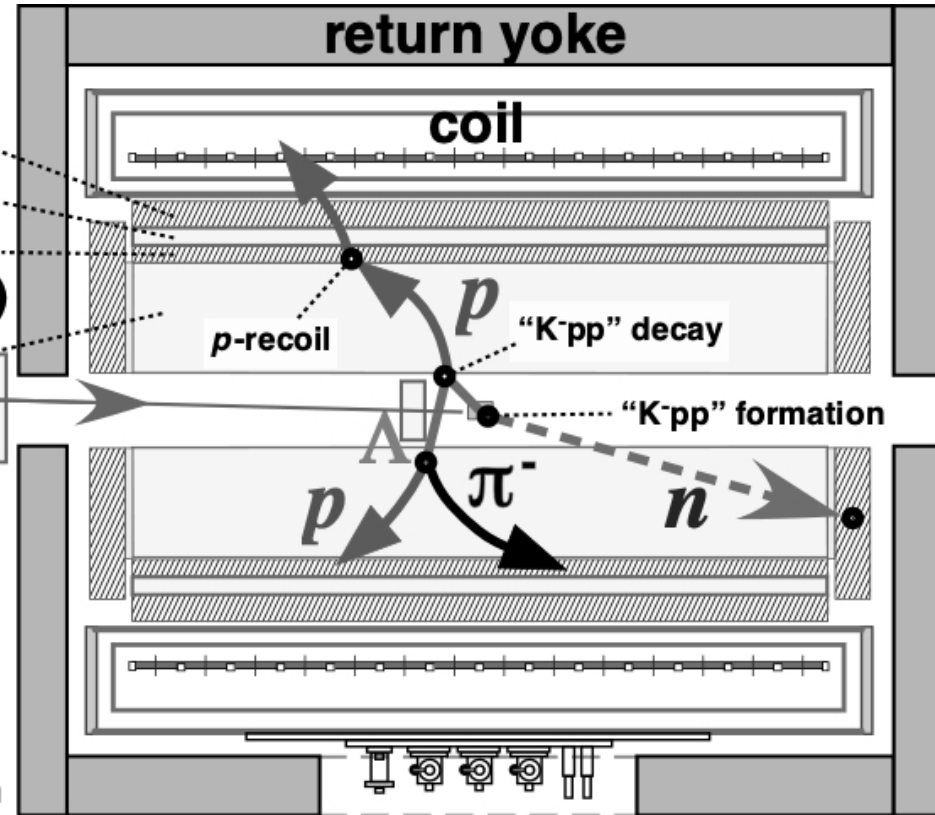
A New Cylindrical Detector System

present setup



n counter
outer-CDC
polarimeter (n counter)
inner-CDC

K^-



- SC Solenoid Magnet
- Cylindrical Drift Chamber
- Neutron Counter
- FWD/BWD Drift Chambers
- Vertex Fiber Tracker
- Electromagnetic Calorimeter (constructed in 2nd-stage)

Solid angle: $\sim \times 1.5$ ($\sim 90\%$)

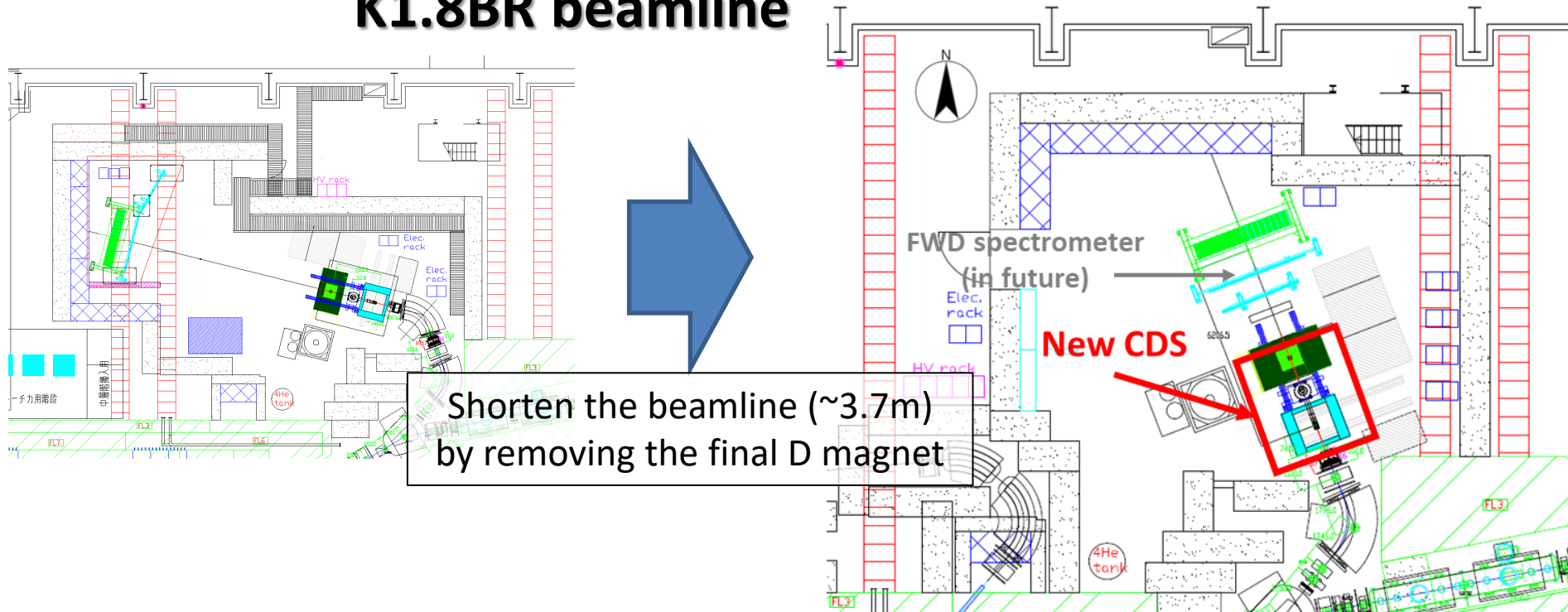
Neutron detection capability:

$\sim \times 10$

($\sim 1.5 \times 15\%$)

Improvement of Kaon Intensity

K1.8BR beamline



- We propose a new configuration of the beamline
 - K- yield is expected to increase by ~ 1.4 times @ 1.0 GeV/c

Expected Yield of $\bar{K}NNN$

$$N = \sigma \times N_{beam} \times N_{target} \times \epsilon,$$

$$\epsilon = \epsilon_{DAQ} \times \epsilon_{trigger} \times \epsilon_{beam} \times \epsilon_{fiducial} \times \Omega_{CDS} \times \epsilon_{CDS},$$

- We assume the K^-ppn cross section of

$$\sigma(K^-ppn) \cdot Br(\Lambda d) \sim 10 \mu b$$

$$\sigma(K^-ppn) \cdot Br(\Lambda pn) \sim 10 \mu b$$

◆ The same CS of “K-pp” $\rightarrow \Lambda p$ in E15

◆ As for Λd decay, we refer to the absorption of stopped K^- on ${}^4\text{He}$
 \rightarrow decay fraction to $\Sigma^-pd : \Sigma^-ppn \sim 1 : 1$

absorption of stopped K^- on ${}^4\text{He}$

Reaction	Events/(stopping K^-) (%)
$K^-He^4 \rightarrow \Sigma^+\pi^-H^3$	9.3 ± 2.3
$\rightarrow \Sigma^+\pi^-dn$	1.9 ± 0.7
$\rightarrow \Sigma^+\pi^-ppn$	1.6 ± 0.6
$\rightarrow \Sigma^+\pi^0nnn$	3.2 ± 1.0
$\rightarrow \Sigma^+nnn$	1.0 ± 0.4
Total $\Sigma^+ = (17.0 \pm 2.7)\%$	
$K^-He^4 \rightarrow \Sigma^-\pi^+H^3$	4.2 ± 1.2
$\rightarrow \Sigma^-\pi^+dn$	1.6 ± 0.6
$\rightarrow \Sigma^-\pi^+ppn$	1.4 ± 0.5
$\rightarrow \Sigma^-\pi^0He^3$	1.0 ± 0.5
$\rightarrow \Sigma^-\pi^0pd$	1.0 ± 0.5
$\rightarrow \Sigma^-\pi^0ppn$	1.0 ± 0.4
$\rightarrow \Sigma^-pd$	1.6 ± 0.6
$\rightarrow \Sigma^-ppn$	2.0 ± 0.7
Total $\Sigma^- = (13.8 \pm 1.8)\%$	
$K^-He^4 \rightarrow \pi^-\Lambda He^3$	11.2 ± 2.7
$\rightarrow \pi^-\Lambda pd$	10.9 ± 2.6
$\rightarrow \pi^-\Lambda ppn$	9.5 ± 2.4
$\rightarrow \pi^-\Sigma^0 He^3$	0.9 ± 0.6
$\rightarrow \pi^-\Sigma^0 (pd, ppn)$	0.3 ± 0.3
$\rightarrow \pi^0 \Lambda (\Sigma^0) (ppn)$	22.5 ± 4.2
$\rightarrow \Lambda (\Sigma^0) (ppn)$	11.7 ± 2.4
$\rightarrow \pi^+\Lambda (\Sigma^0)nnn$	2.1 ± 0.7
Total $\Lambda (\Sigma^0) = (69.2 \pm 6.6)\%$	
Total $= \Lambda + \Sigma = (100_{-7}^{+0})\%$	

Expected Yield of $\bar{K}NNN$

$$N = \sigma \times N_{beam} \times N_{target} \times \epsilon,$$

$$\epsilon = \epsilon_{DAQ} \times \epsilon_{trigger} \times \epsilon_{beam} \times \epsilon_{fiducial} \times \Omega_{CDS} \times \epsilon_{CDS},$$

- $N_{beam} = \mathbf{100 G K^-}$ on target
 - under the MR beam power of **90 kW** with **5.2 s** repetition cycle.
 - 3.2×10^5 K- on target / spill @ 1.0 GeV/c around 2024
 - **3 weeks** data taking (90% up-time)
- $N(K^-ppn \rightarrow \Lambda d) \sim \mathbf{2 \times 10^4}$
- $N(K^-ppn \rightarrow \Lambda pn) \sim \mathbf{3 \times 10^3}$
 - c.f. 1.7×10^3 “K-pp” $\rightarrow \Lambda p$ accumulated in E15-2nd (40 G K-)

	$\Lambda d / \Lambda pn$
$\sigma(K^-ppn)*Br$	10 μb
$N(K^- \text{ on target})$	100 G
$N(\text{target})$	2.65×10^{23}
$\epsilon(\text{DAQ})$	0.9
$\epsilon(\text{trigger})$	0.93
$\epsilon(\text{beam})$	0.55
$\Omega(\text{CDC})$	0.27 / 0.077
$\epsilon(\text{CDC})$	0.6 / 0.3
$N(K^-ppn)$	19 k / 2.8 k

* improved from E15

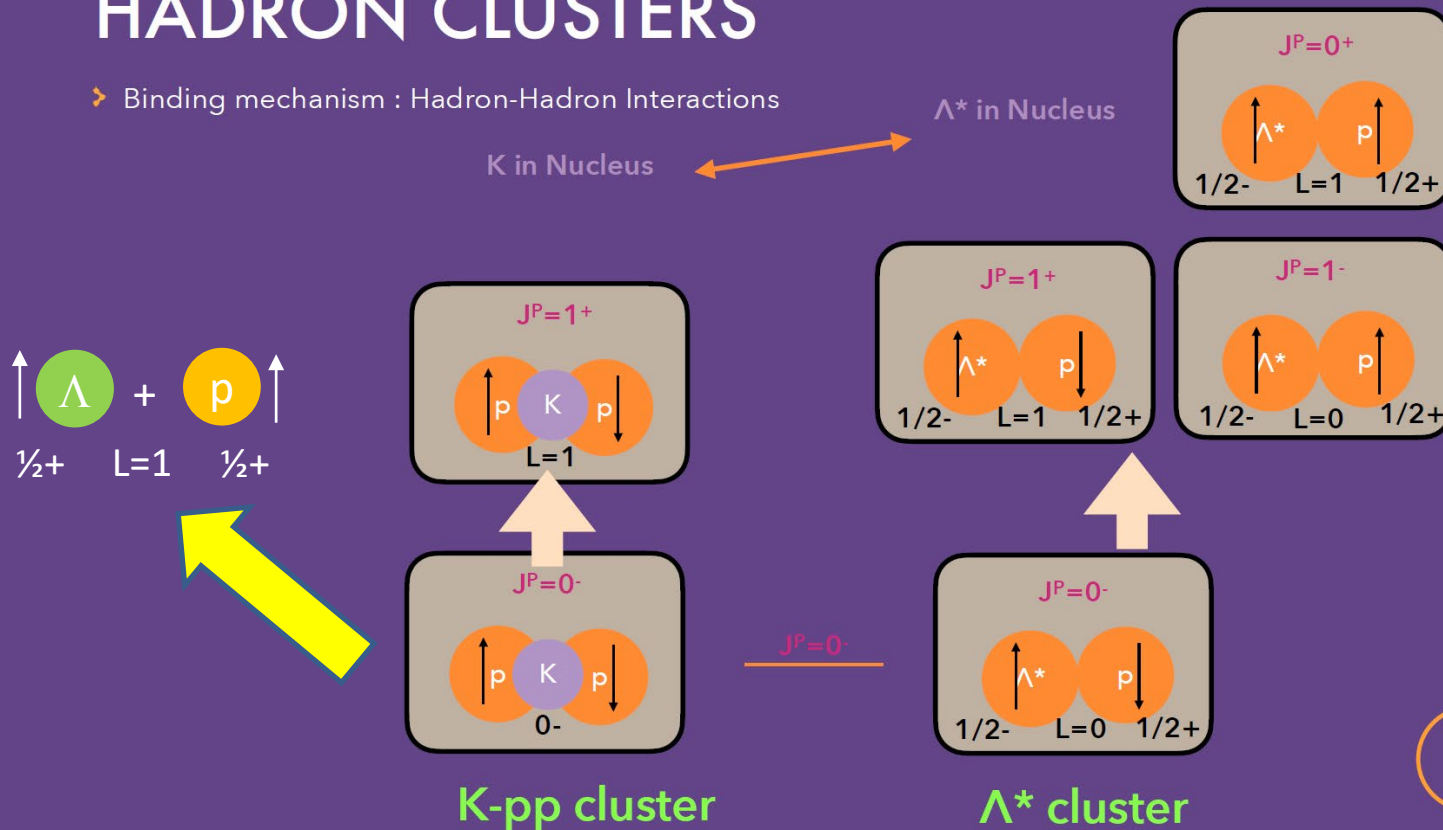
Spin-Parity of $\bar{K}NN$

5

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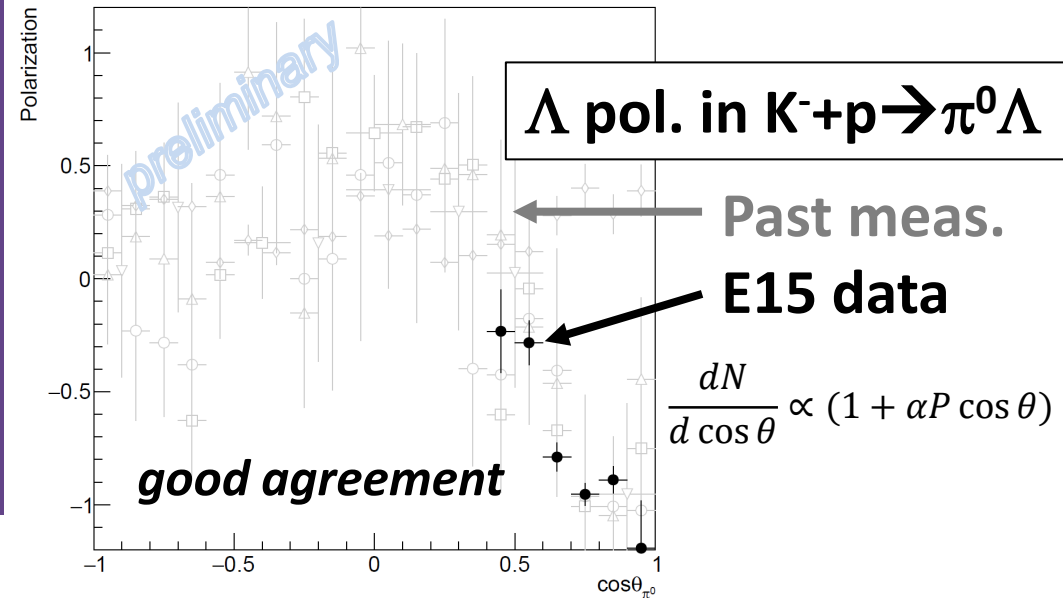
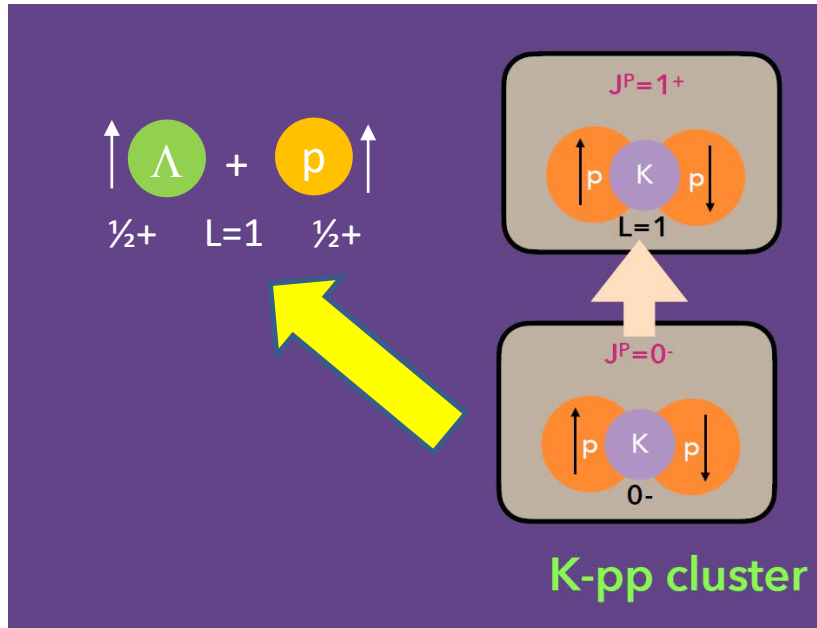
➤ Binding mechanism : Hadron-Hadron Interactions



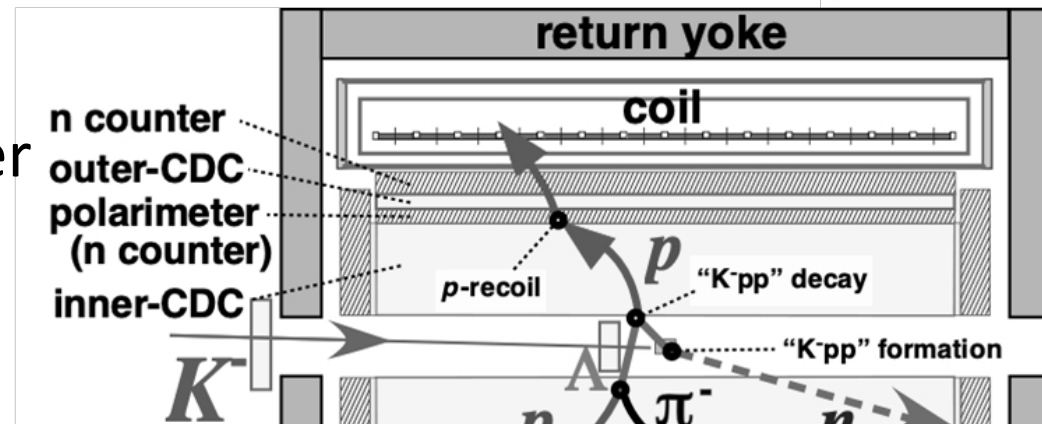
T.Nagae, "The 4th Symposium on Clustering as a window on the hierarchical structure of quantum systems", May 28th, 2020

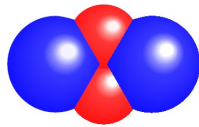
Spin-Parity of $\bar{K}NN$

- Λ polarization can be measured via π^-p decay



- Proton polarization is measured with a polarimeter
 - Tracking system
 - Plastic scintillator



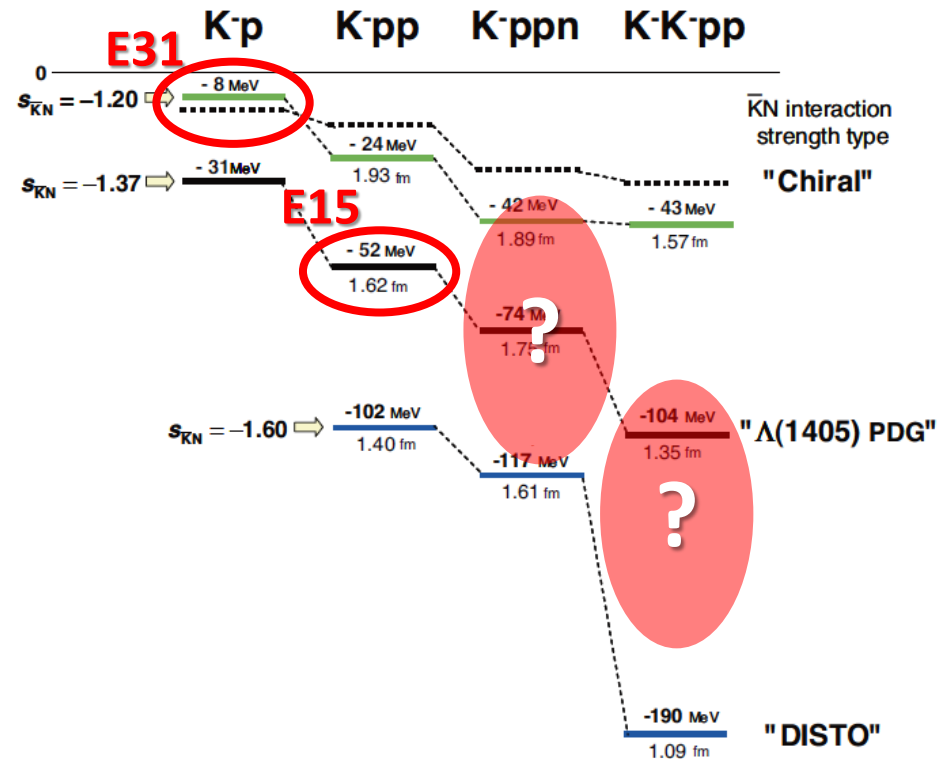
K⁻K⁻pp

Toward $\bar{K}\bar{K}NN$

- We also wish to access the $S = -2$ kaonic nuclei such as the theoretically predicted "K⁻K⁻pp" state

- ✓ as previously submitted Lol
- ✓ A good probe to the $\bar{K}N$ int.

- The $\bar{K}\bar{K}NN$ system could give us a chance to access much higher density than the $S = -1$ kaonic nuclei



Proc. Jpn. Acad., **B89** (2013) 418.



**The $\bar{K}\bar{K}NN$ production cross section would be quite small
 → roughly 1/1000 of that of the $\bar{K}NN$**