



International Conference on Exotic Atoms and Related
Topics and conference on Low Energy Antiprotons (EXA/
LEAP 2024)

Prospect of Hadronic-Molecule / Cluster with Strangeness

- via the detailed study of the kaonic nucleus -

- 28 / 08 / 2024 -

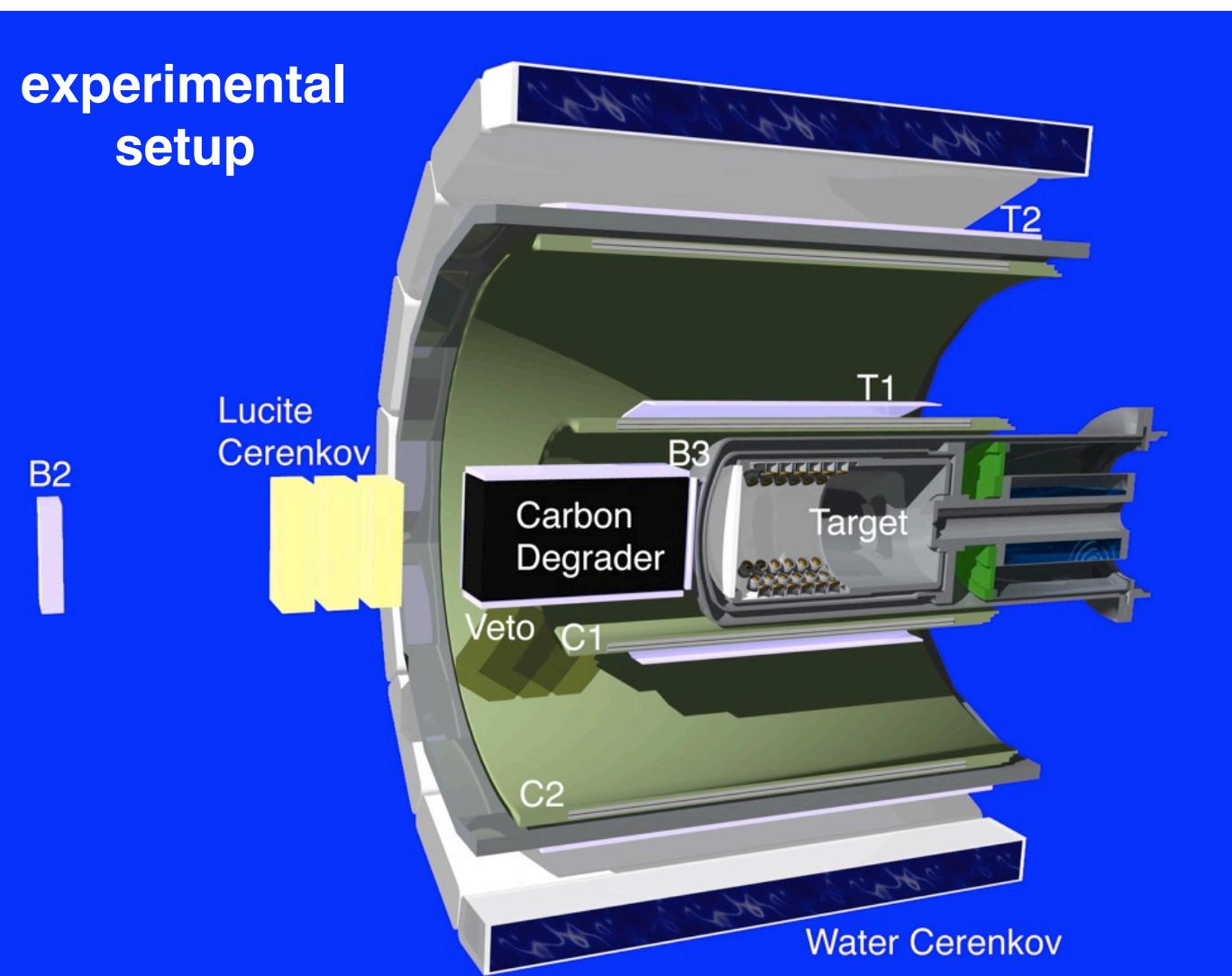
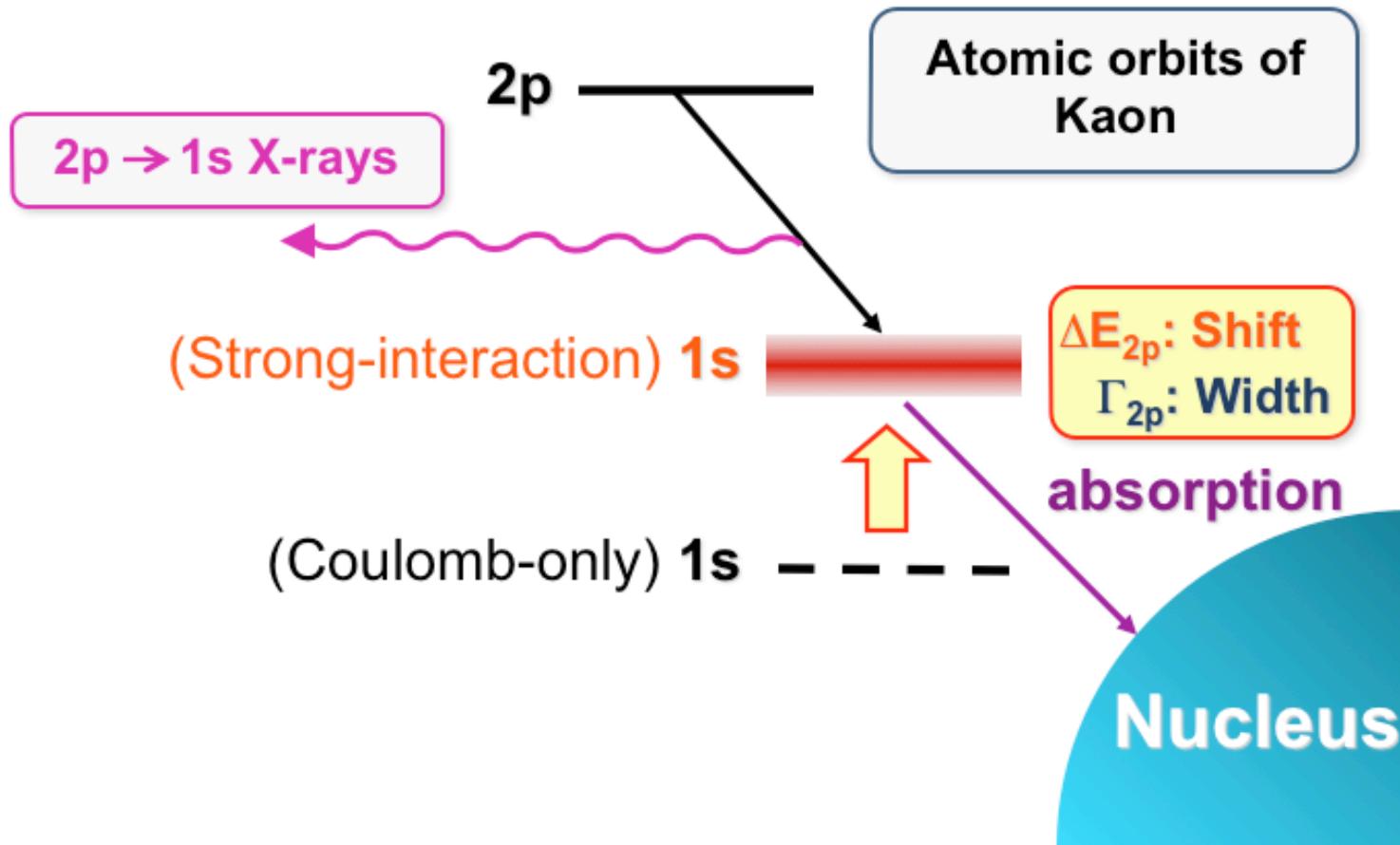
M. Iwasaki / RIKEN

$\bar{K}N$ interaction study via Kaonic atom

Succeeded in Kaonic Hydrogen x-ray Measurement

KpX: obtained X-ray spectrum 3

KpX experiment at KEK in 1997:



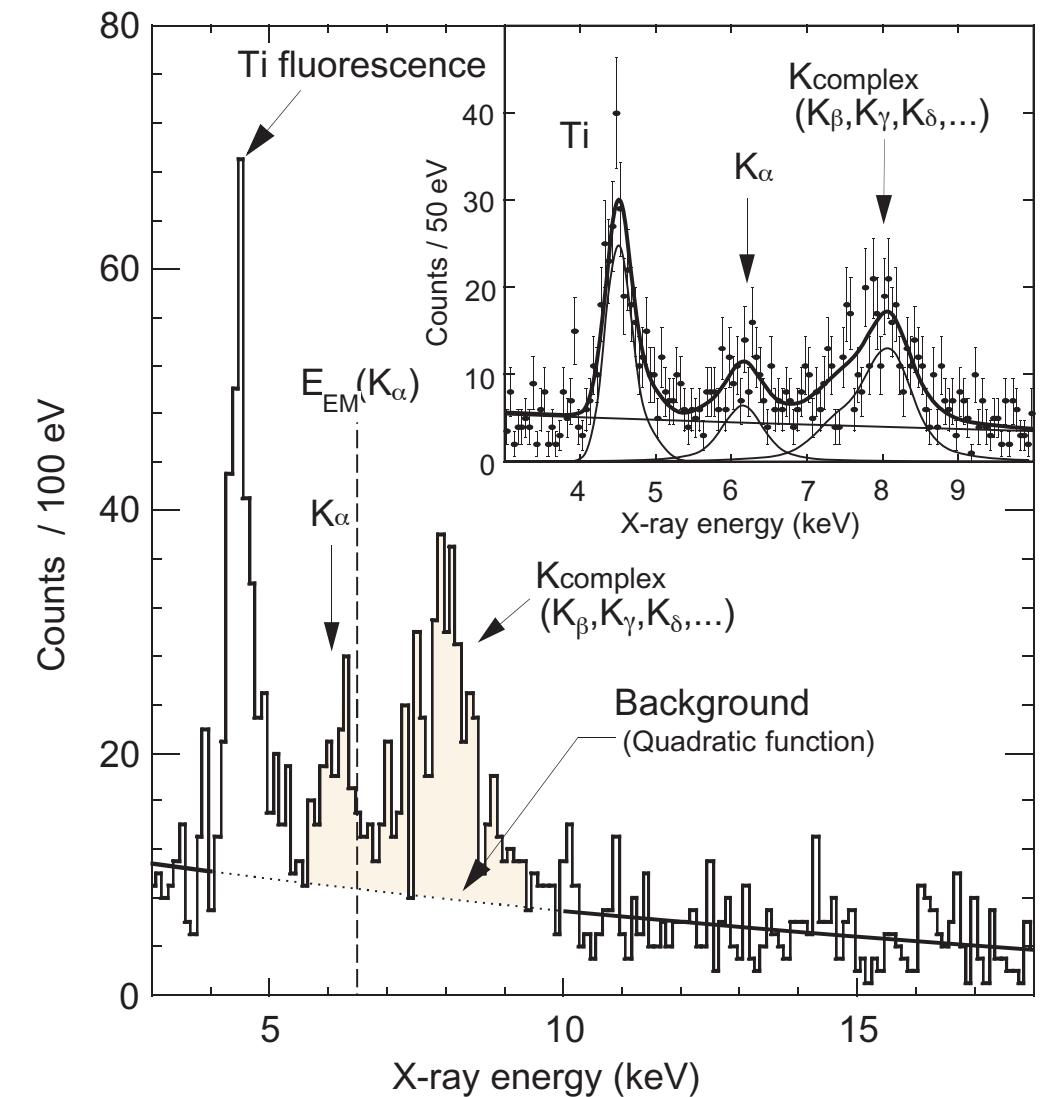
The European Physical Journal C
Volume 15 · Number 1-4 · 2000
THE $\Lambda(1405)$
Revised March 1998 by R.H. Dalitz, Oxford University

.....

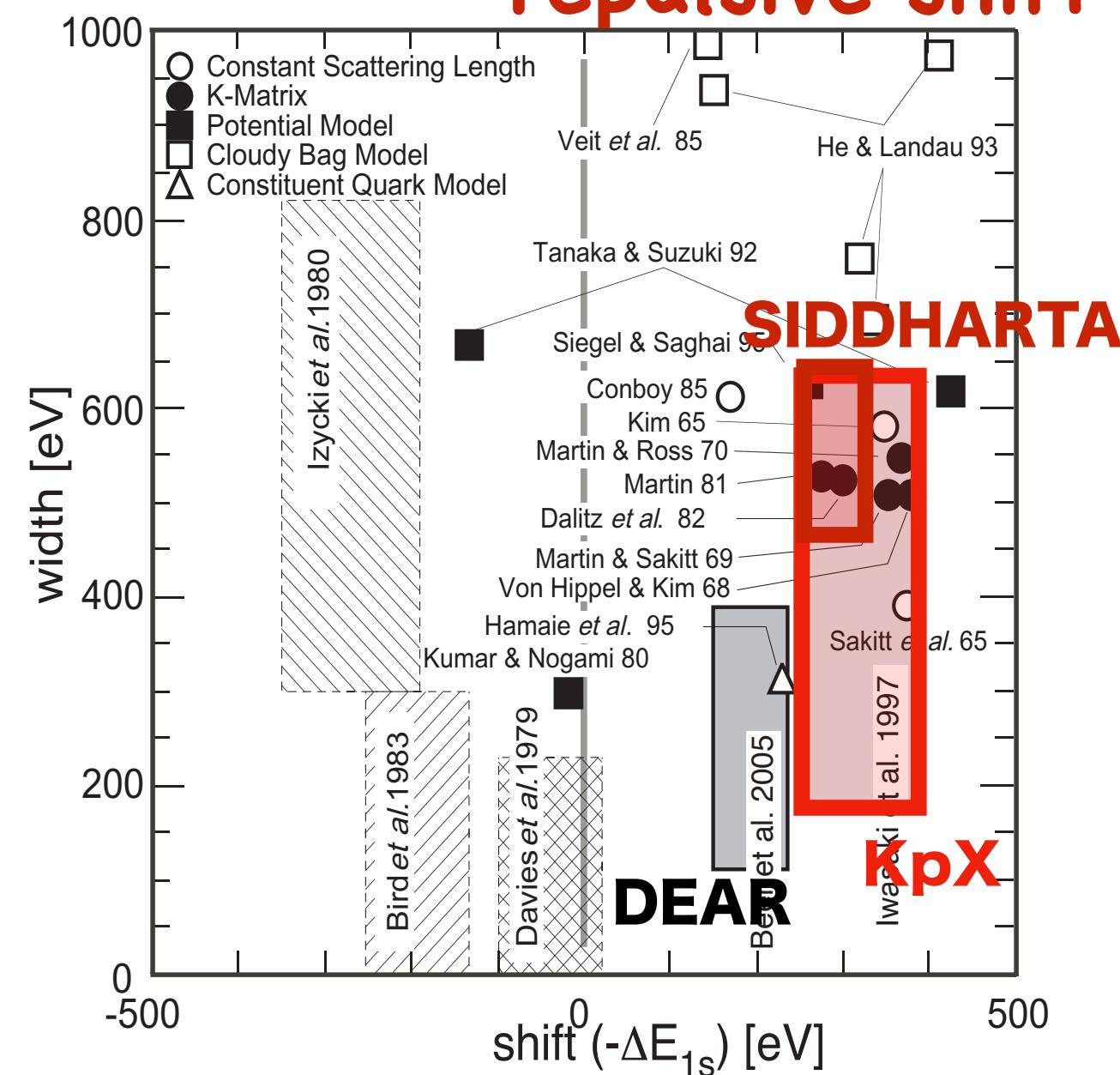
From the measurement of $2p - 1s$ x rays from kaonic-hydrogen, the energy-level shift ΔE and width Γ of its $1s$ state can give us two further constraints on the $(\bar{\Sigma}\pi, NK)$ system, at an energy roughly midway between those from the low-energy hydrogen bubble chamber studies and those from qR($\Sigma\pi$) observations below pK^- threshold. IWASAKI 97 have reported the first convincing observation of this x ray, with a good initial estimate:

$$\Delta E - i\Gamma/2 = (-323 \pm 63 \pm 11) - i(204 \pm 104 \pm 50) \text{ eV. (2)}$$

the errors here encompass about half of the predictions made following various analyses and/or models for the in-flight K^-p and sub-threshold qR($\Sigma\pi$) data. Better measurements will be needed to discriminate between the analyses and predictions., perhaps from the DAΦNE storage ring at Frascati, information vital for our quantitative understanding of the $(\Sigma\pi, NK)$ system in this region.



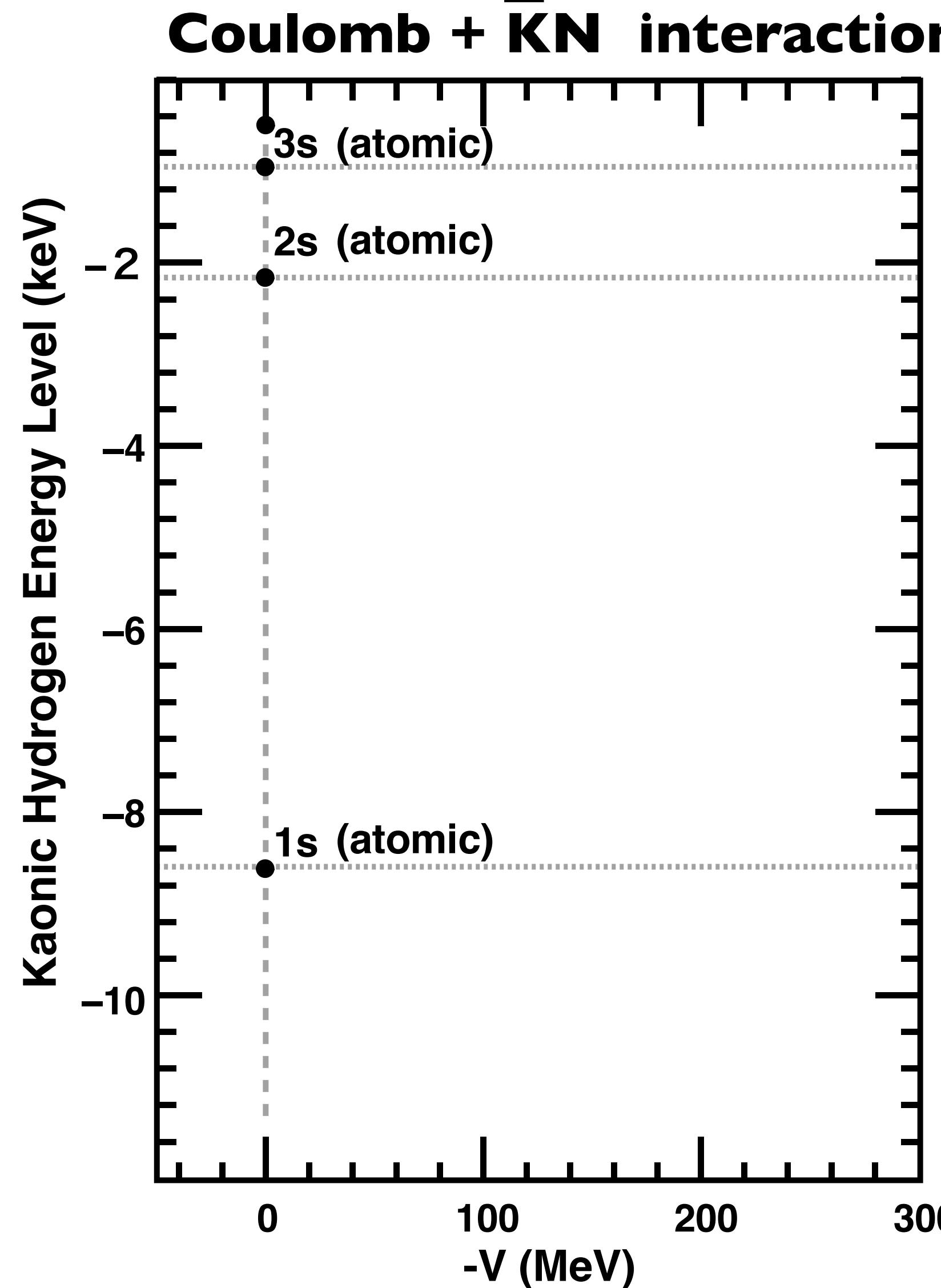
repulsive shift



What's next in physics?

Observed Shift was REPULSIVE!

Is $\bar{K}N$ interaction repulsive?



Let's study how atomic level shifts depending on $\bar{K}N$ interaction.

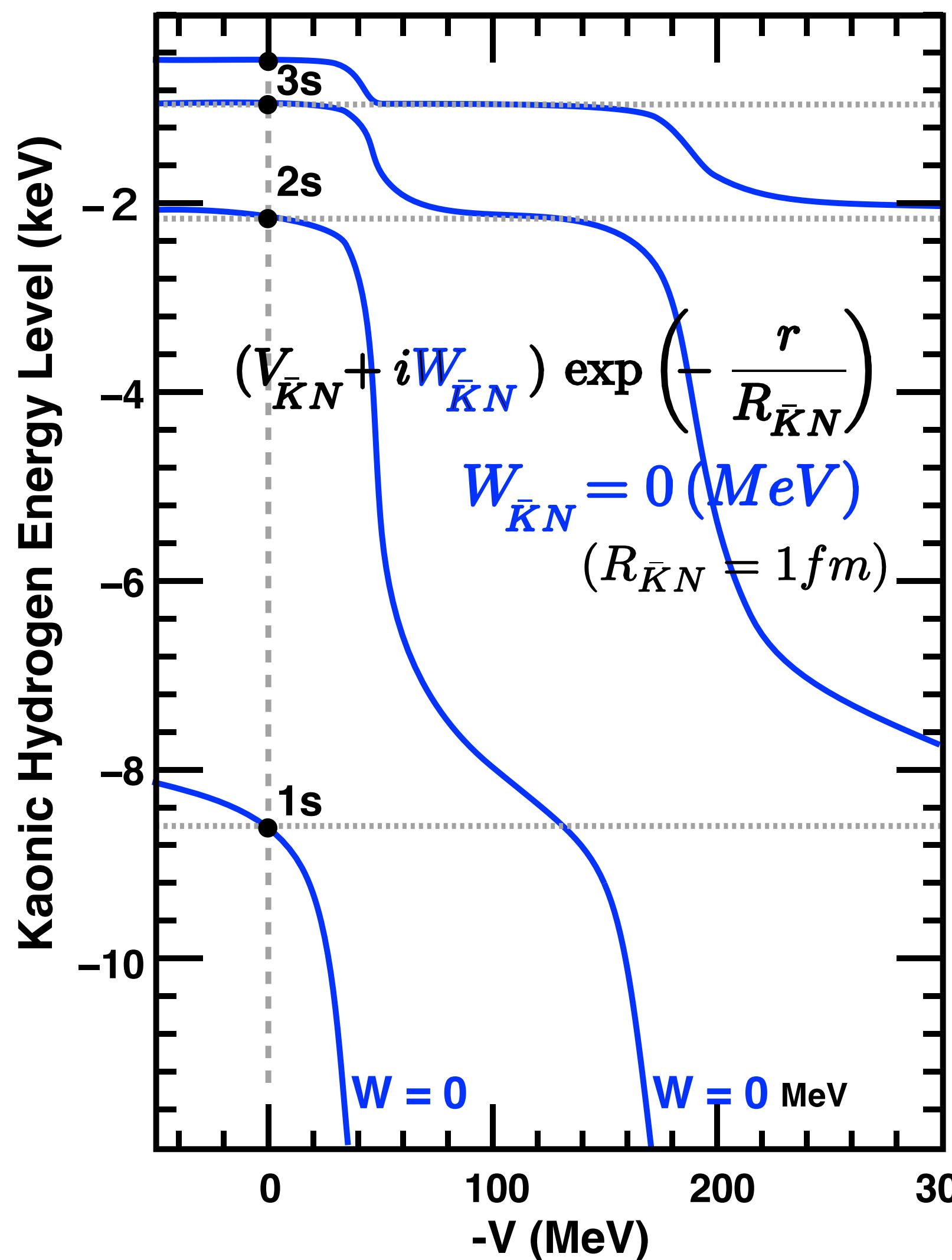
Without $\bar{K}N$ interaction

$(V_{\bar{K}N} = 0, W_{\bar{K}N} = 0)$, Coulomb potential forms $K^- p$ atomic levels.

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Coulomb + $\bar{K}N$ interaction



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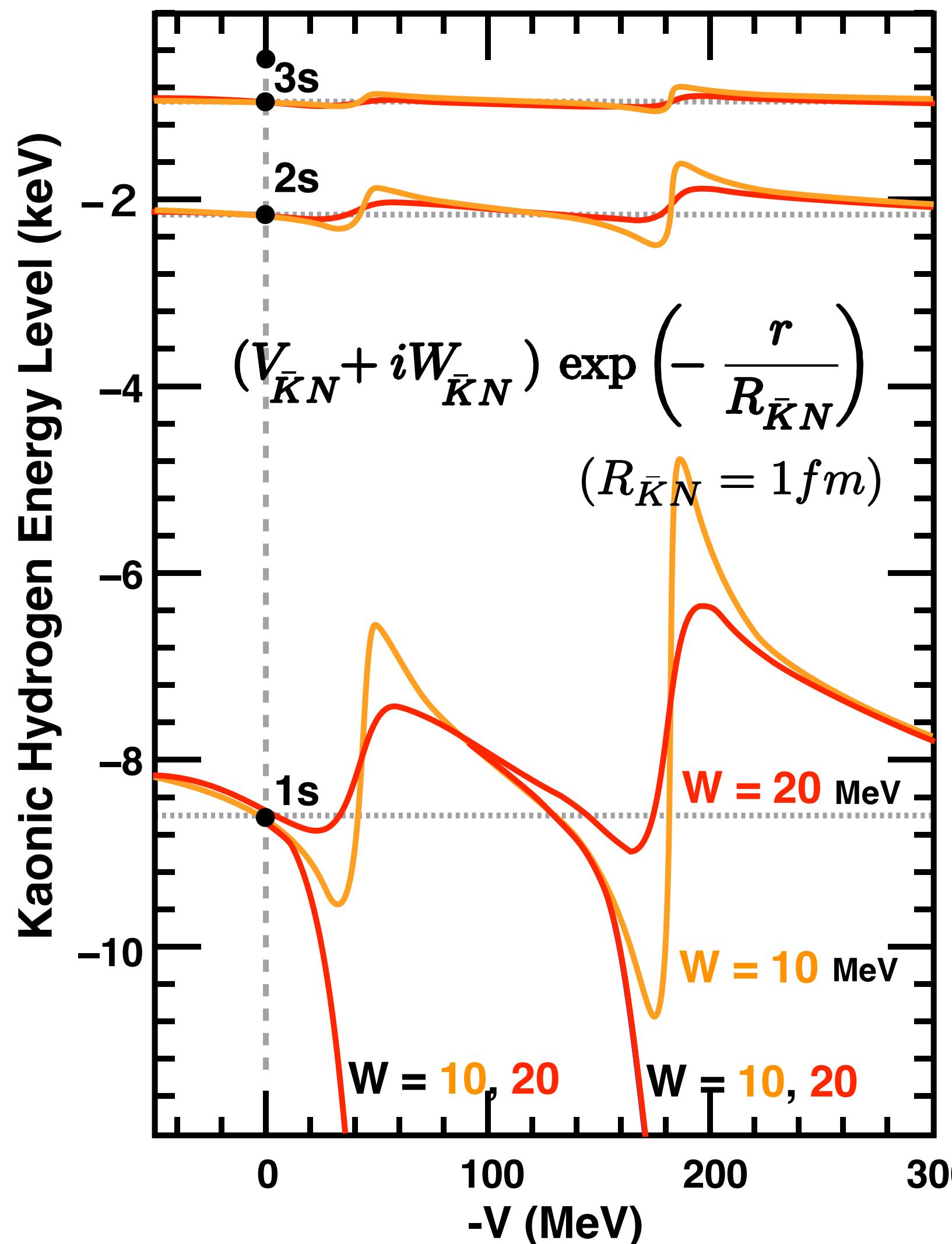
$(V_{\bar{K}N} = 0, W_{\bar{K}N} = 0)$, Coulomb potential forms $K^- p$ atomic levels.

Each atomic level shifts downward as a function of $-V_{\bar{K}N}$ (at $W_{\bar{K}N} = 0$) in a step-like function, and atomic 1s change its nature to be a nuclear ground state.

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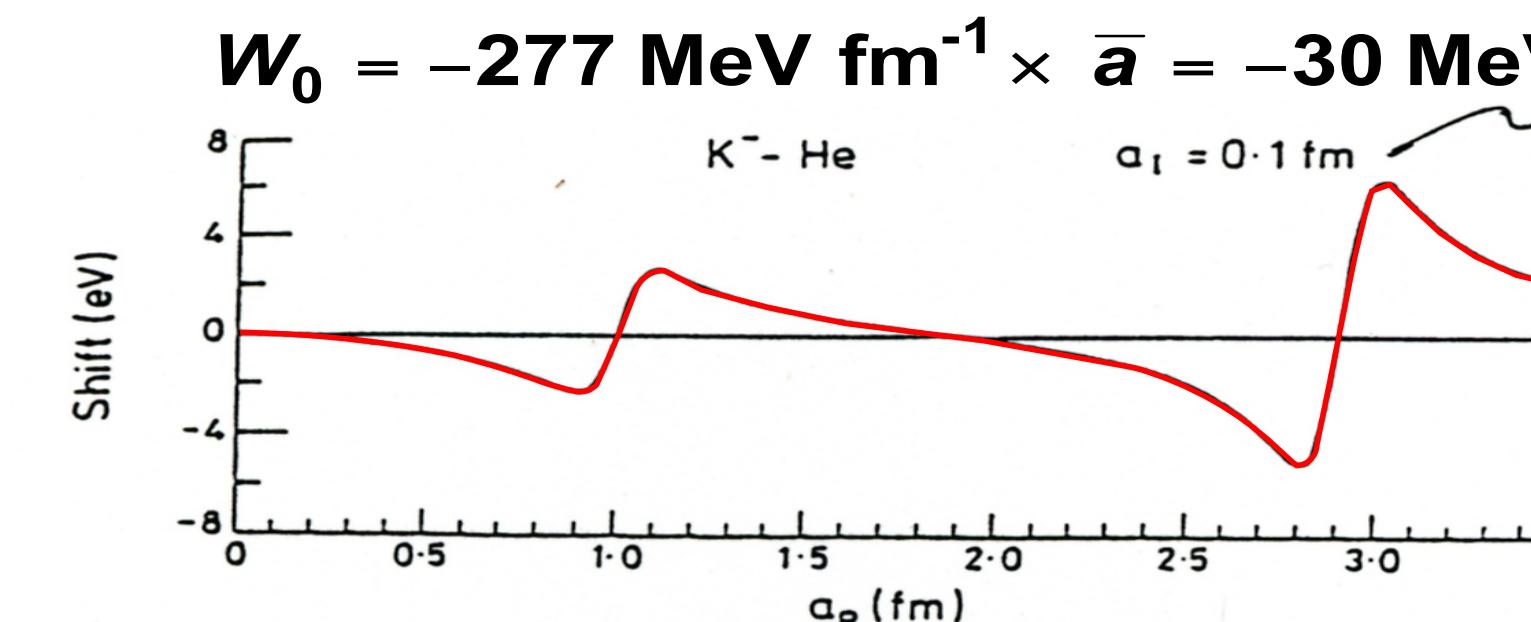
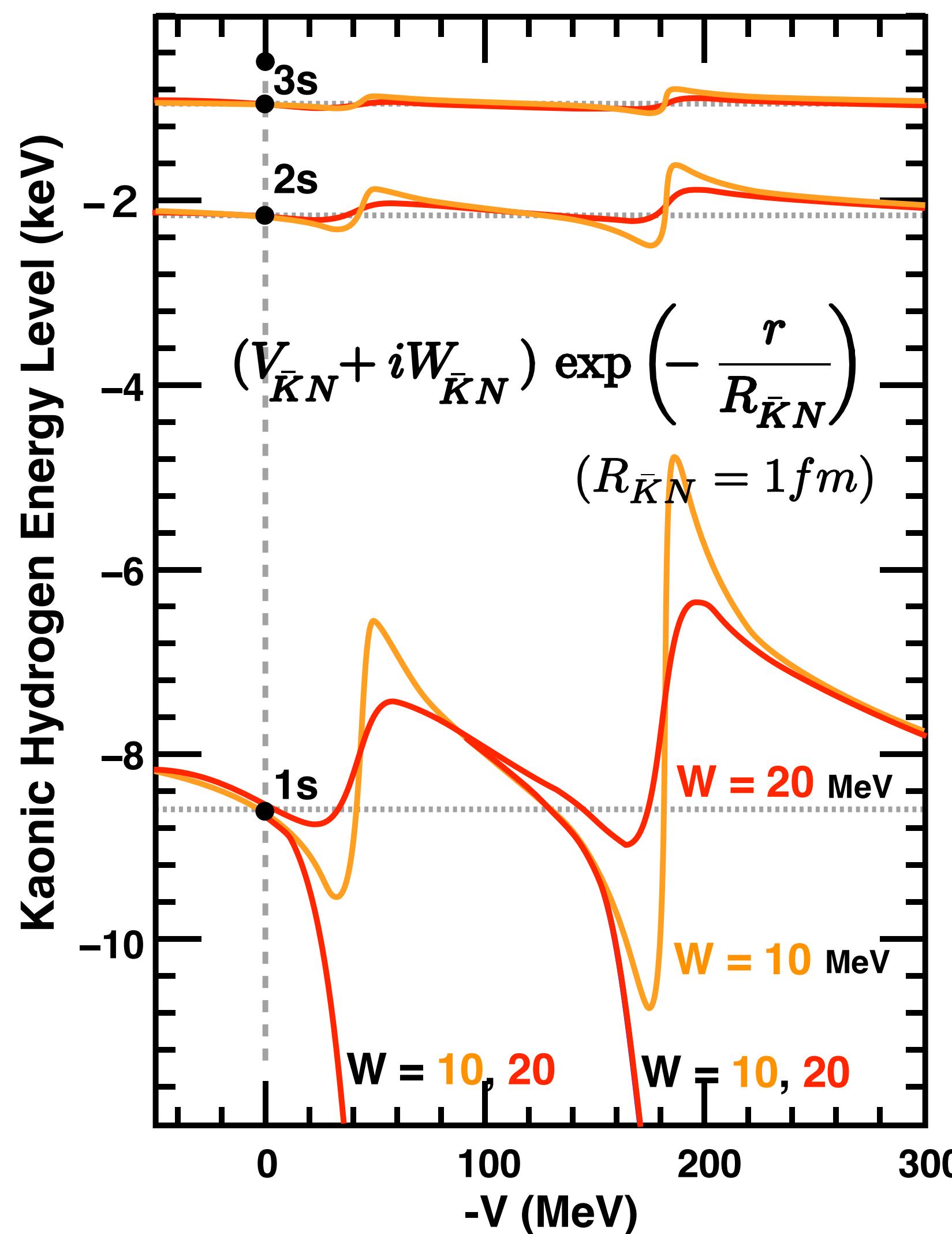
When switching on the absorption part $W_{\bar{K}N}$, a level crossing happens and connects to another level.

As a result, a wiggling pattern appears in energy levels (red lines), and the nuclear bound state is branched and separated from the atomic ground state.

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Theoretical calculations for realistic $\bar{K}N$ interaction

R. Seki, Phys. Rev. C5 (1972) 1196

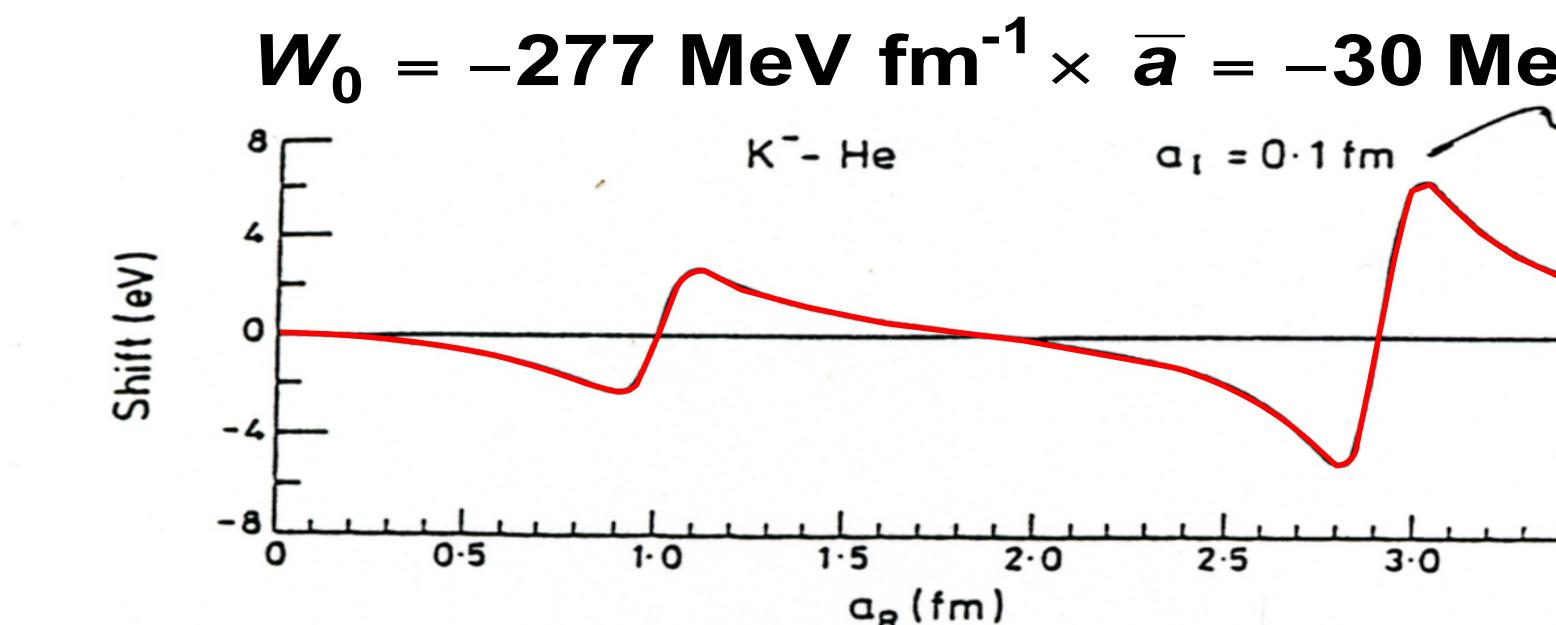
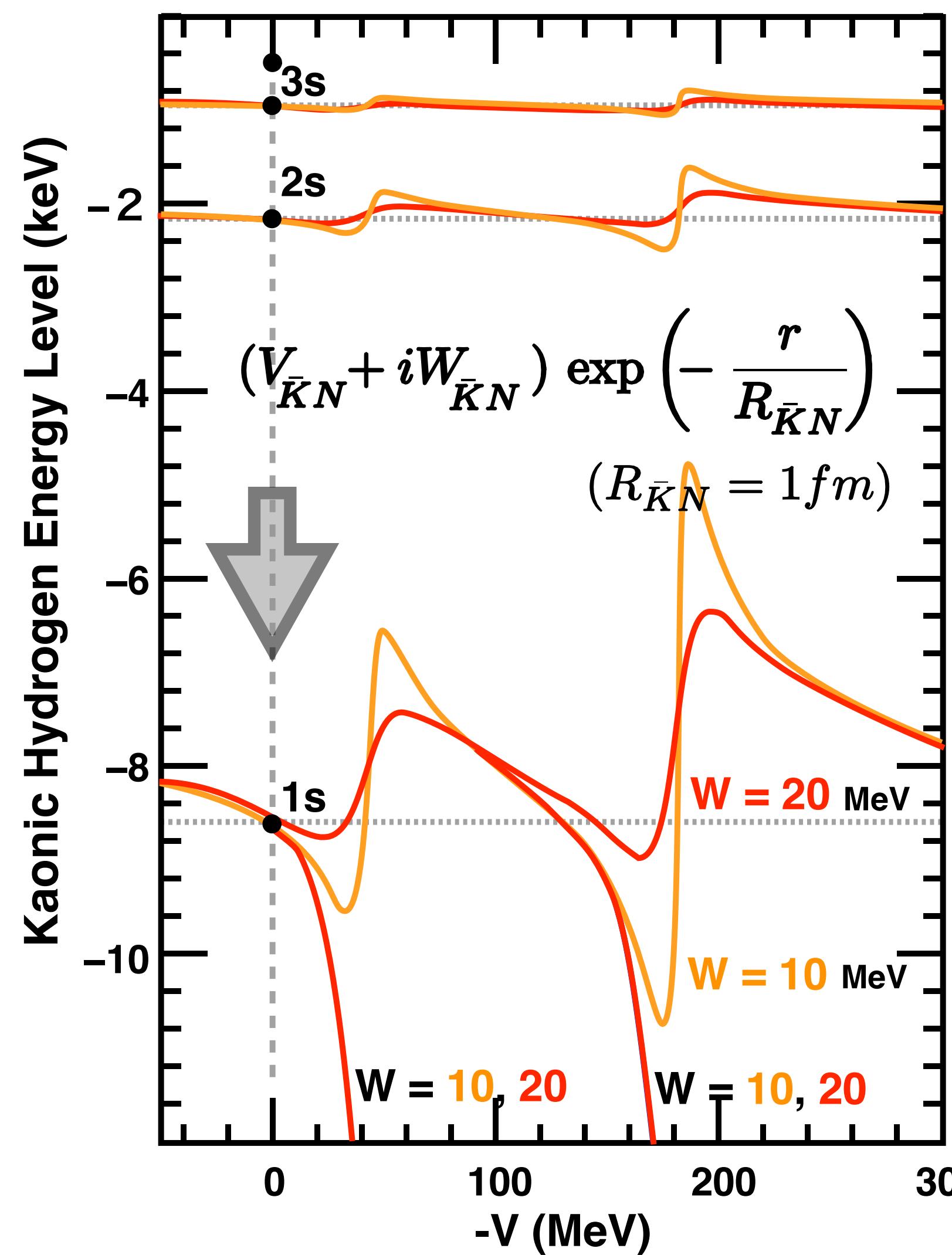
S. Baird et al., Nucl. Phys. A392 (1983) 297

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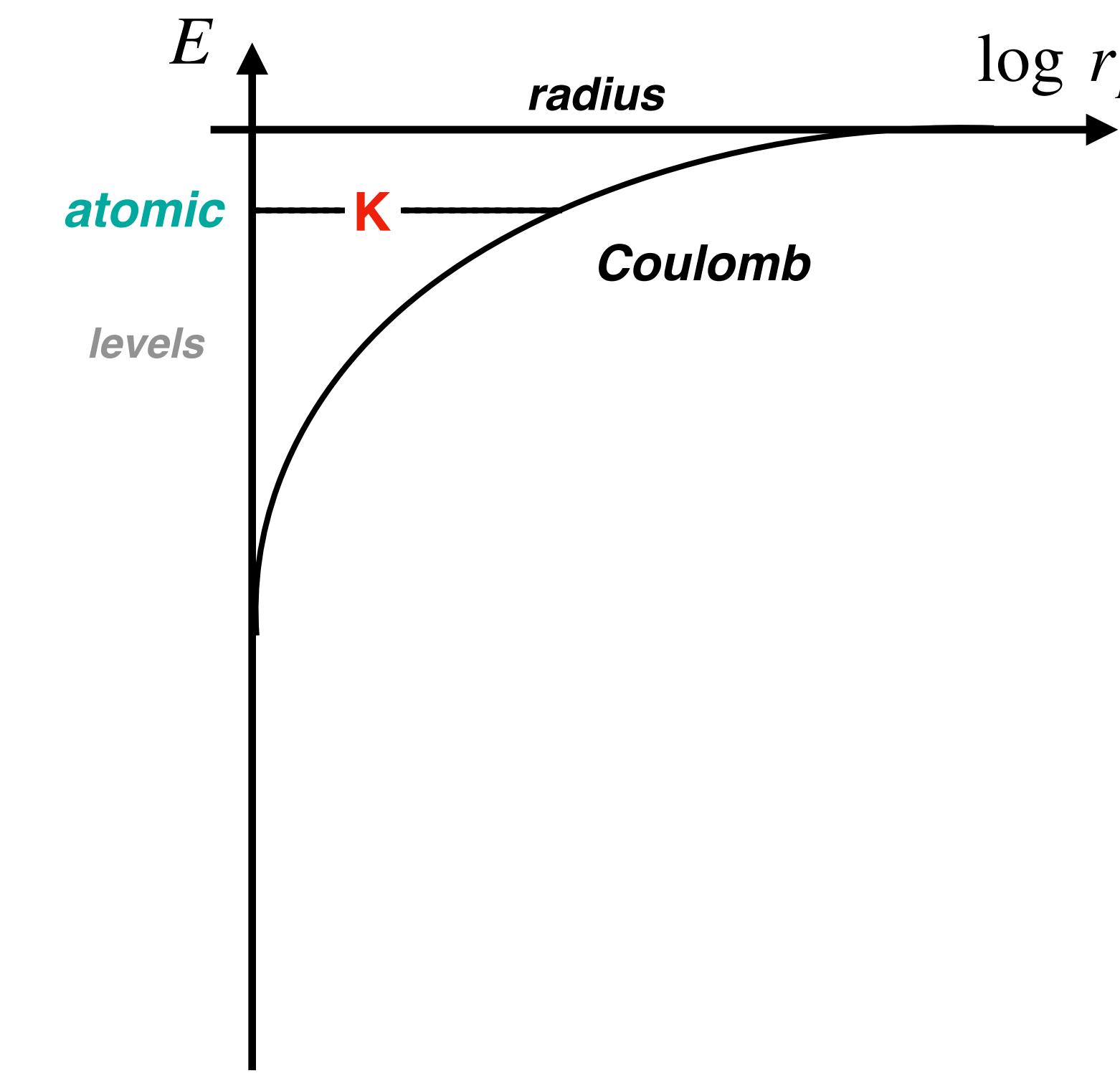


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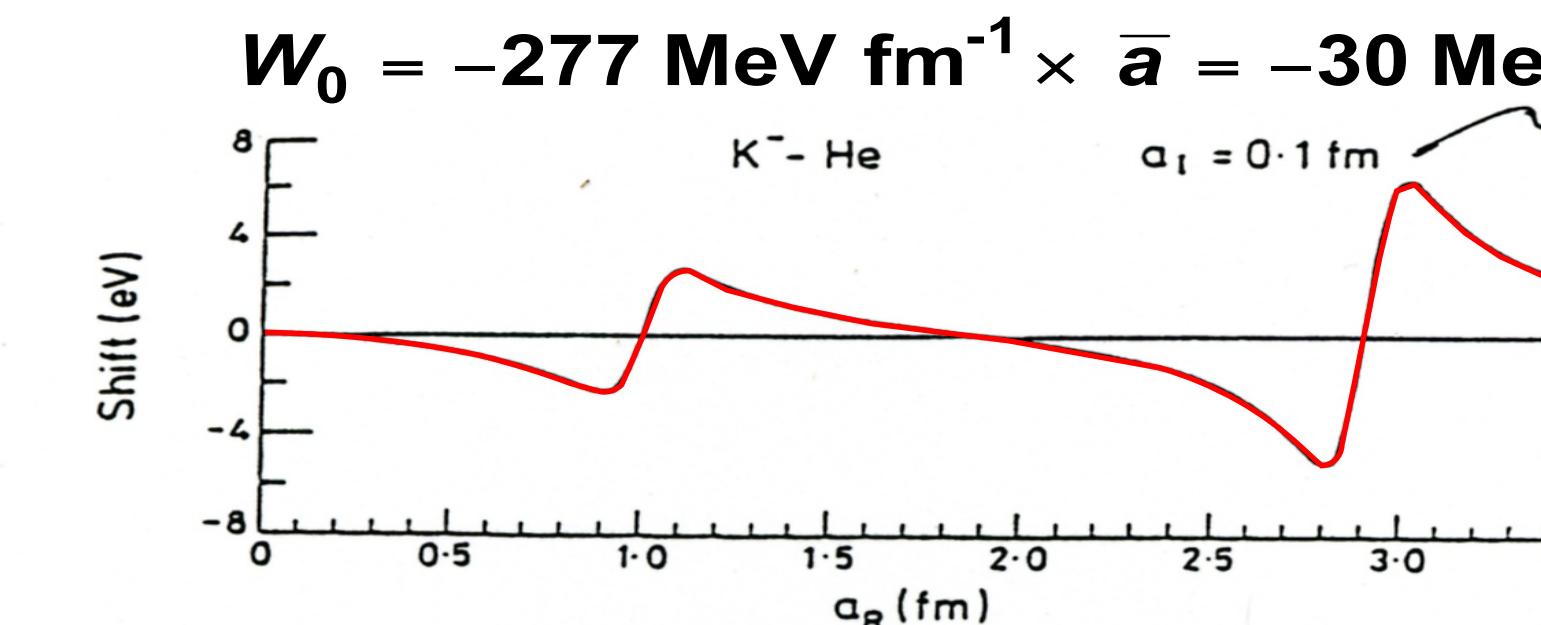
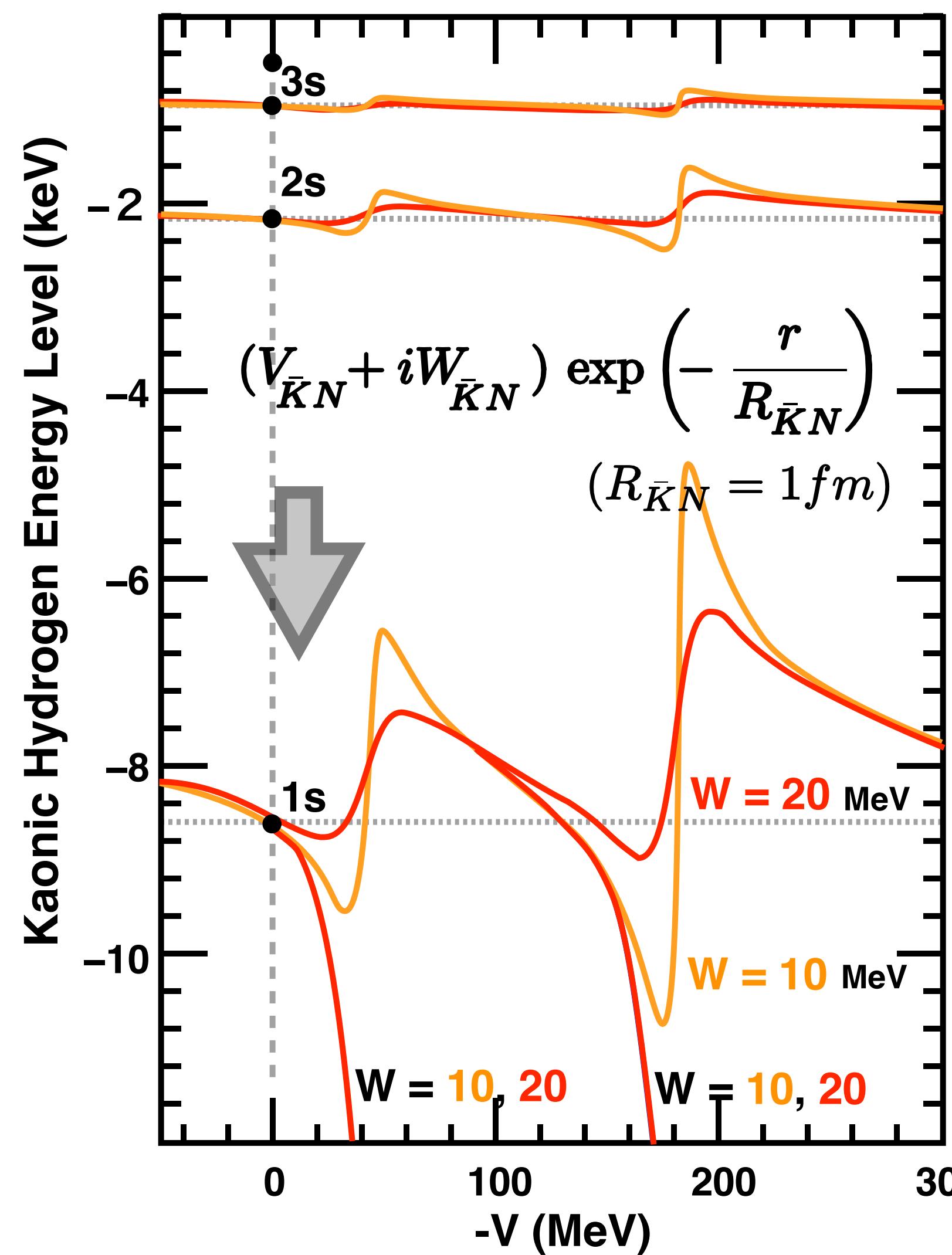


Illustrative Animation

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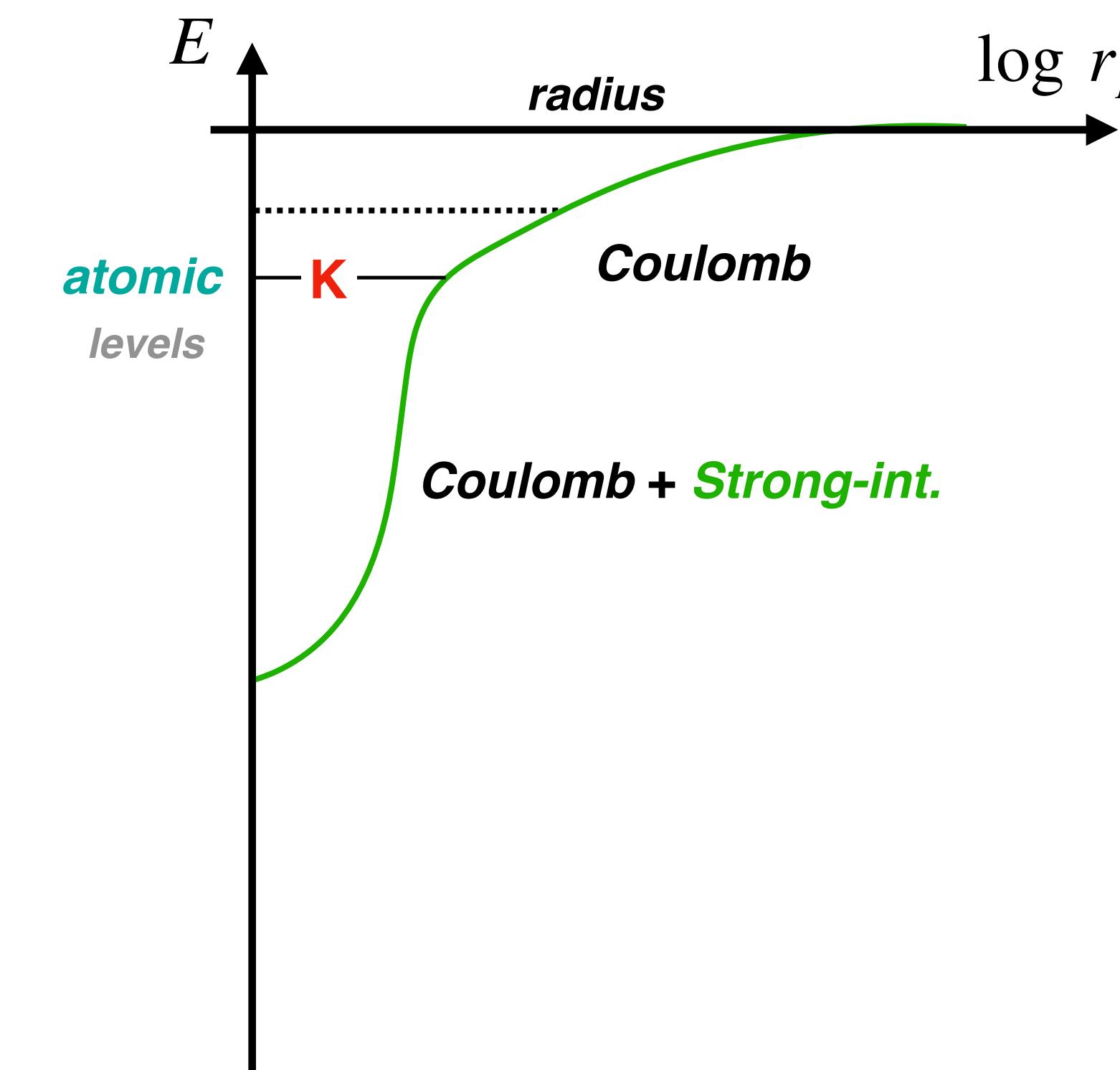


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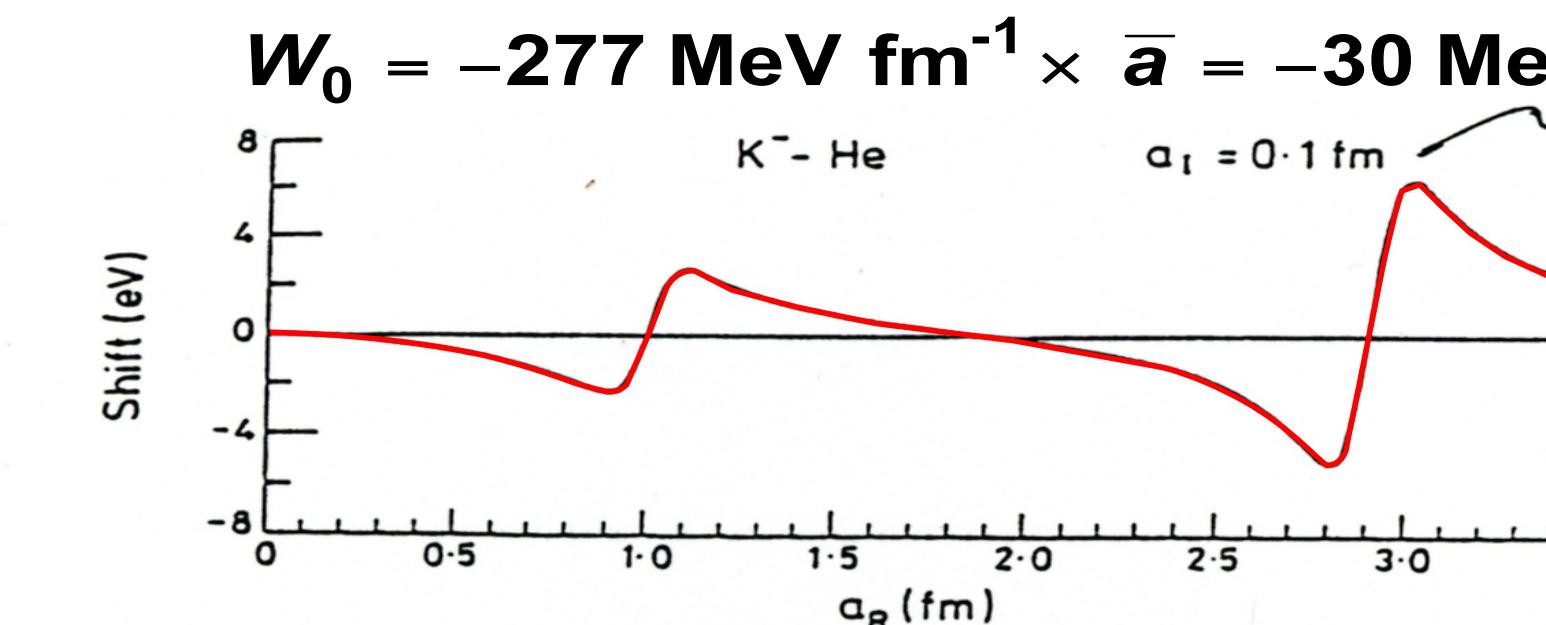
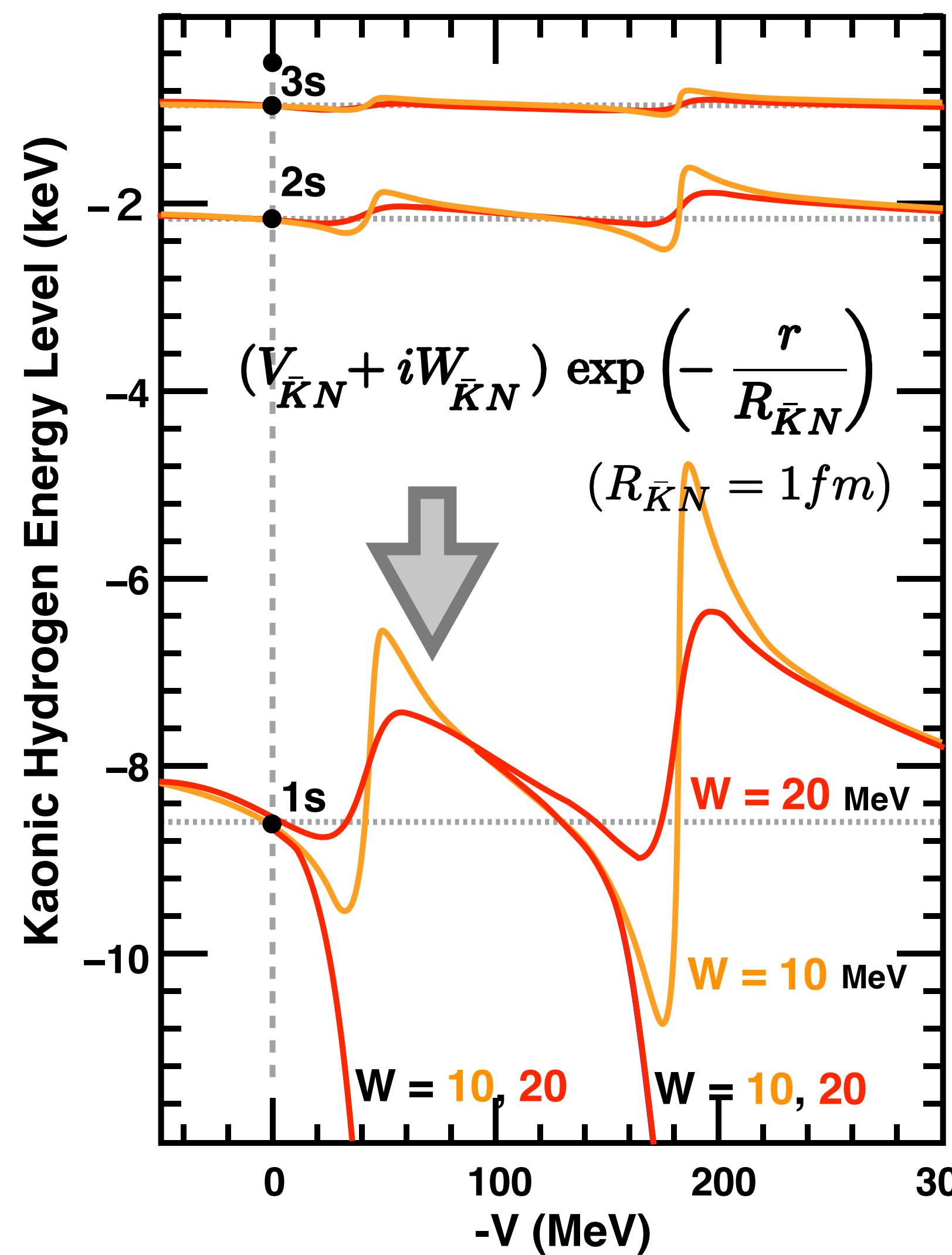


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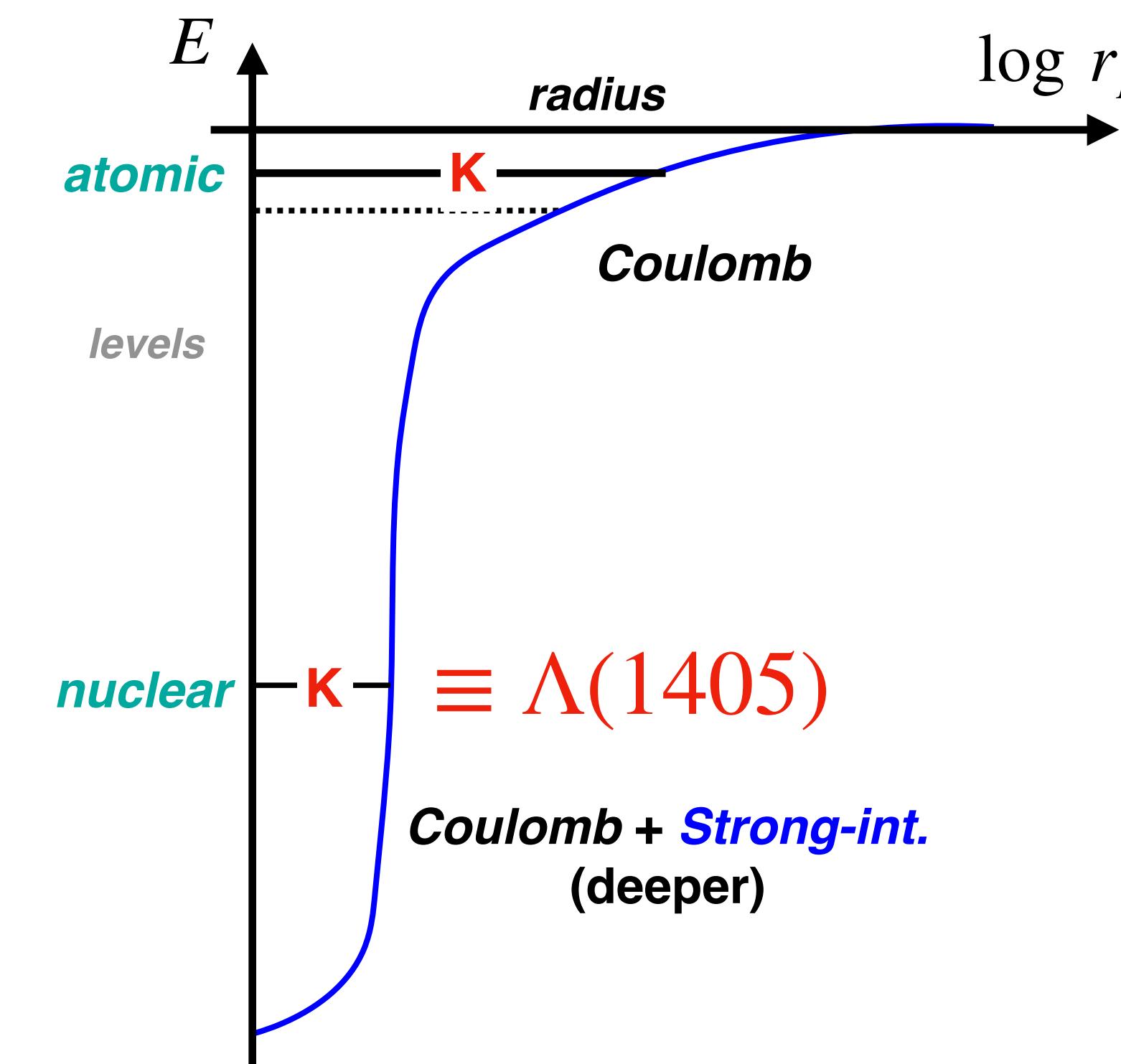


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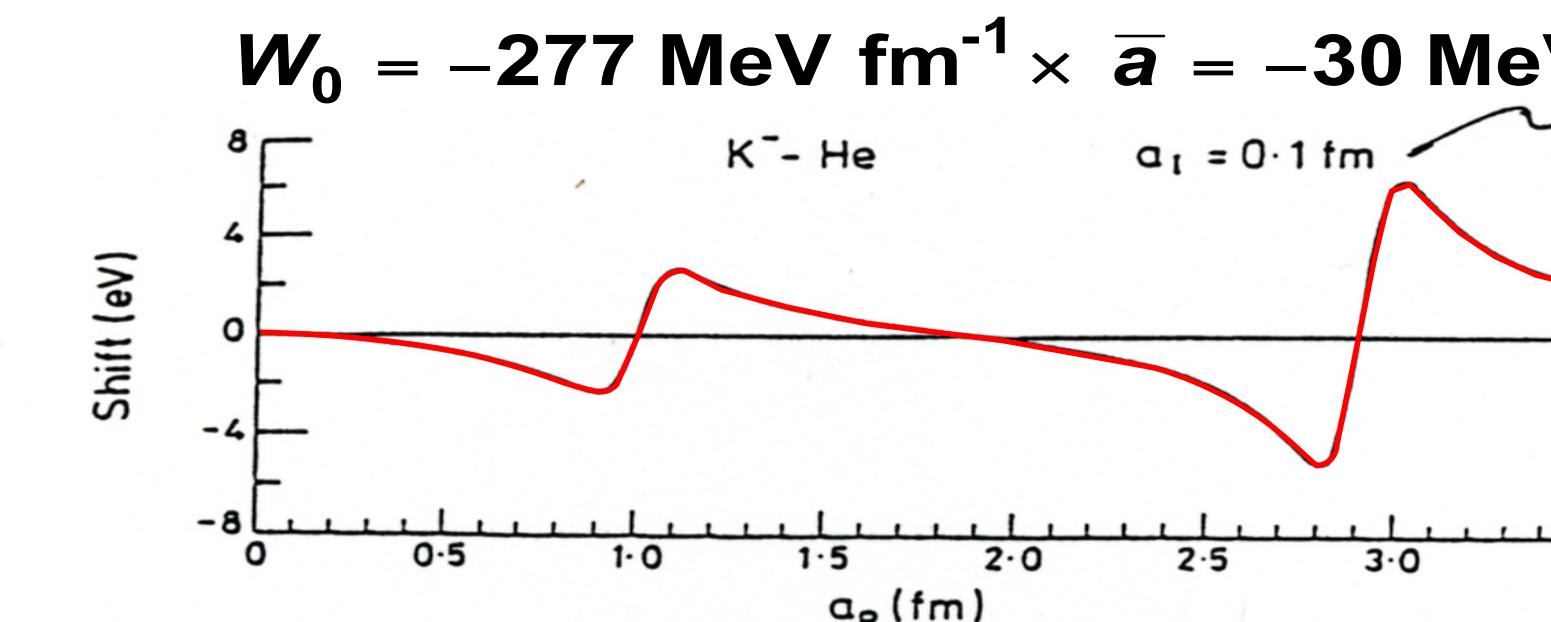
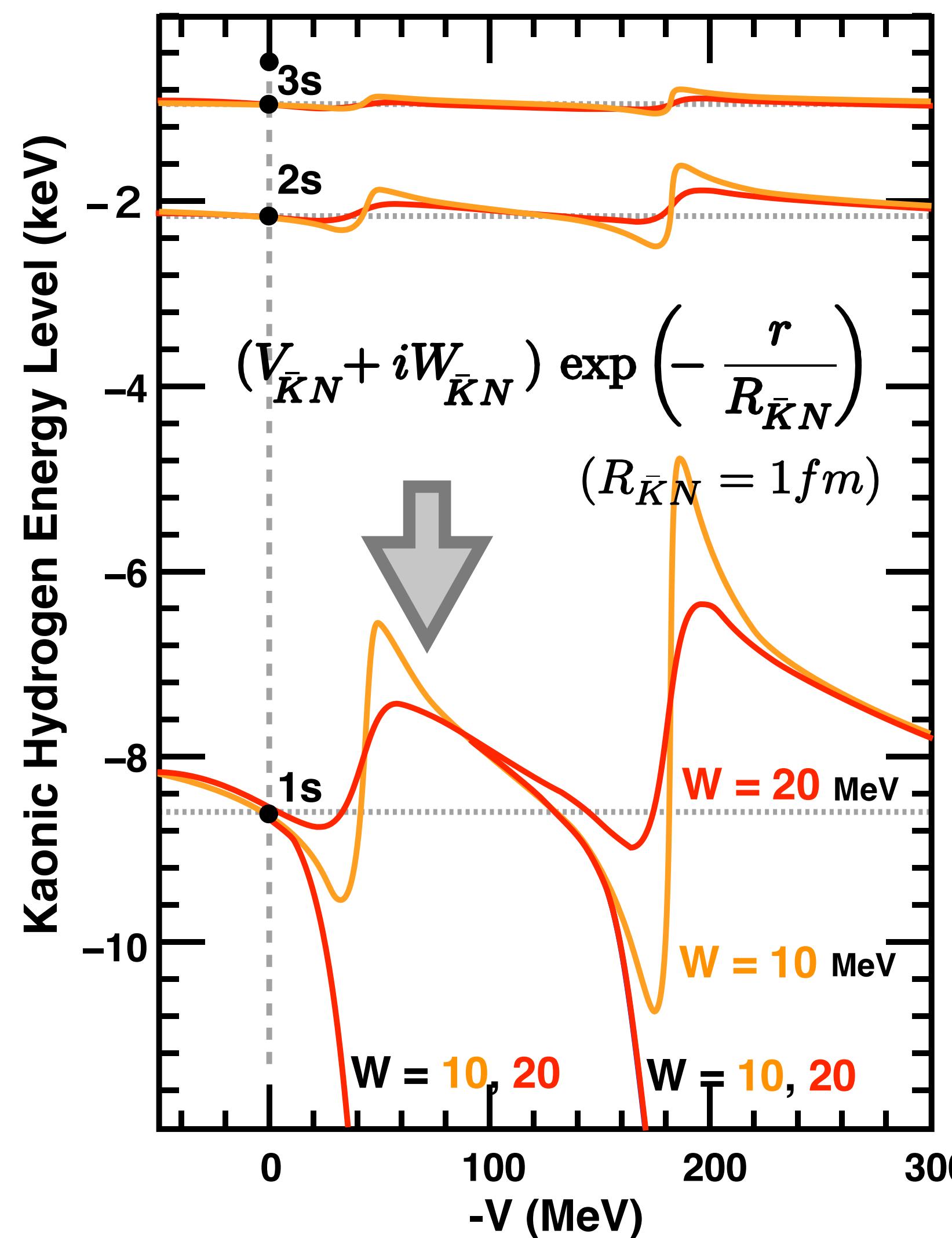


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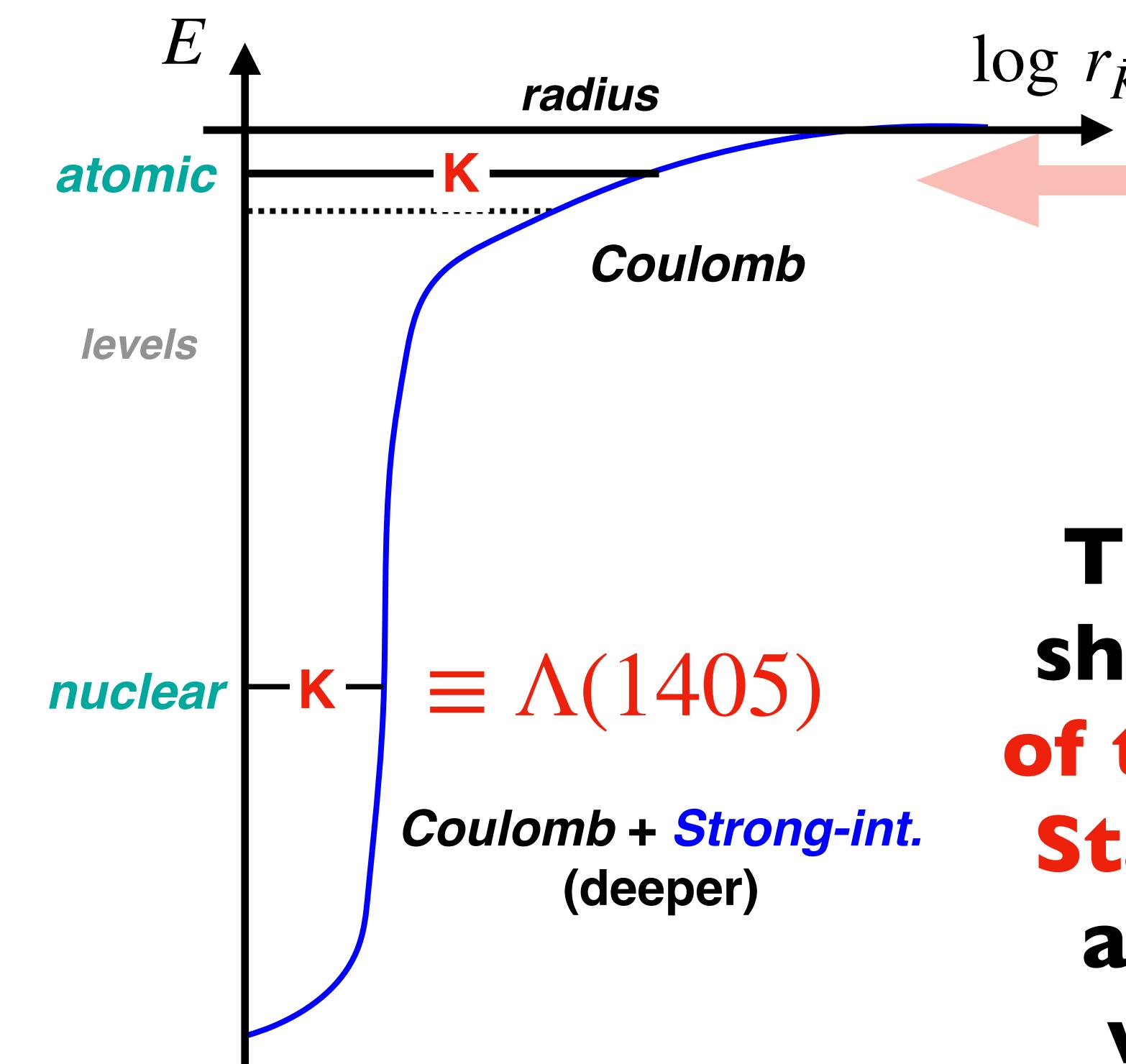


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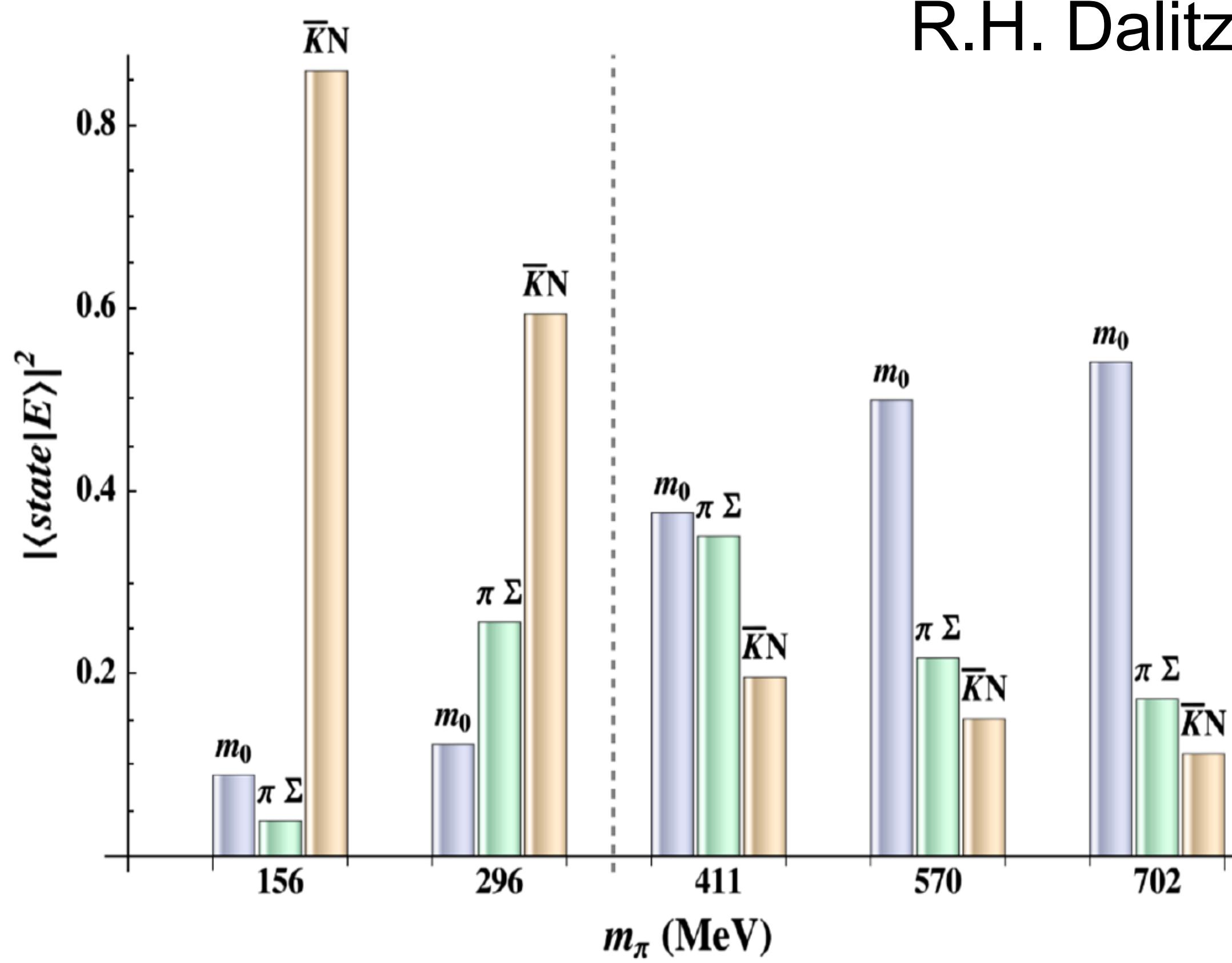
The sign flip of the energy shift suggests the existence of the Kaonic Nuclear Bound State because the repulsive atomic shift appears only when the bound state is formed!

Illustrative Animation

As a candidate of K^-p bound state, $\Lambda(1405)$ is the most natural

- Is it quark excited state of Λ baryon (qqq)?

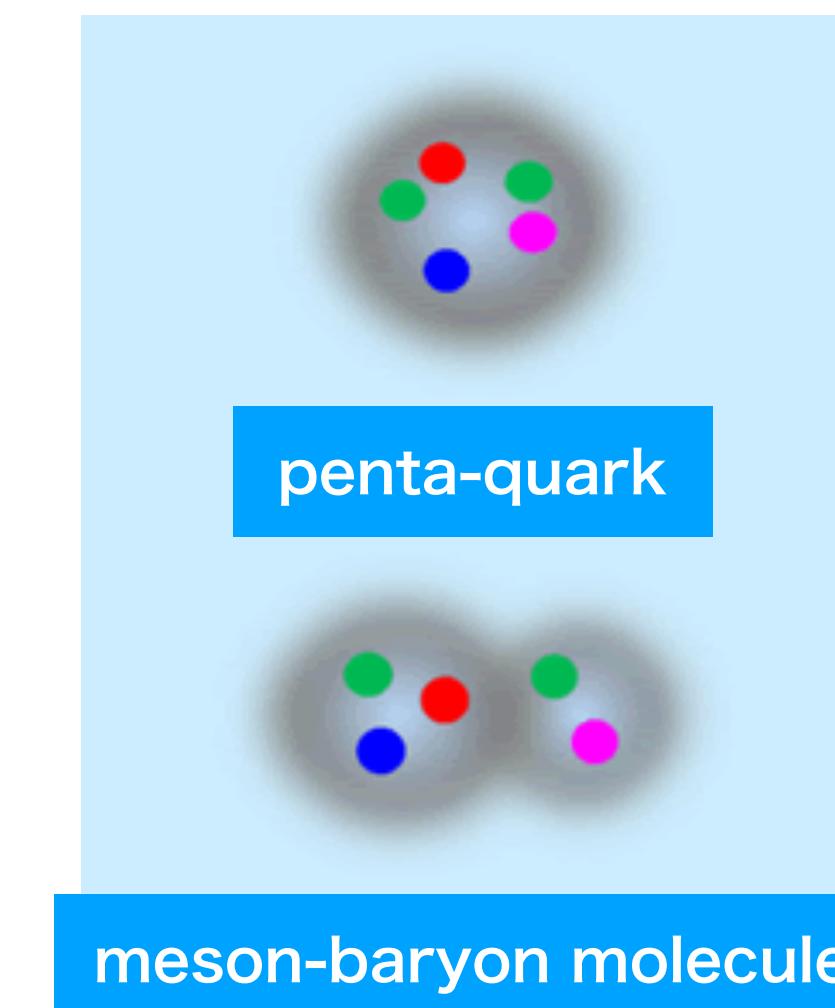
$\Lambda(1405) = \bar{K}N$... a “molecule-like hadron composite”



R.H. Dalitz and S.F. Tuan, Ann. Phys., 3, 307 (1960)

- ◆ supported by kaonic hydrogen data
Phys. Rev. Lett., 78, 3067 (1997)
- ◆ supported by Lattice QCD

J.M.M. Hall et al., Phys. Rev. Lett. 114(2015)132002.



why not $\bar{K}NN$?

forming a nuclear
bound state

Search for $\bar{K}NN$ bound state

- T. Hashimoto, S. Ajimura, G. Beer, PTEP (2015) ptep/ptv076.
S. Ajimura, H. Asano, G. Beer et al. Phys. Lett. B789 (2019) 620.
T. Yamaga, S. Ajimura, H. Asano et al. Phys. Rev. C102 (2020) 044002.

Basic understanding of nuclei

- Nuclei consist of nucleons bound by nuclear force

nucleons (N): qqq

$q = u$ or d

Fermion:

Pauli exclusion

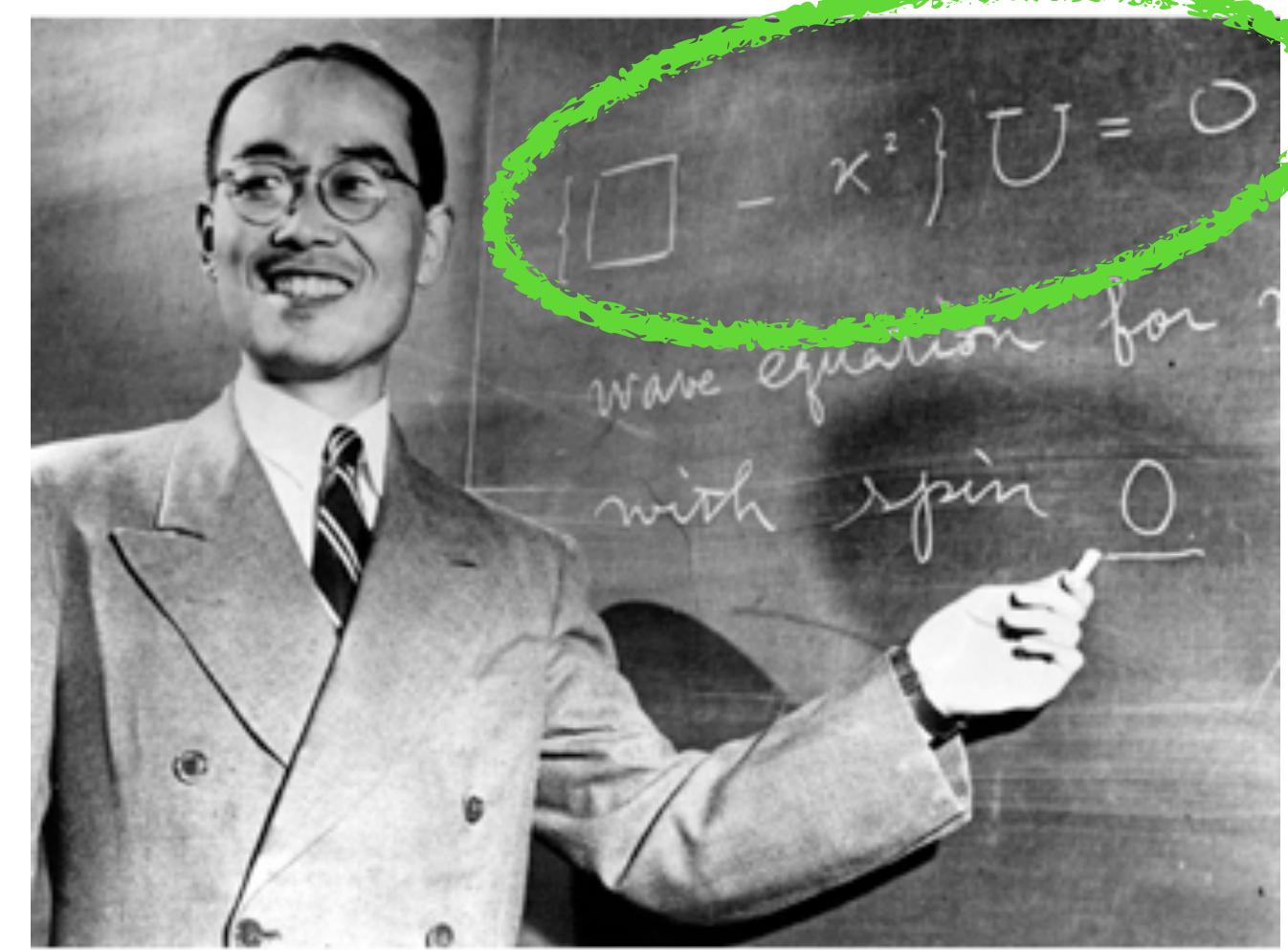
meson: $\bar{q}q$

Boson:

particles can share a quantum state

Z [e]	1st	2nd	3rd
2	u	c	t
+ 3			
1	d	s	b
- 3			

quark flavor



Yukawa Theorem tells :

- in nuclei, mesons are virtual particles and form nuclear potential
- in vacuum, mesons are real particles having own intrinsic masses

Long standing question :

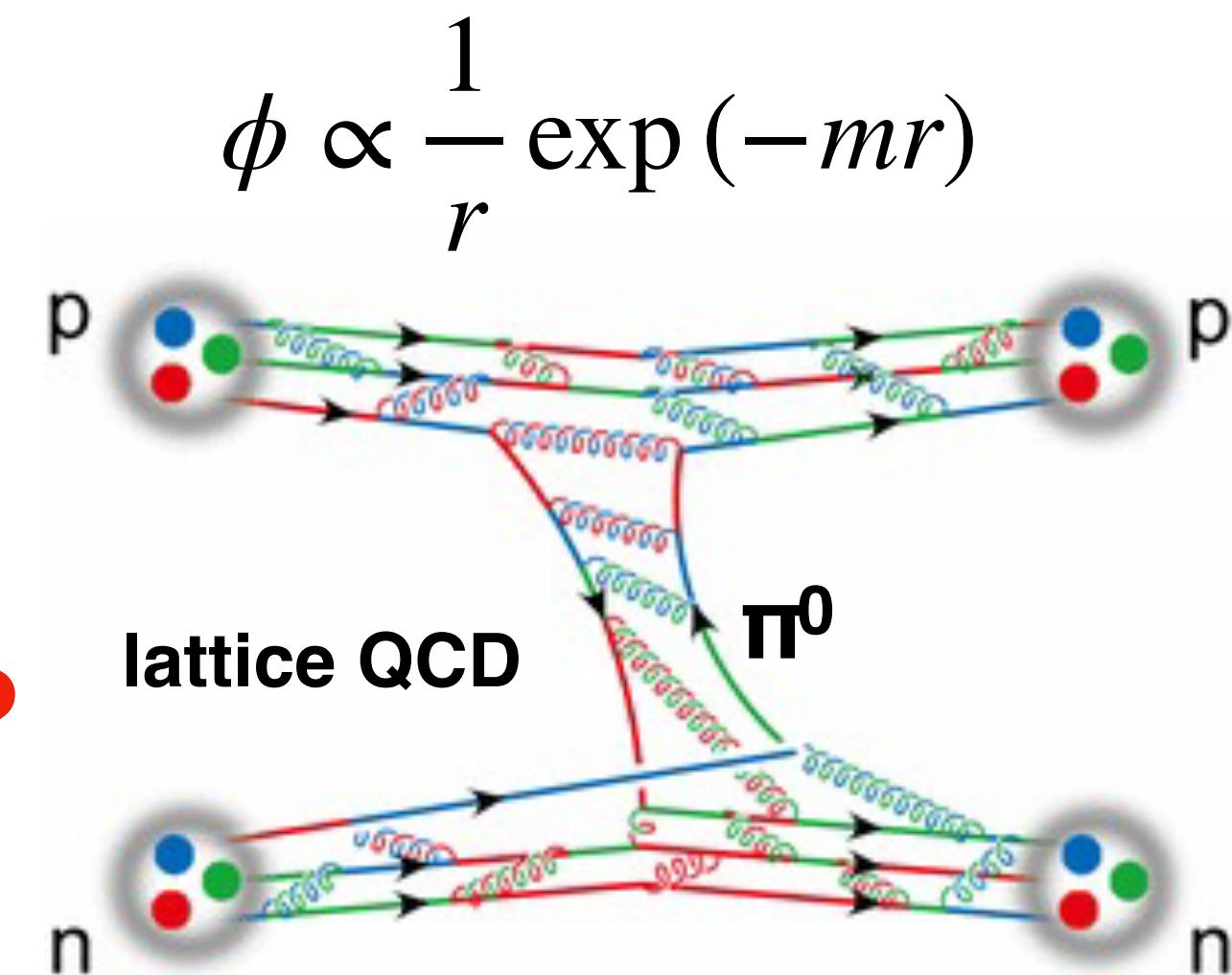
Can meson be a constituent particle forming nuclei?

— Can meson form a quantum state as a particle ? —

... in the case of K-meson ...

\bar{K} ($\bar{q}s$) forms a bound state
with two nucleons

\bar{K} meson (K^- : $\bar{u}s$, \bar{K}^0 : $\bar{d}s$)

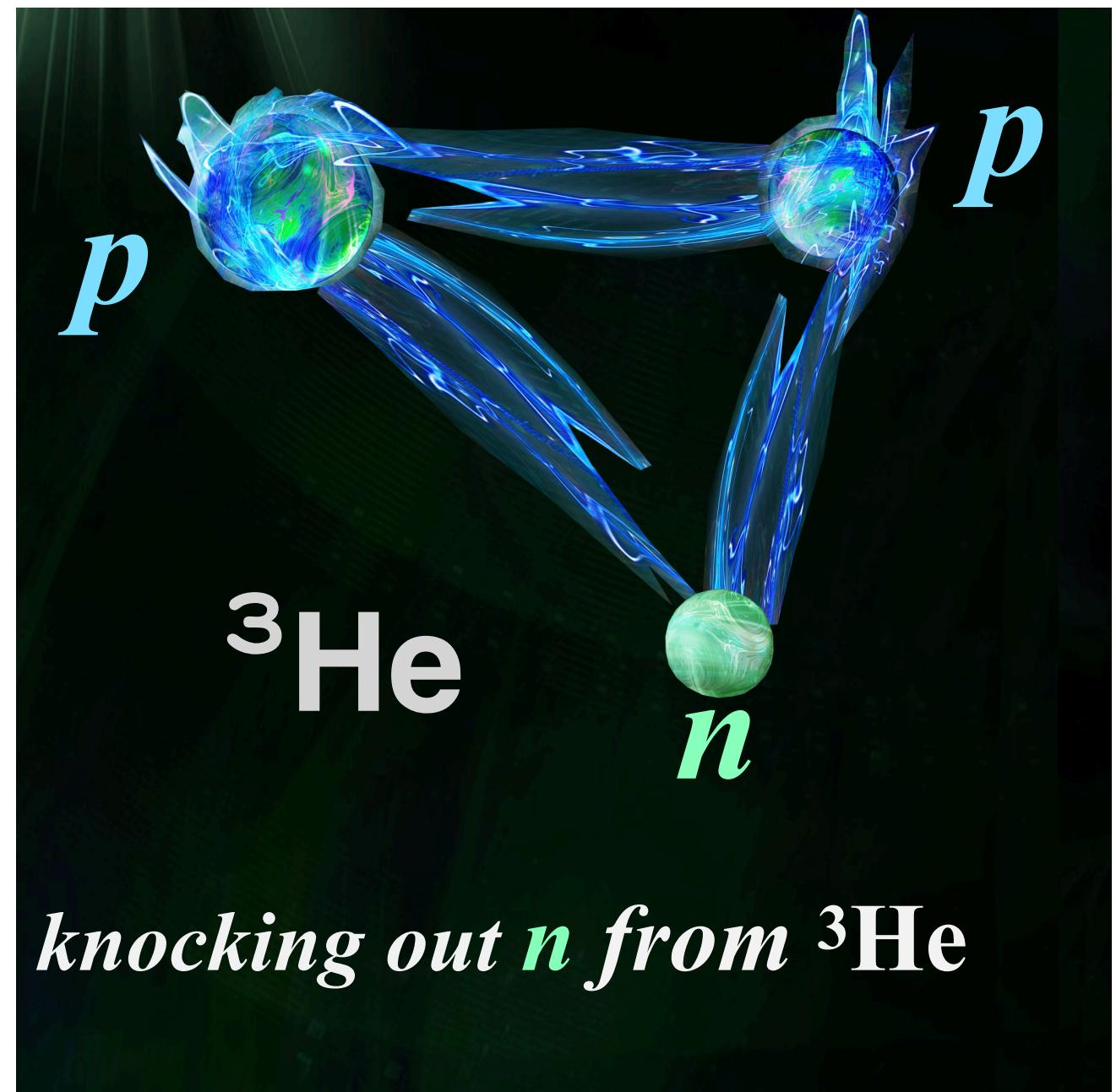


Beyond Conventional
Understanding

totally new probe (impurity)
to study inside nuclei

J-PARC E15: “ K^- -pp” Exploration

$\text{K}^- + {}^3\text{He}$ (ppn)

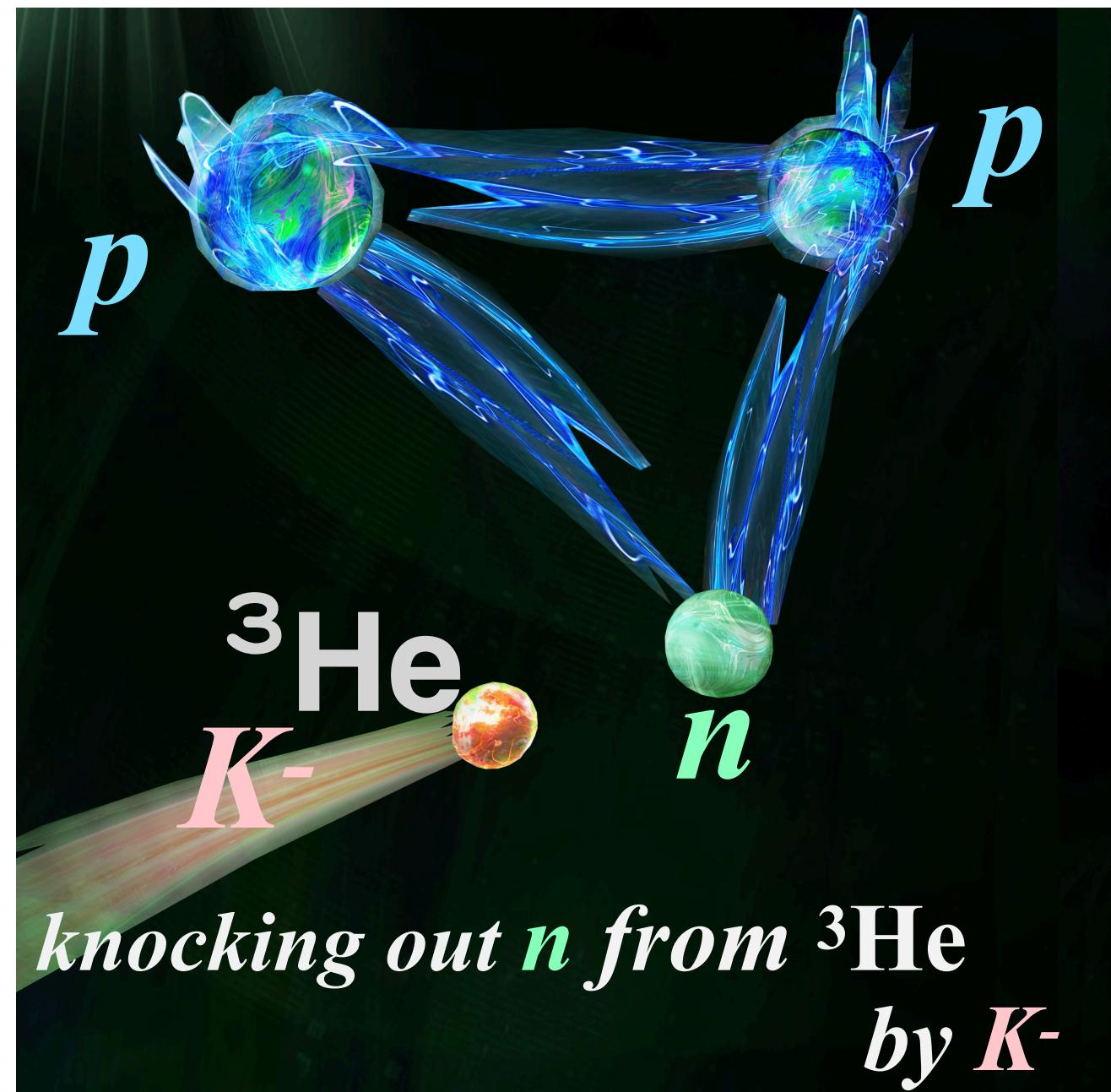


$(\text{K}^- + \text{pp}) + \text{n}$

substitute n in ${}^3\text{He}$ by K^-

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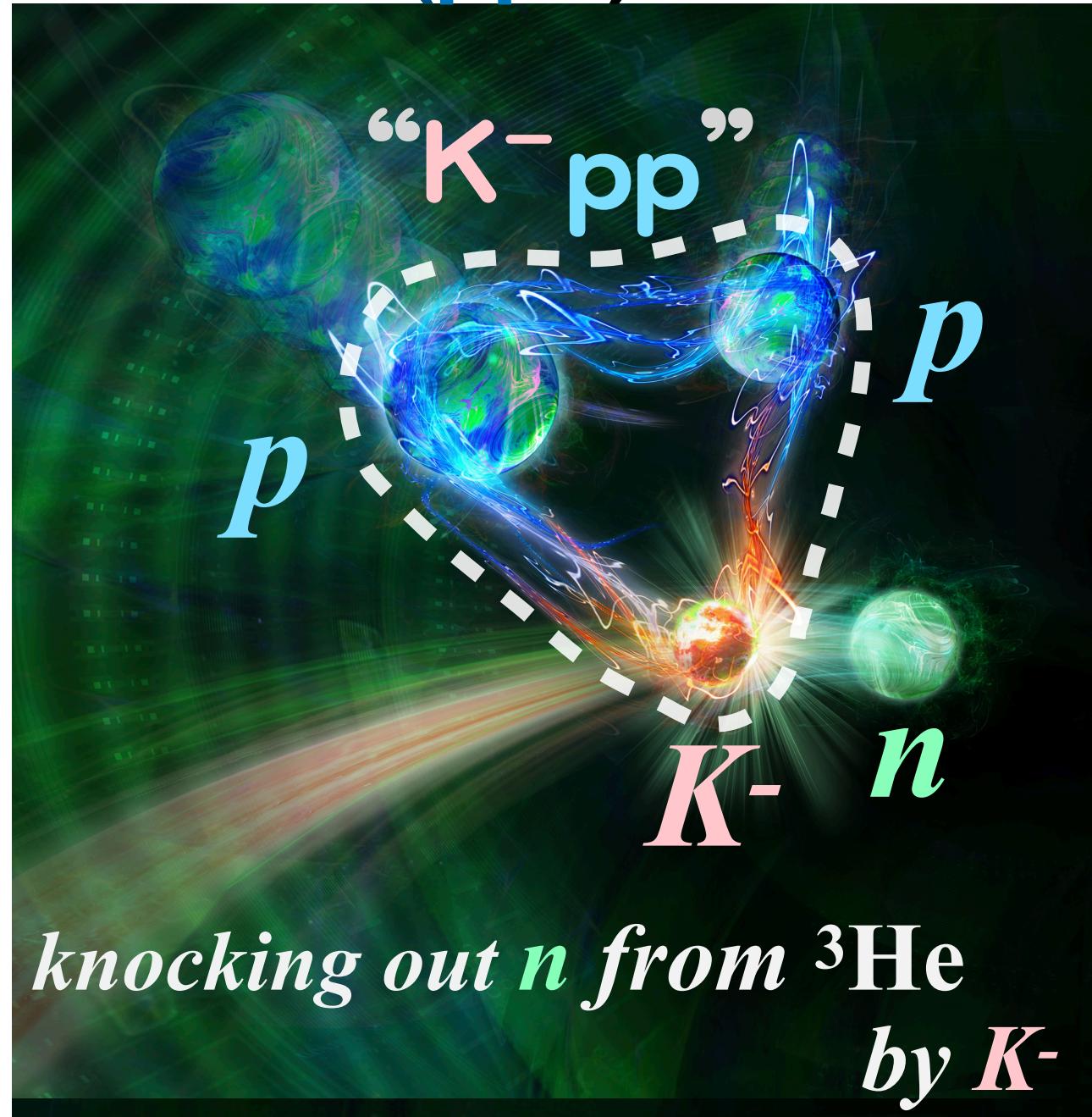


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substitute n in 3He by K^-

If “ K^-pp ” exists, a peak will be formed in invariant mass spectrum below $M(K^-pp)$

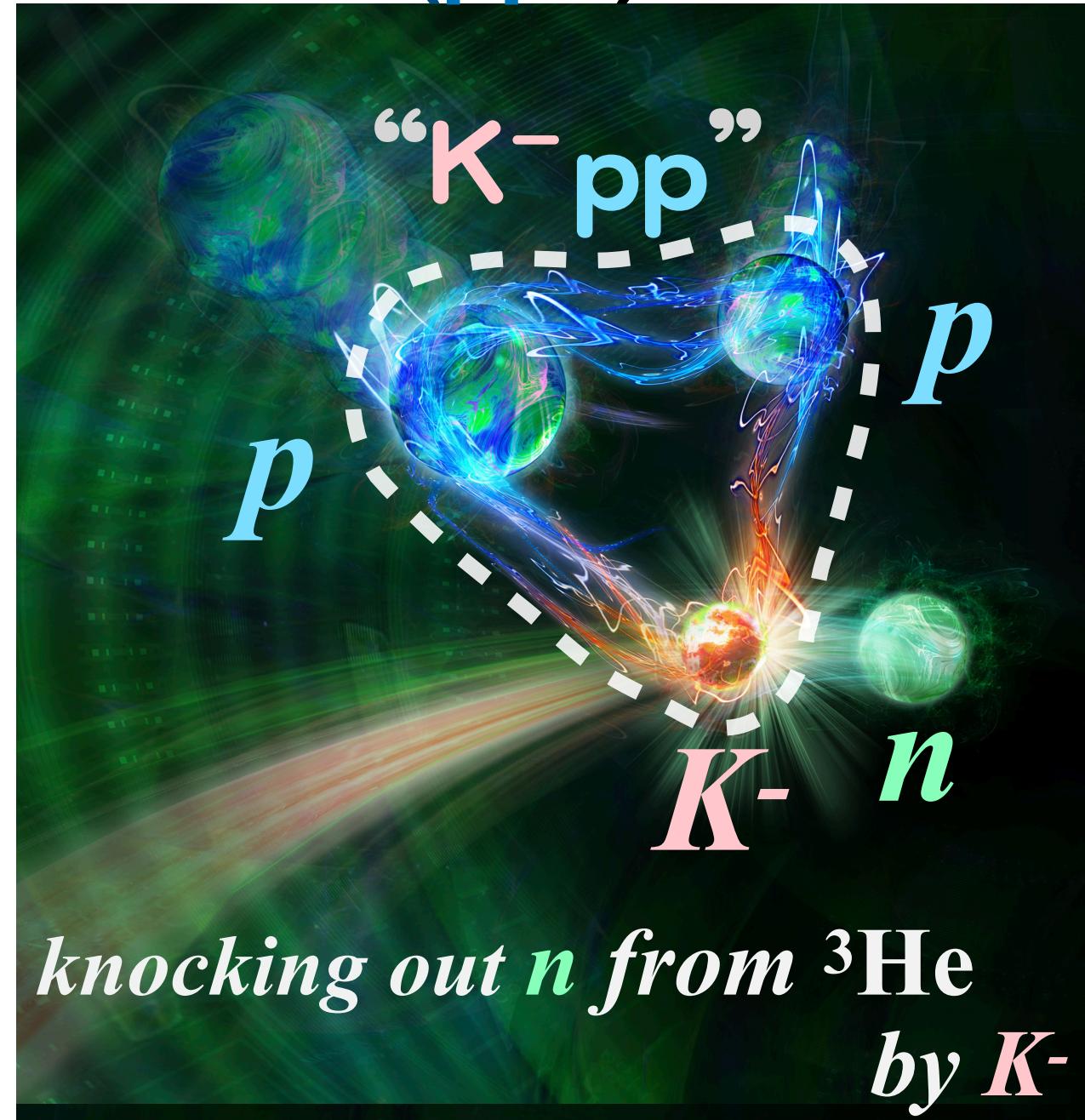
$$M(K^-pp) \equiv m_{K^-} + 2m_p$$



: formation

J-PARC E15: “ K^-pp ” Exploration

$K^- + {}^3He$ (ppn)

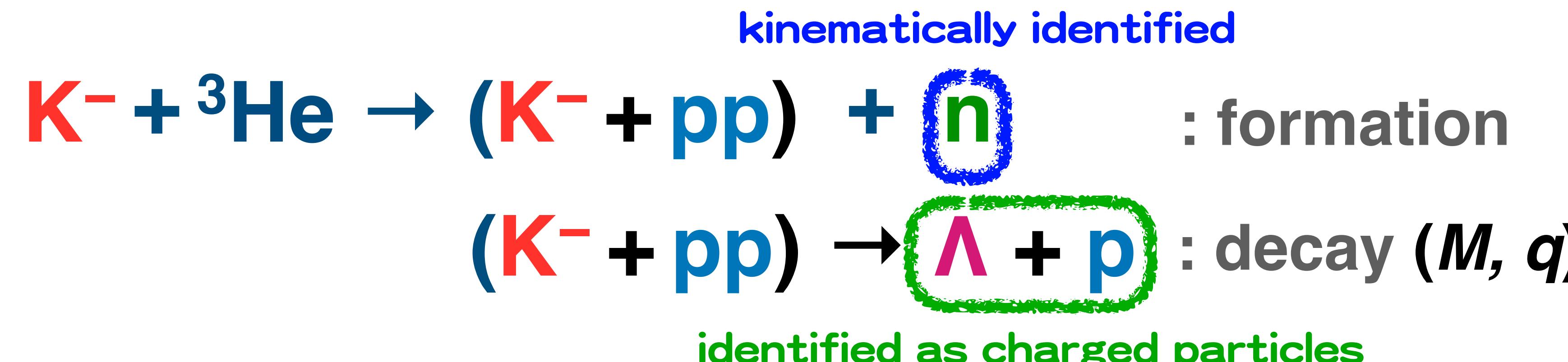


$(K^- + pp) + n$

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select $K^- + {}^3He \rightarrow (\Lambda + p) + n$ events,
 analyze *invariant mass M* of $(K^- + pp)$ -system
 and *momentum transfer q* to the system

provides multi-dimensional kinematical information



on (M, q) -plane

q -distribution: system size

- high- q capture happens if the system is compact –

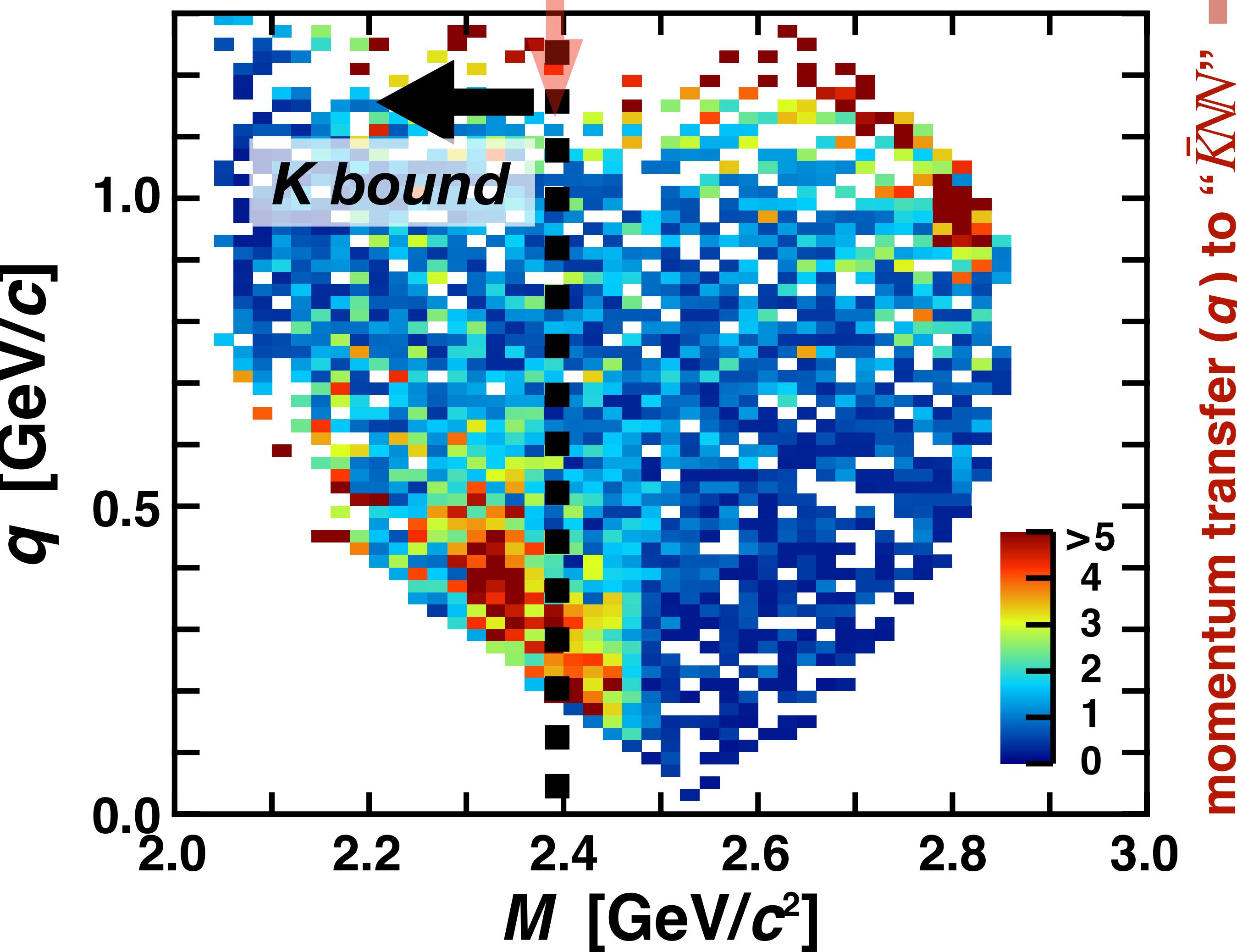
M -distribution: binding energy & absorption width

- sensitive to $\bar{K}N$ interaction –

Acceptance corrected event distribution on (M, q)

reconstructed “ $\bar{K}N\bar{N}$ ” mass (M) →

binding threshold




 Δp + n events
 on (M, q) -plane

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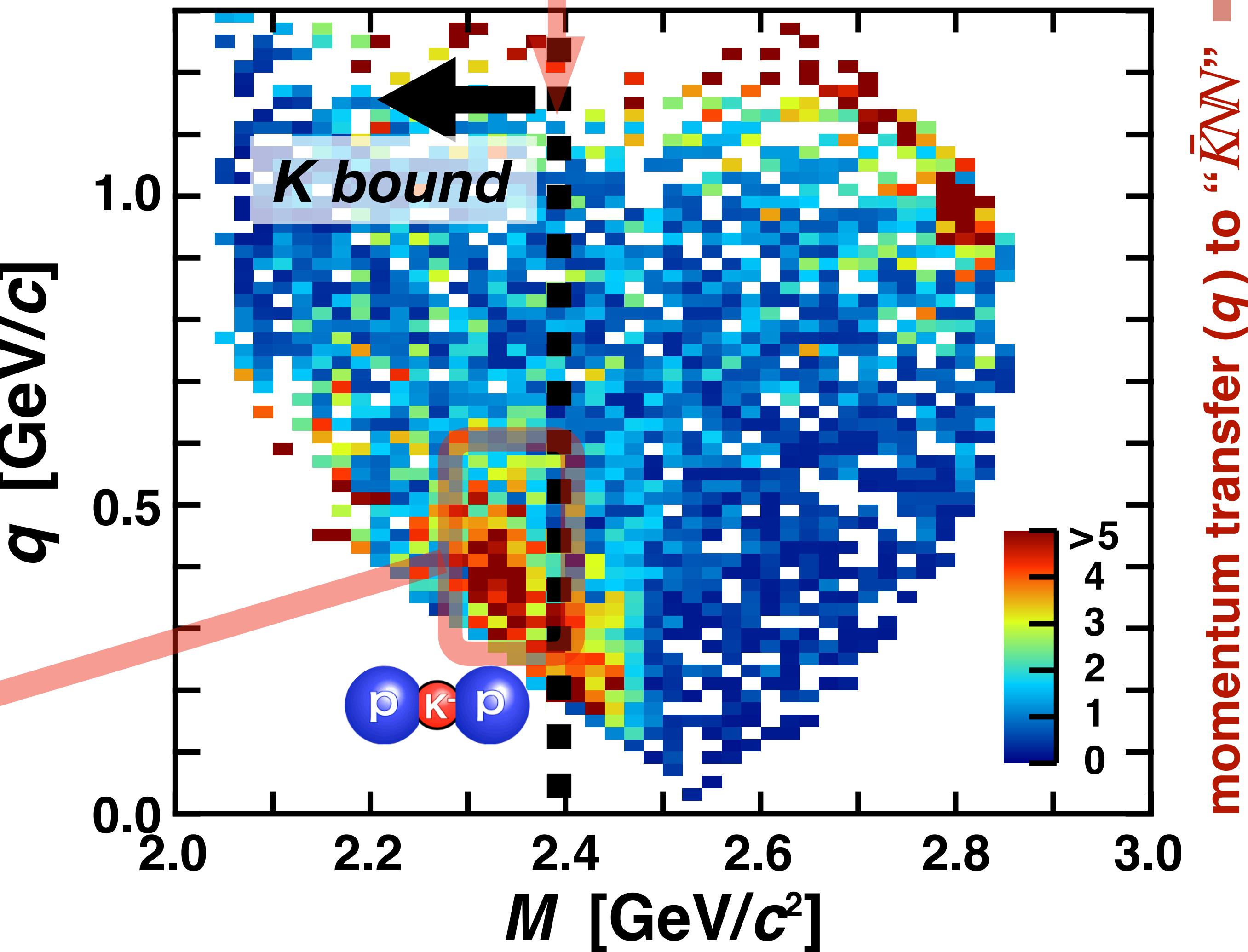
**The K-pp signal
is clearly seen on (M, q) -plane!**

- relatively deep and wide, and extended to high- q region –

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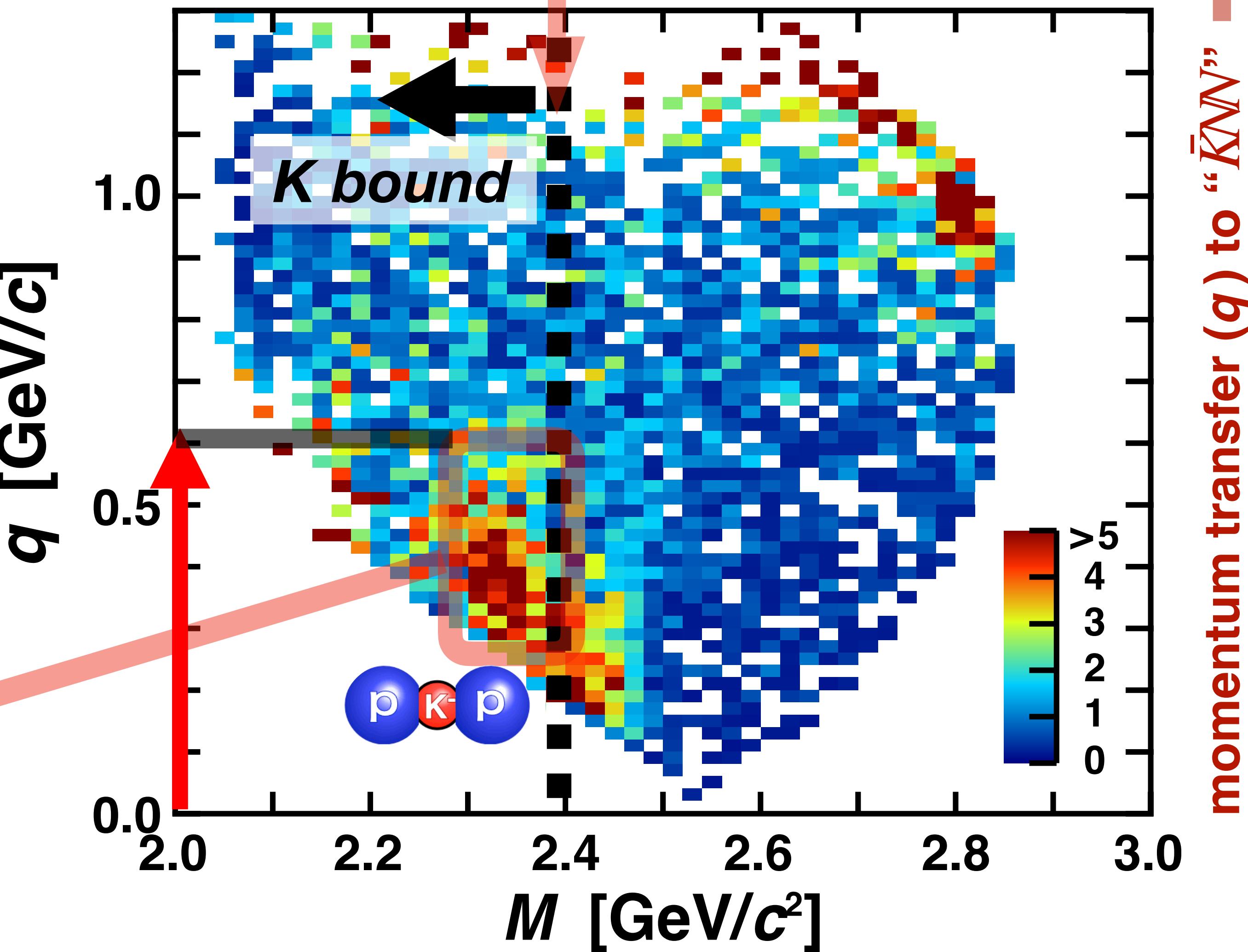
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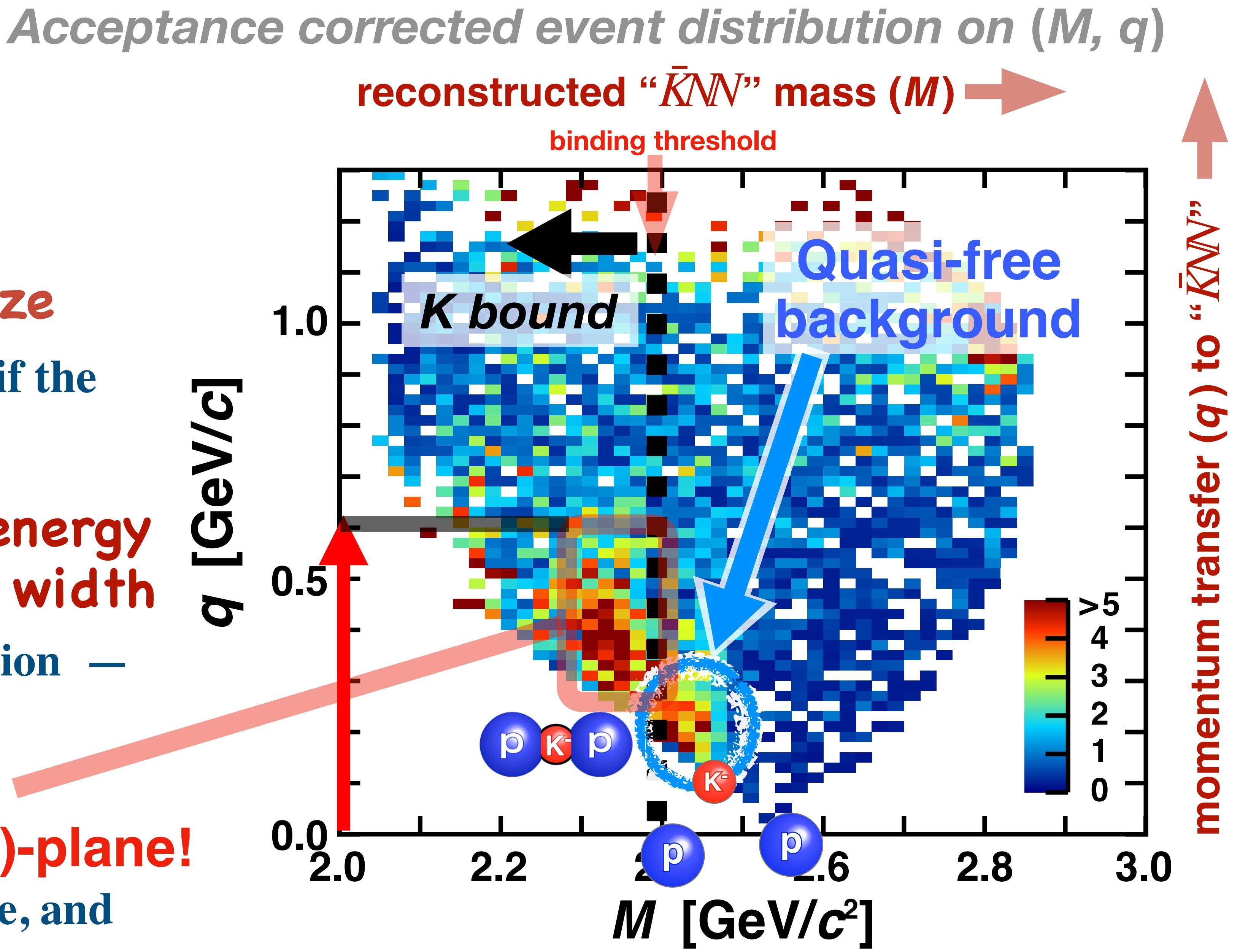
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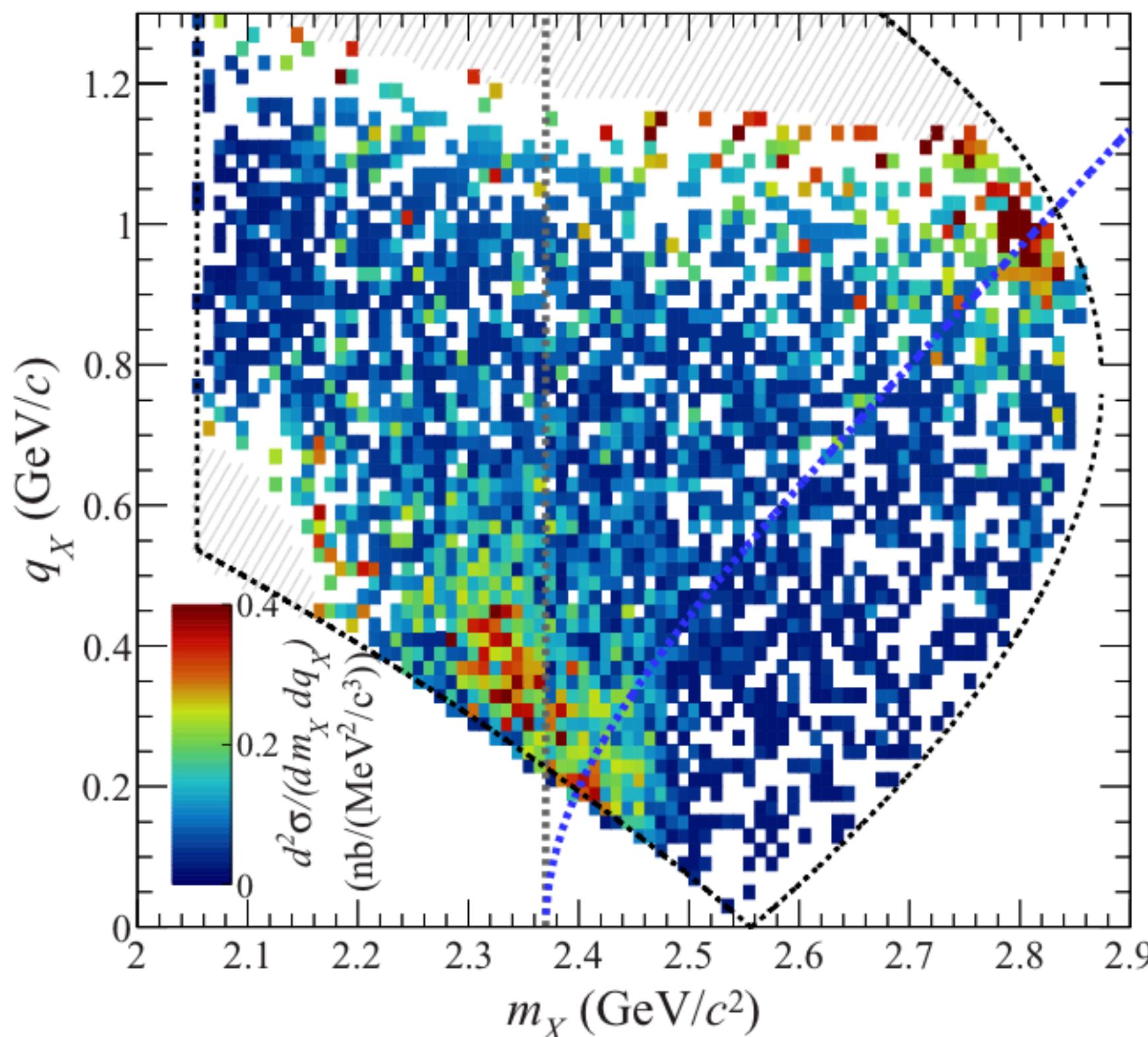


Λp + n_{miss}

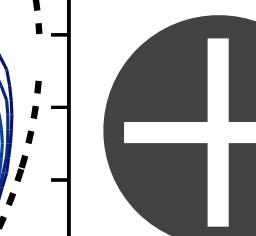
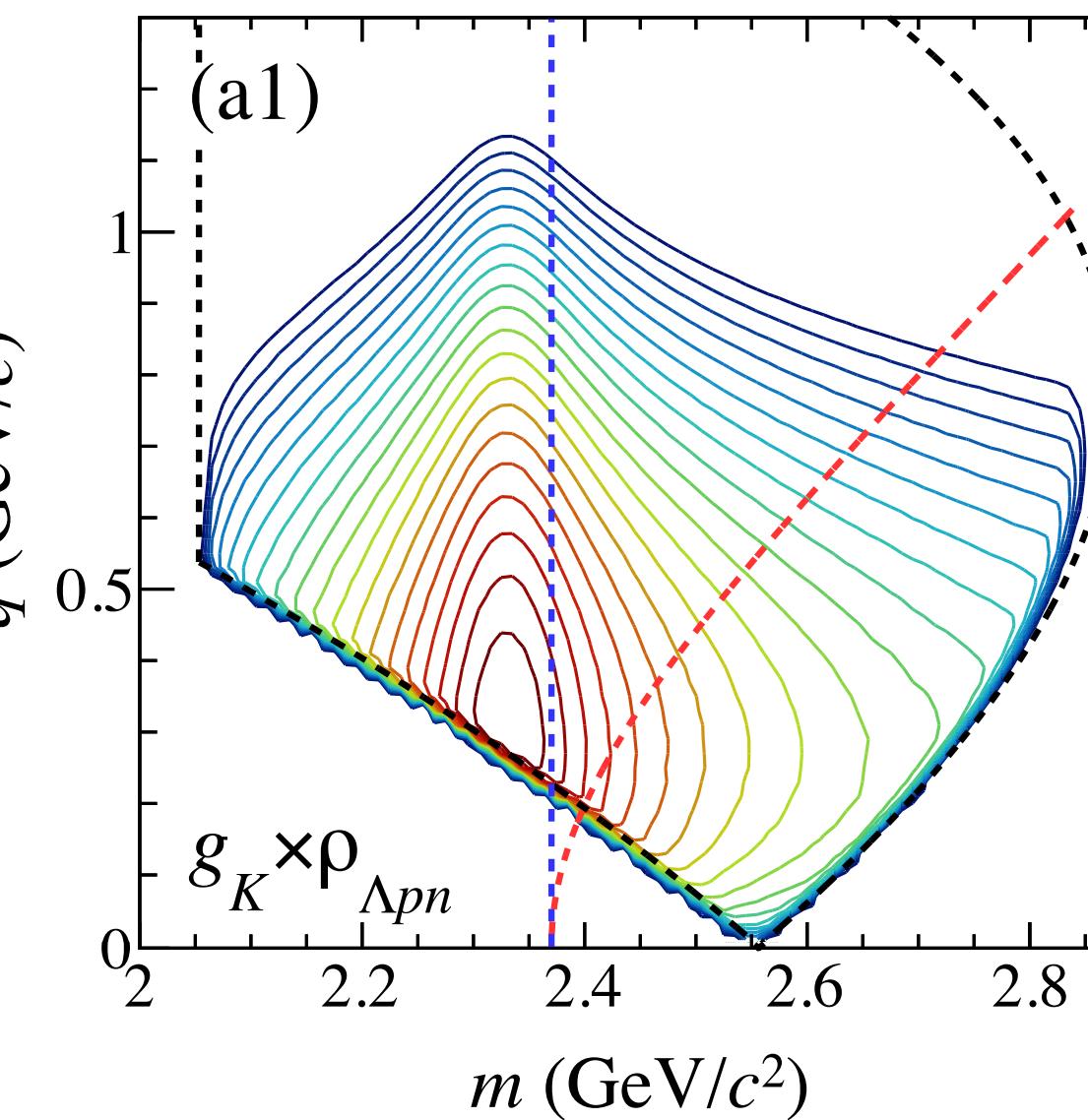
Phenomenological model fitting function in (m, q) -plane

ρ : Lorentz-invariant phase-space

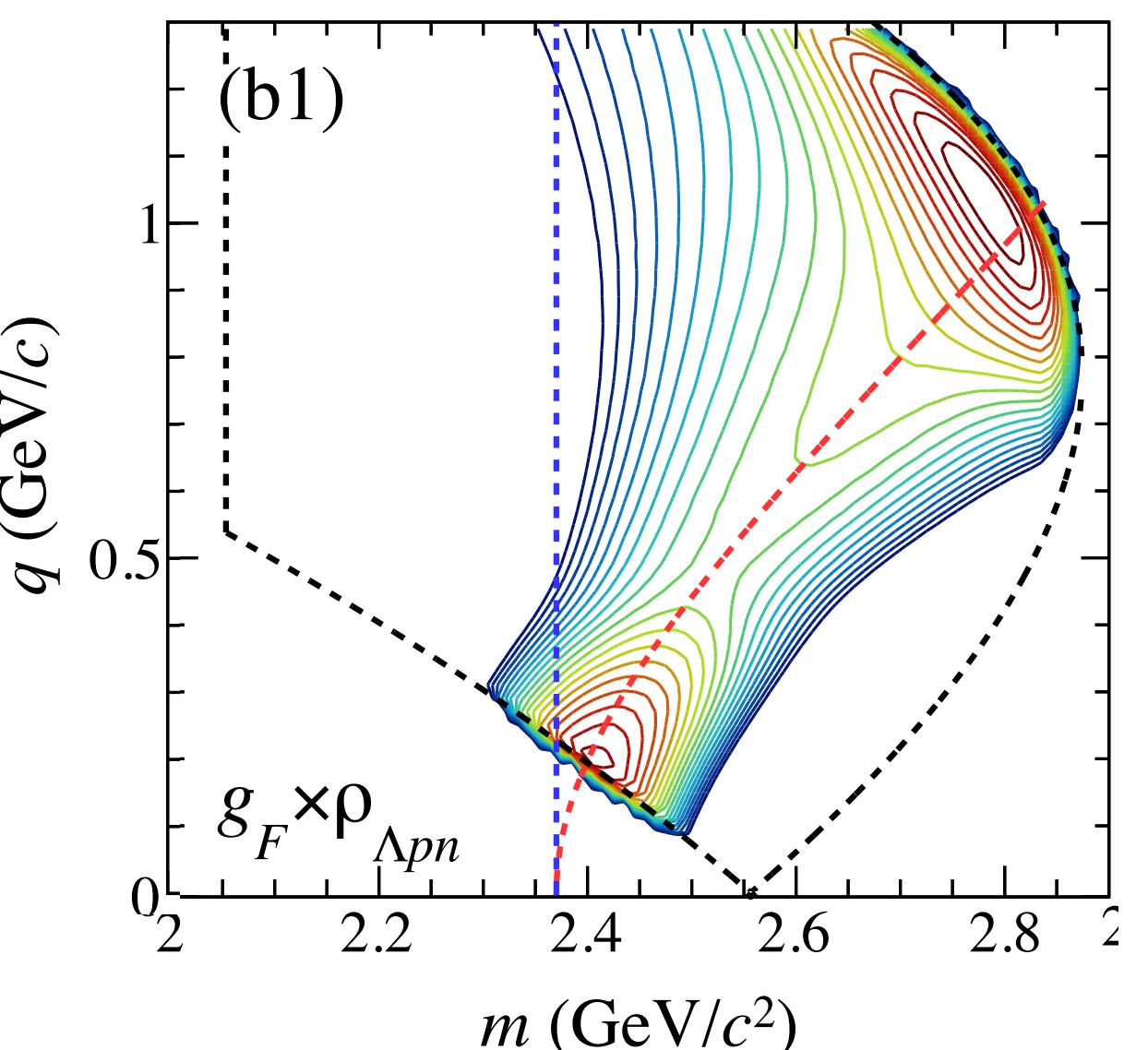
$$f_{\bar{K}NN}(m, q) \times \rho_{\{\Lambda p n\}}(m, q) \quad f_{QF-\bar{K}}(m, q) \times \rho_{\{\Lambda p n\}}(m, q)$$



$\bar{K}NN$ production



QF- \bar{K} absorption



$f_{\bar{K}NN}(m, q) :$ $B.W.(m) \times$
form factor(q)

$f_{QF-\bar{K}}(m, q) :$
phenomenological

PWIA based interpretation

(plane wave impulse approximation)

$$\sigma(M, q) \propto \rho_{3B}(M, q) \times \frac{(\Gamma_{Kpp}/2)^2}{(M - M_{Kpp})^2 + (\Gamma_{Kpp}/2)^2} \times \exp\left(-\frac{q^2}{Q_{Kpp}^2}\right)$$

Differential
cross section

PWIA based interpretation

(plane wave impulse approximation)

$$\sigma(M, q) \propto \frac{\rho_{3B}(M, q)}{\text{Lorentz invariant phase space } (\Lambda p n)} \times \frac{(\Gamma_{Kpp}/2)^2}{(M - M_{Kpp})^2 + (\Gamma_{Kpp}/2)^2} \times \exp\left(-\frac{q^2}{Q_{Kpp}^2}\right)$$

Differential cross section

PWIA based interpretation

— from time integral —

B.W. / Lorentzian

(plane wave impulse approximation)

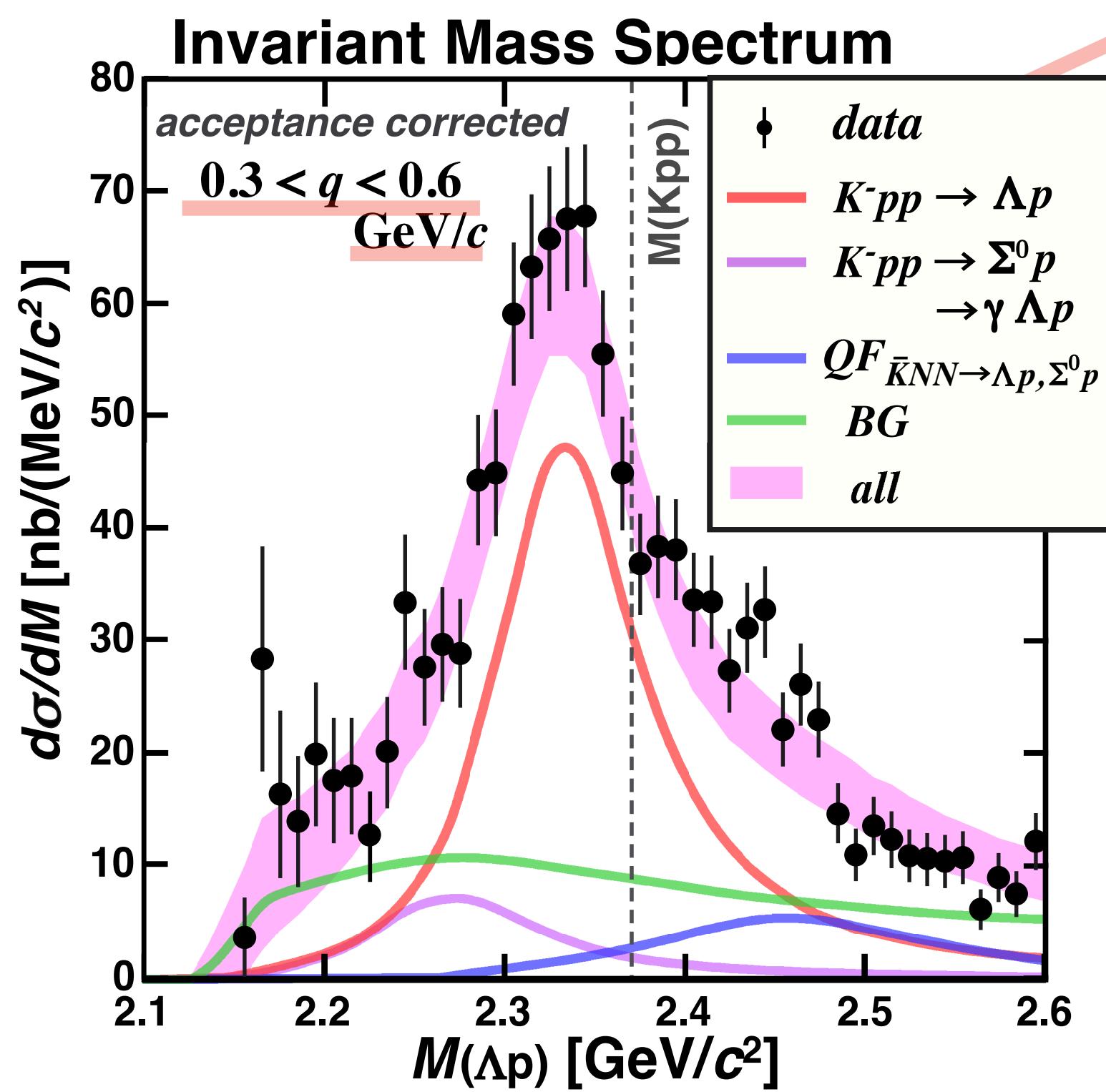
$$\sigma(M, q) \propto$$

Differential
cross section

$$\rho_{3B}(M, q) \times$$

Lorentz invariant
phase space ($\Lambda p n$)

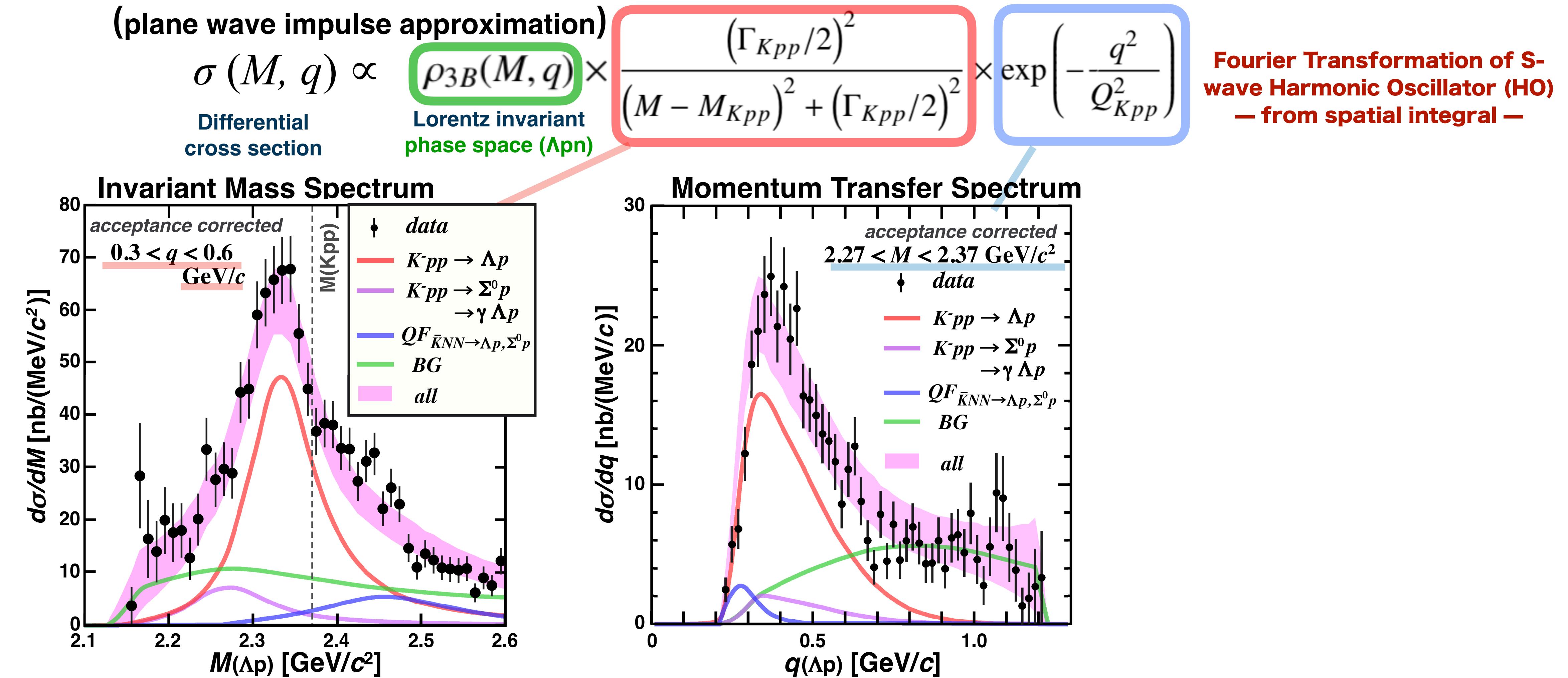
$$\frac{(\Gamma_{Kpp}/2)^2}{(M - M_{Kpp})^2 + (\Gamma_{Kpp}/2)^2} \times \exp\left(-\frac{q^2}{Q_{Kpp}^2}\right)$$



strong binding ($\bar{K}N$ attraction)

$B_{Kpp} \sim 40 \text{ MeV}, \Gamma_{Kpp} \sim 100 \text{ MeV}$

PWIA based interpretation



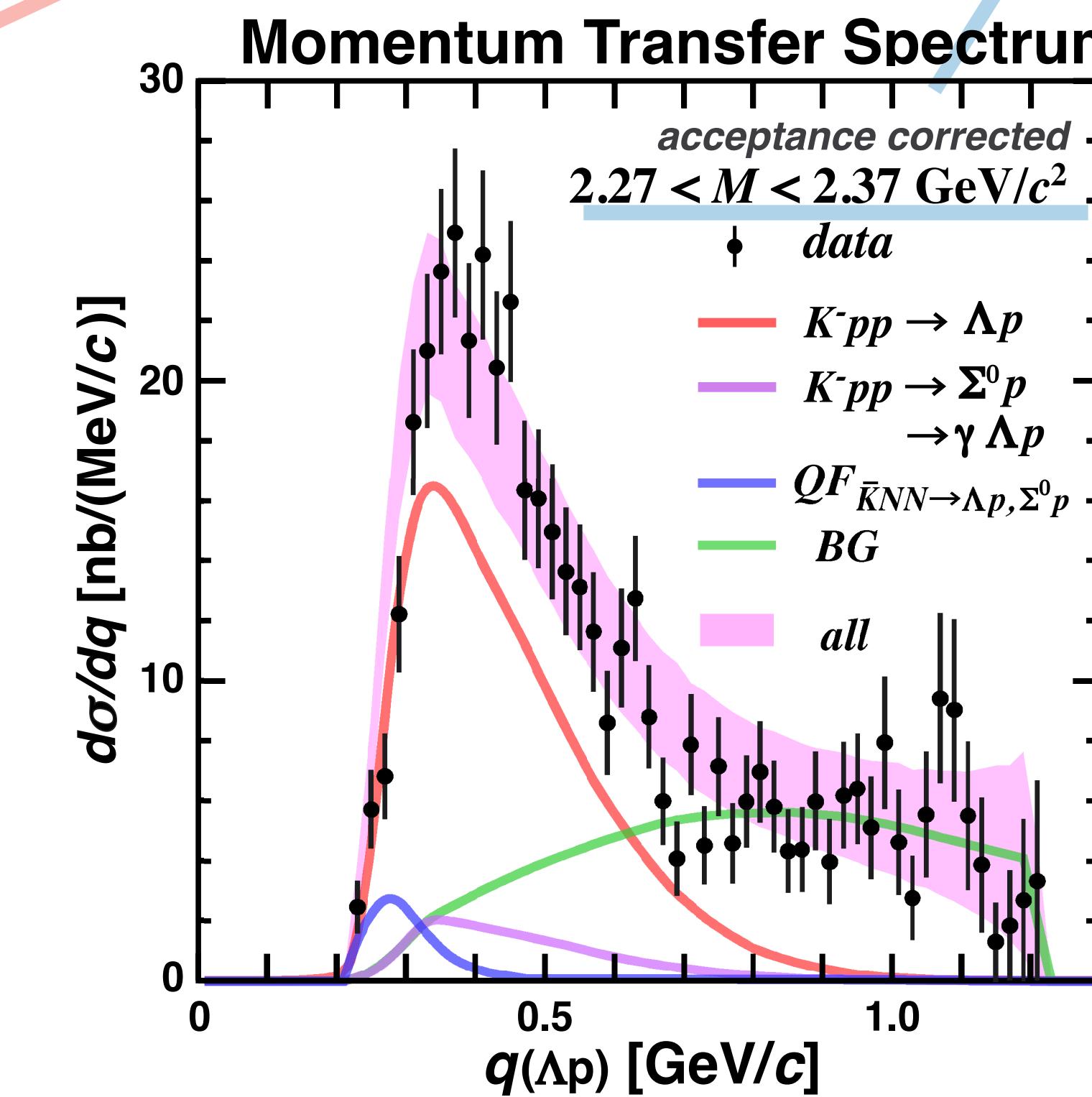
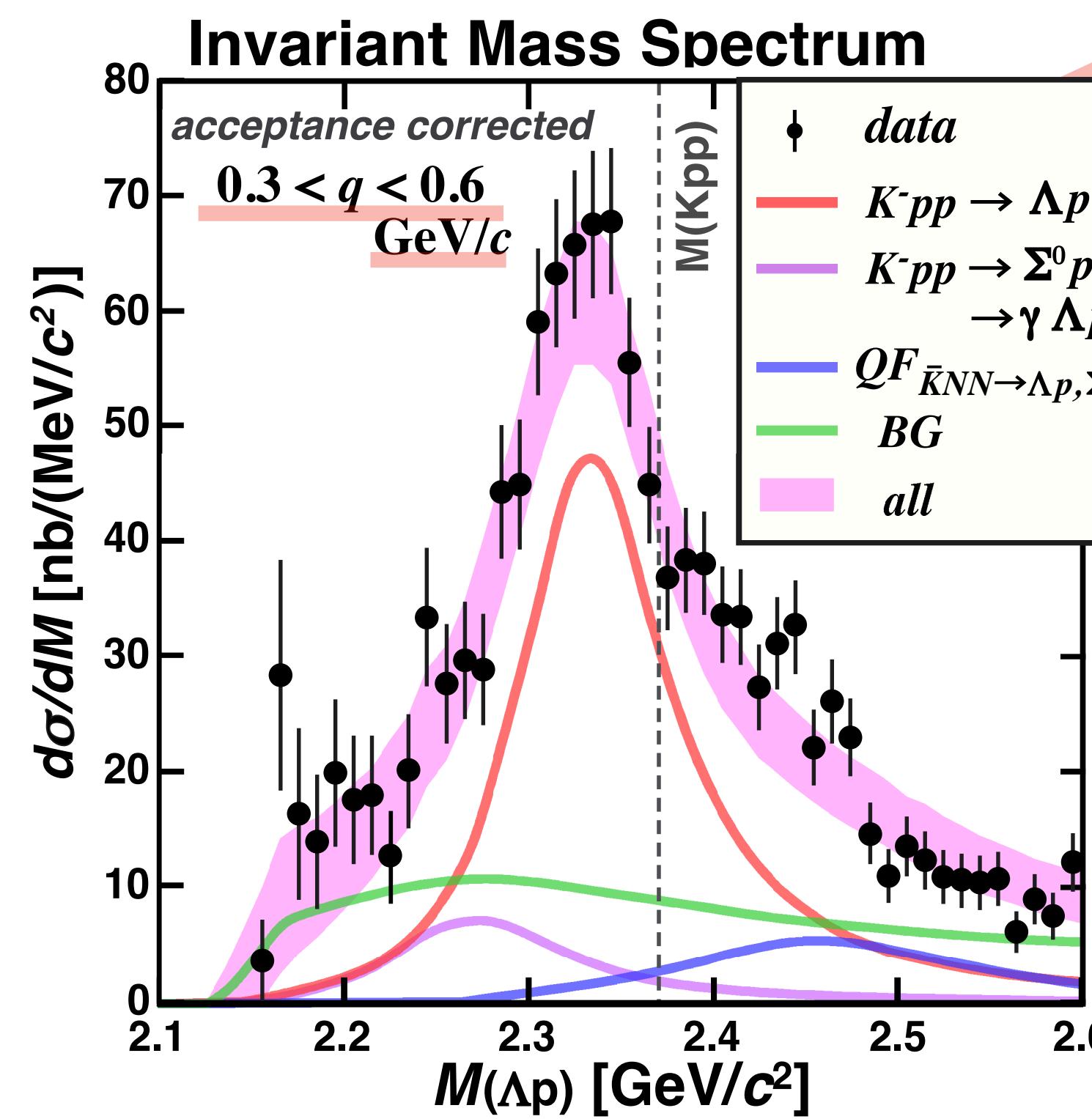
strong binding ($\bar{K}N$ attraction)
 $B_{Kpp} \sim 40 \text{ MeV}, \Gamma_{Kpp} \sim 100 \text{ MeV}$

wide momentum width
 $Q_{Kpp} \sim 400 \text{ MeV}/c$

PWIA based interpretation

(plane wave impulse approximation)

$$\sigma(M, q) \propto \text{Differential cross section} \times \rho_{3B}(M, q) \times \text{Lorentz invariant phase space } (\Lambda p n)$$



strong binding ($\bar{K}N$ attraction)
 $B_{Kpp} \sim 40 \text{ MeV}, \Gamma_{Kpp} \sim 100 \text{ MeV}$

wide momentum width
 $Q_{Kpp} \sim 400 \text{ MeV}/c$

... could be quite compact ...
 $(R_{Kpp} \sim 0.6 \text{ fm (H.O.)})$

— from time integral — — from space integral —

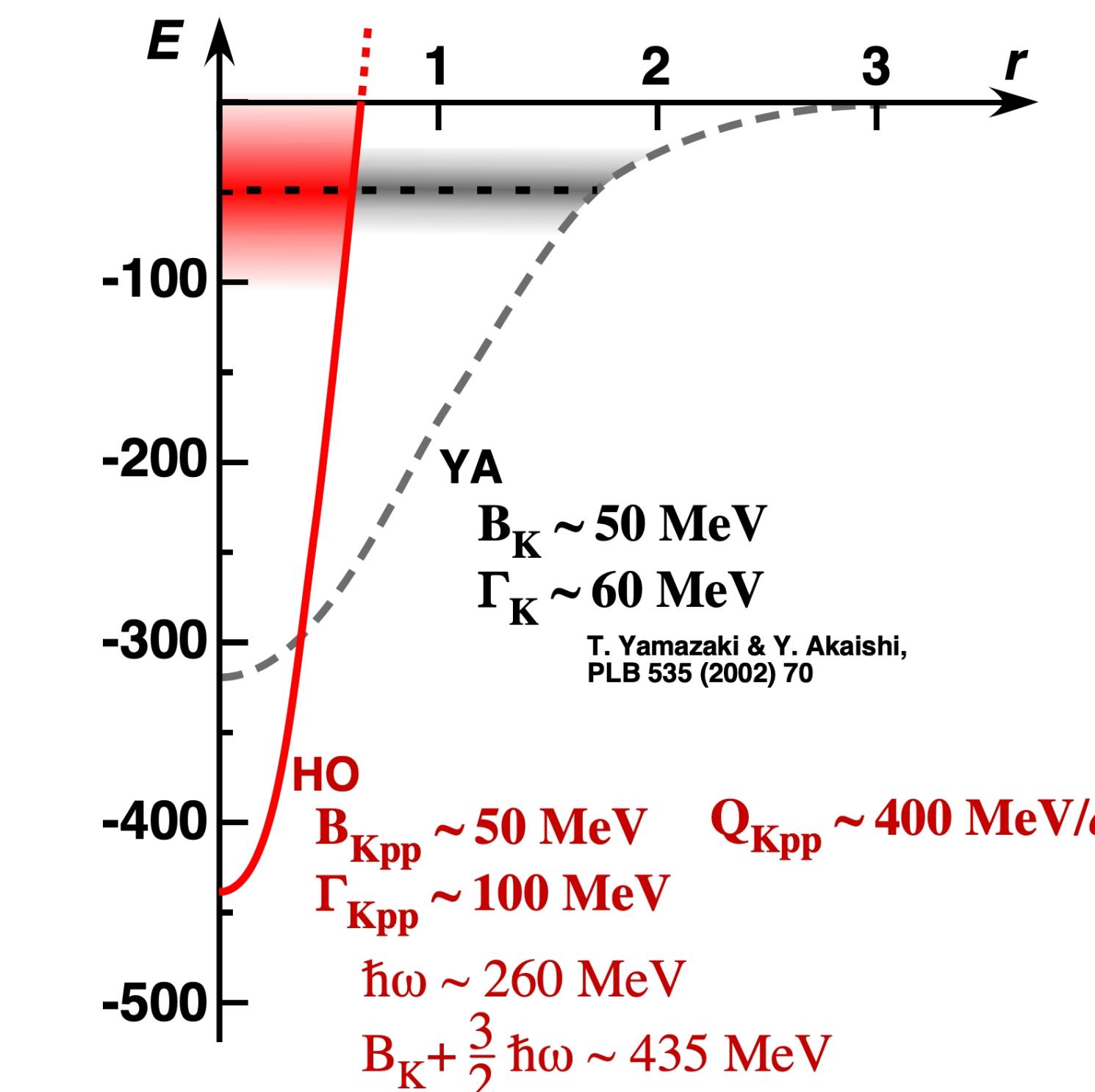
B.W. / Lorentzian

Gaussian form factor / structure factor

Fourier Transformation of S-wave Harmonic Oscillator (HO)

— from spatial integral —

$$\frac{\left(\Gamma_{Kpp}/2\right)^2}{\left(M - M_{Kpp}\right)^2 + \left(\Gamma_{Kpp}/2\right)^2} \times \exp\left(-\frac{q^2}{Q_{Kpp}^2}\right)$$

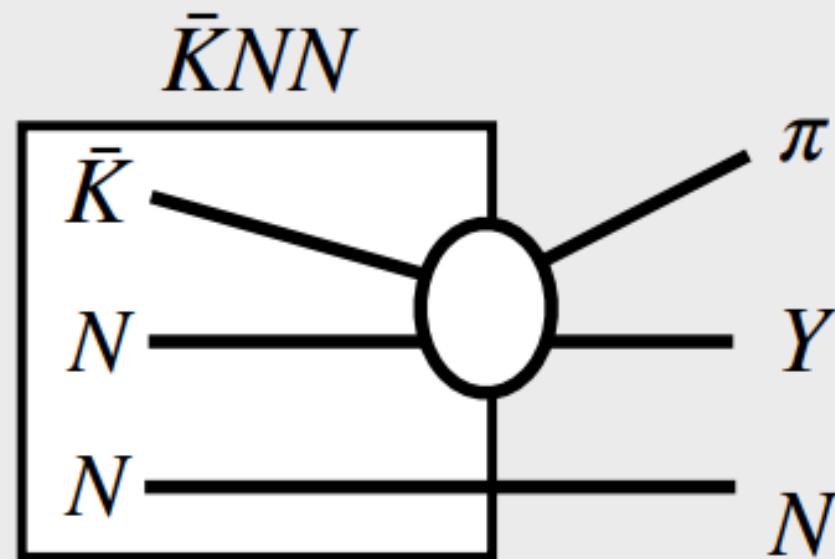


Mesonic Decay Modes of $\bar{K}NN$

T. Yamaga et.al., Phys. Rev. C 110, 014002 (2024)

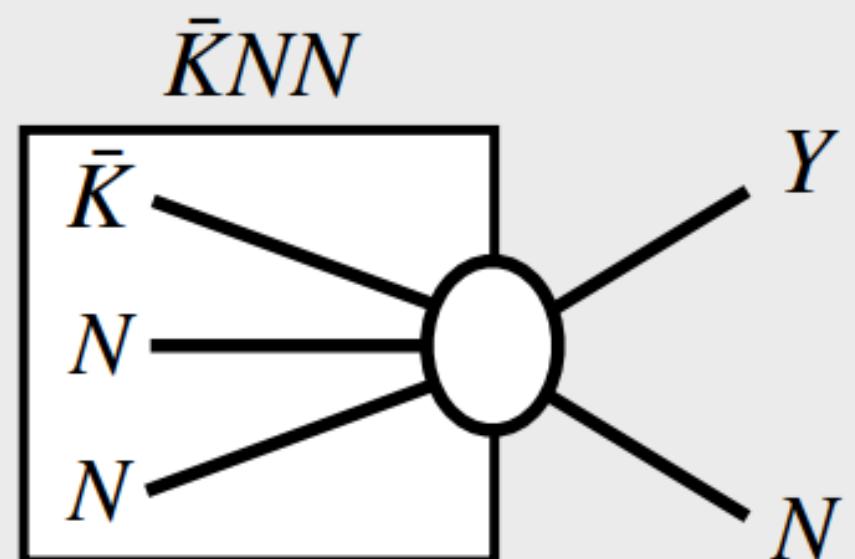
- Mesonic decays will give us further information on $\bar{K}NN$
 - ✓ internal structure
 - ✓ $\bar{K}N$ interaction below the threshold

Mesonic



1N absorption

Non-mesonic



2N absorption

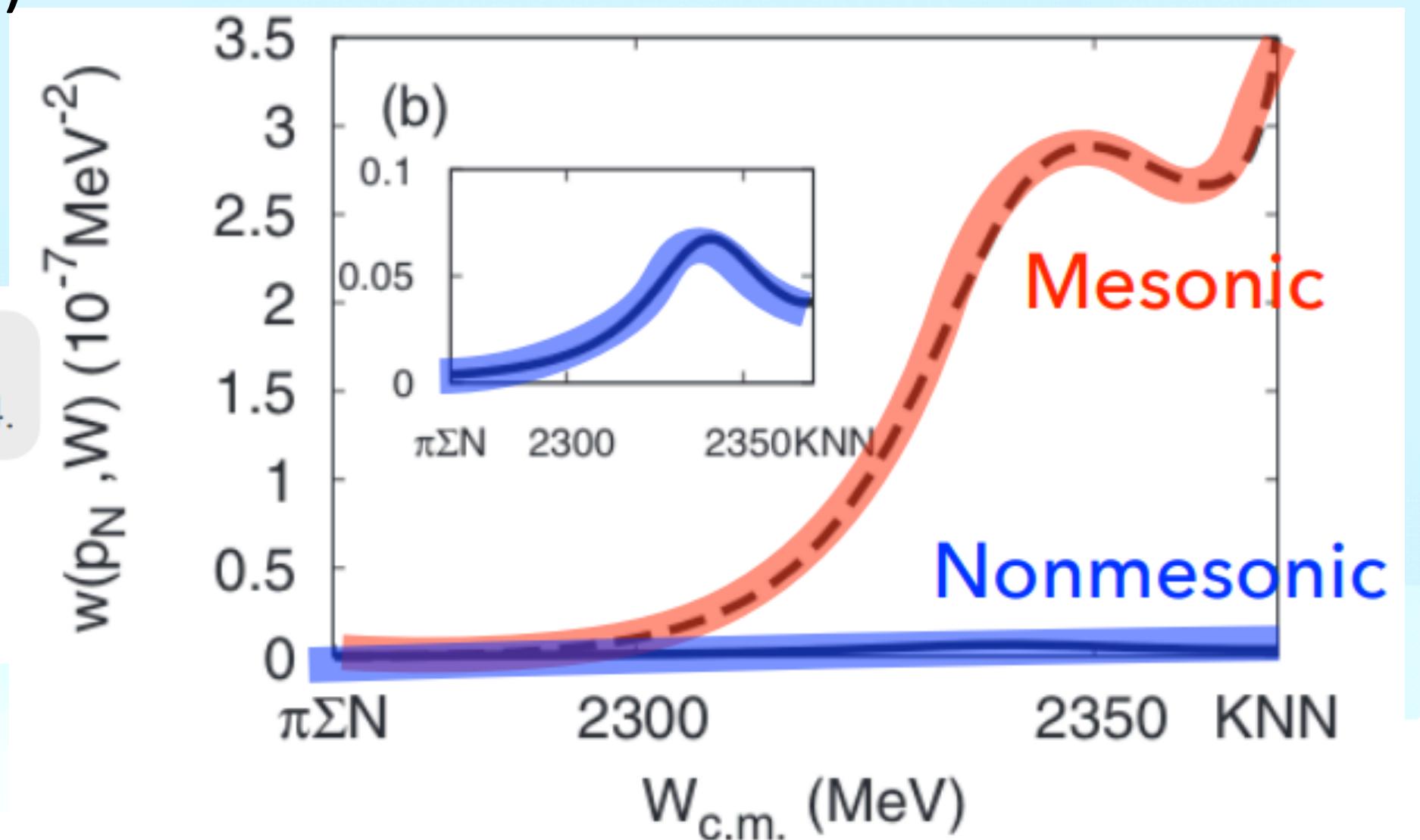
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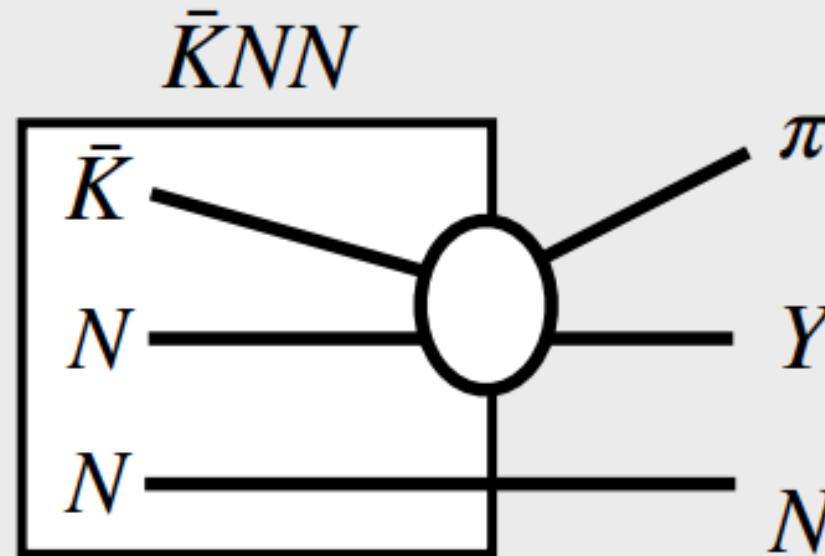
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 - ✓ internal structure
 - ✓ $\bar{K}N$ interaction below the threshold

S. Ohnishi, et al.,
Phys. Rev. C 88 (2013) 025204.

$$\Gamma_{YN} \ll \Gamma_{\pi YN}$$

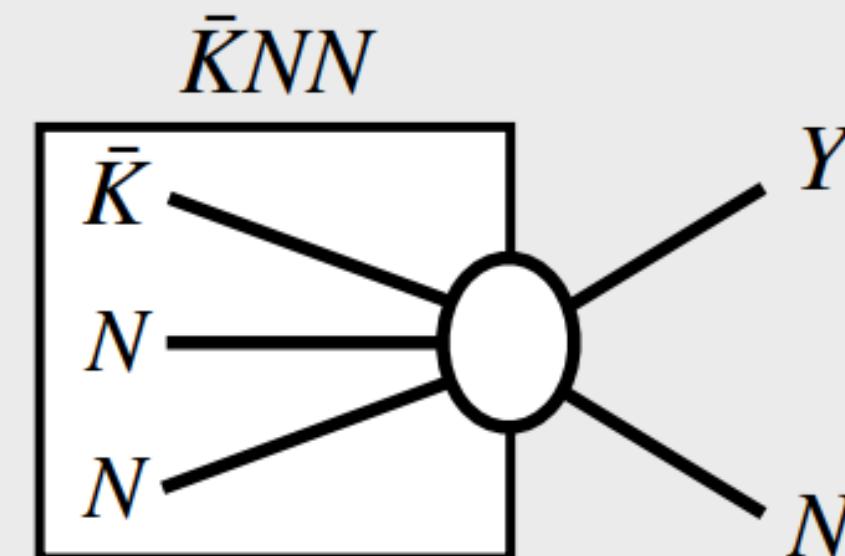


Mesonic



1N absorption

Non-mesonic



2N absorption

Mesonic Decay Modes of $\bar{K}NN$

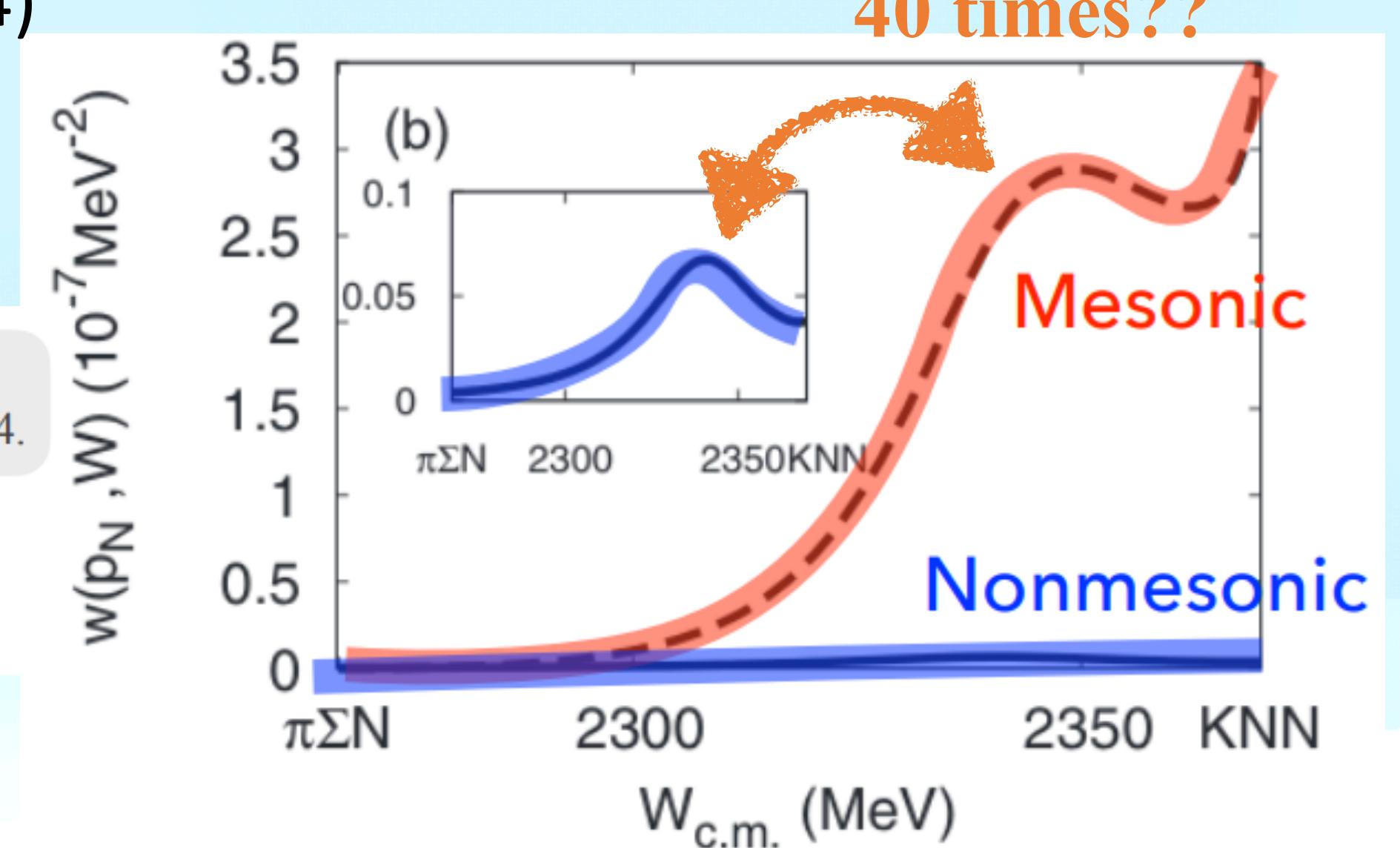
T. Yamaga et.al., Phys. Rev. C 110, 014002 (2024)

- Mesonic decays will give us further information on $\bar{K}NN$
 - ✓ internal structure
 - ✓ $\bar{K}N$ interaction below the threshold

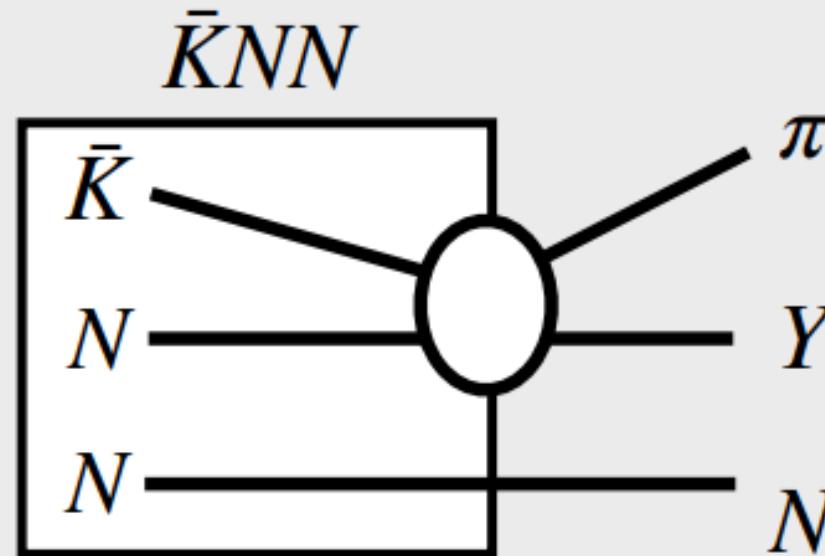
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Phys. Rev. C 88 (2013) 025204.

$$\Gamma_{YN} \ll \Gamma_{\pi YN}$$

40 times??

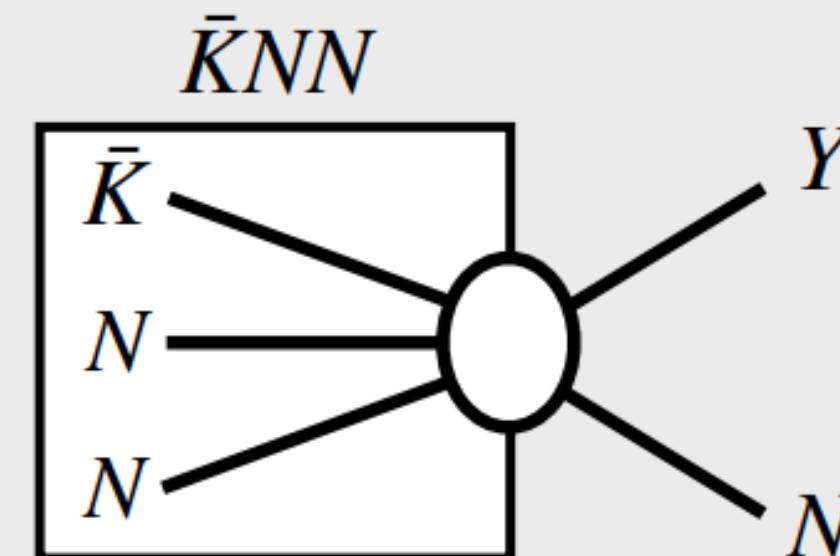


Mesonic



1N absorption

Non-mesonic



2N absorption

Mesonic Decay Modes of $\bar{K}NN$

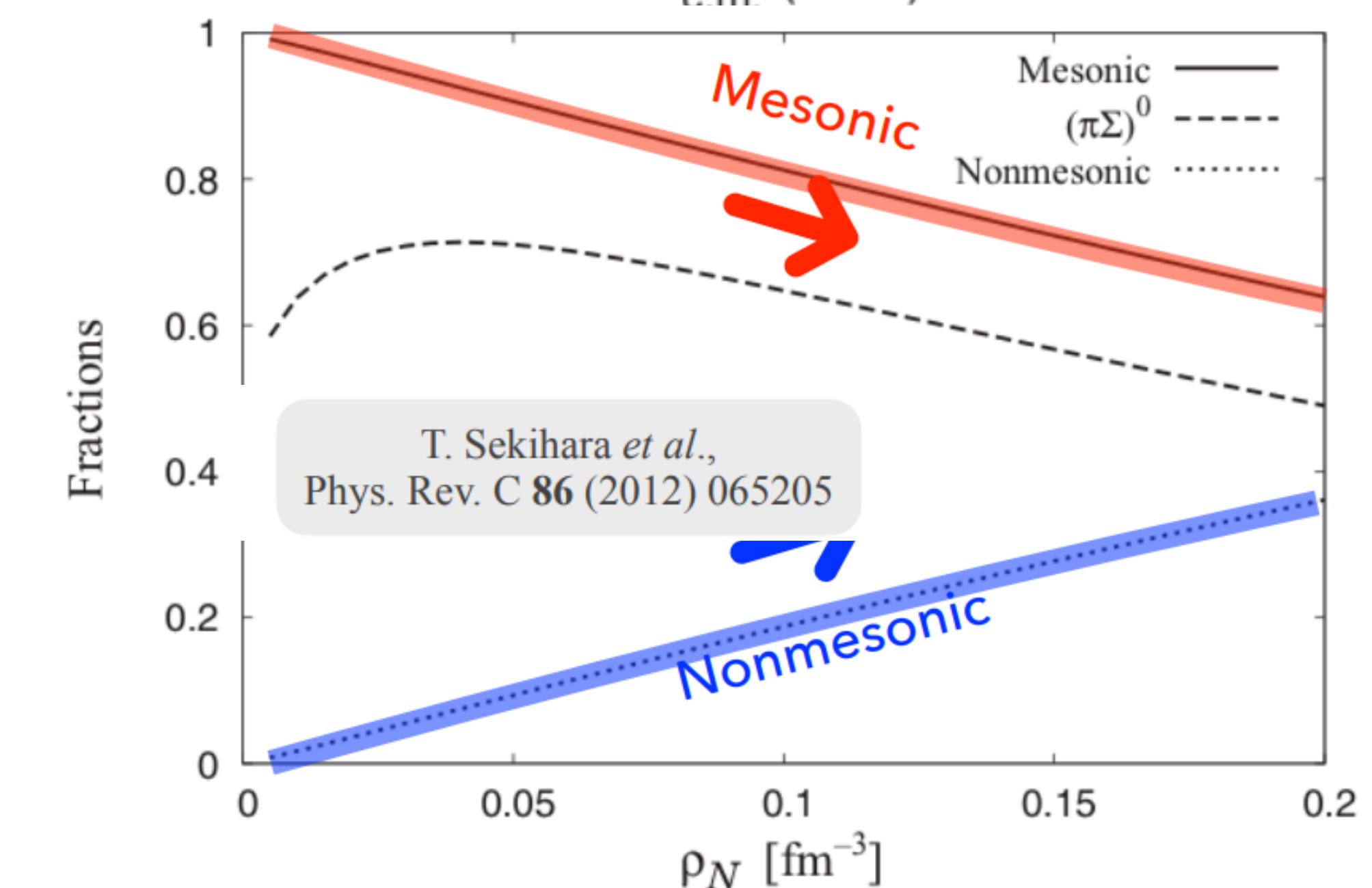
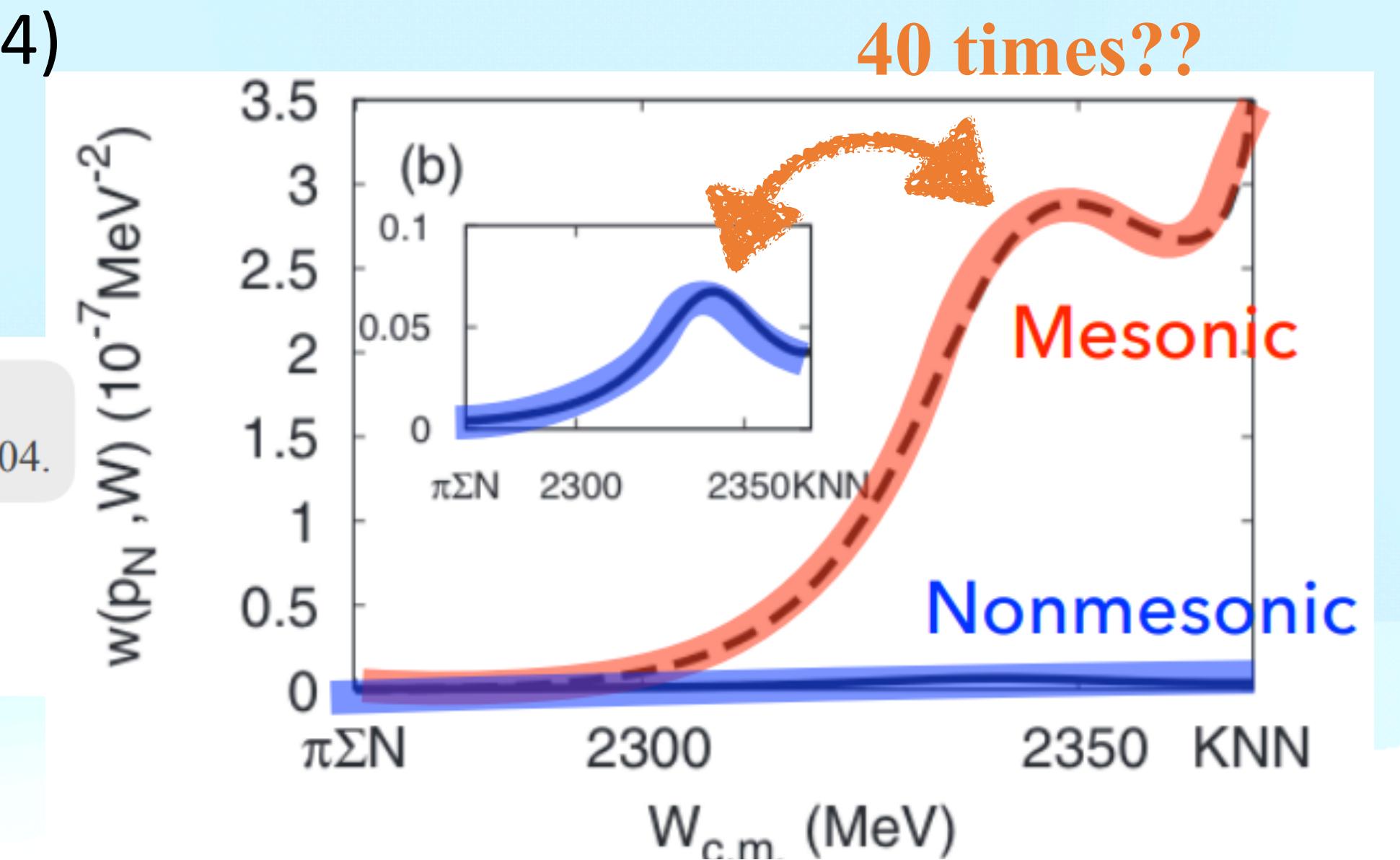
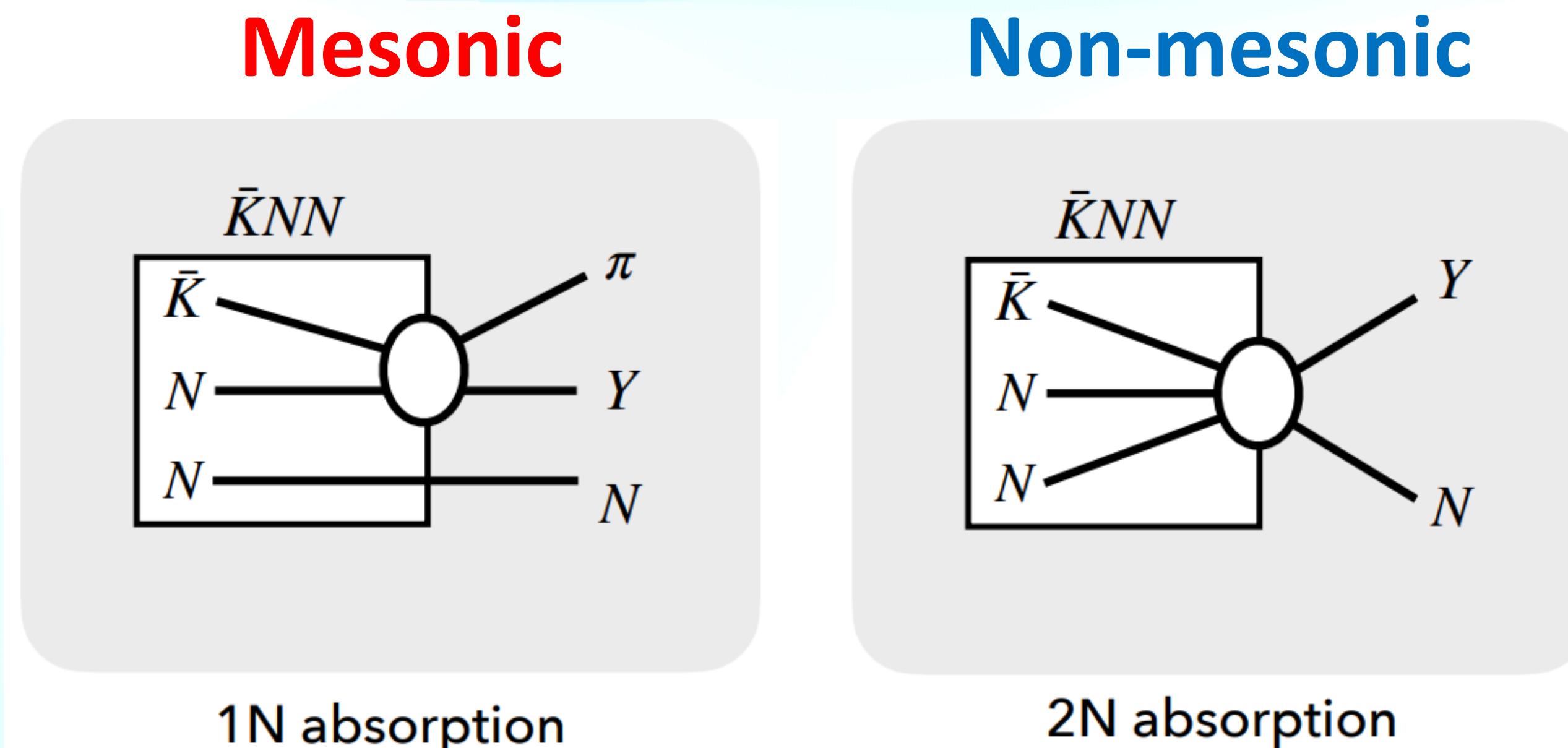
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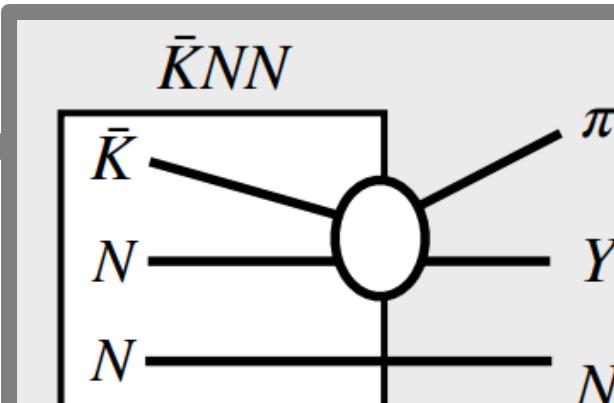
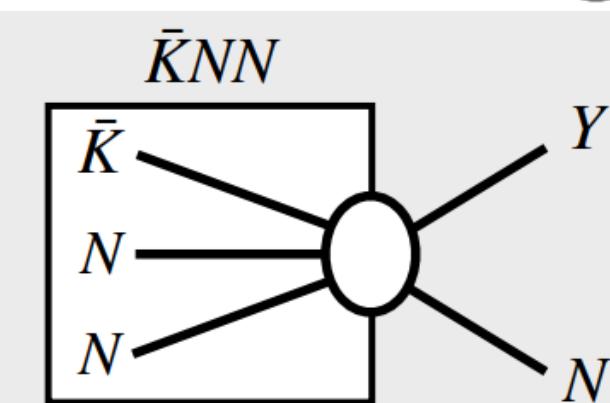
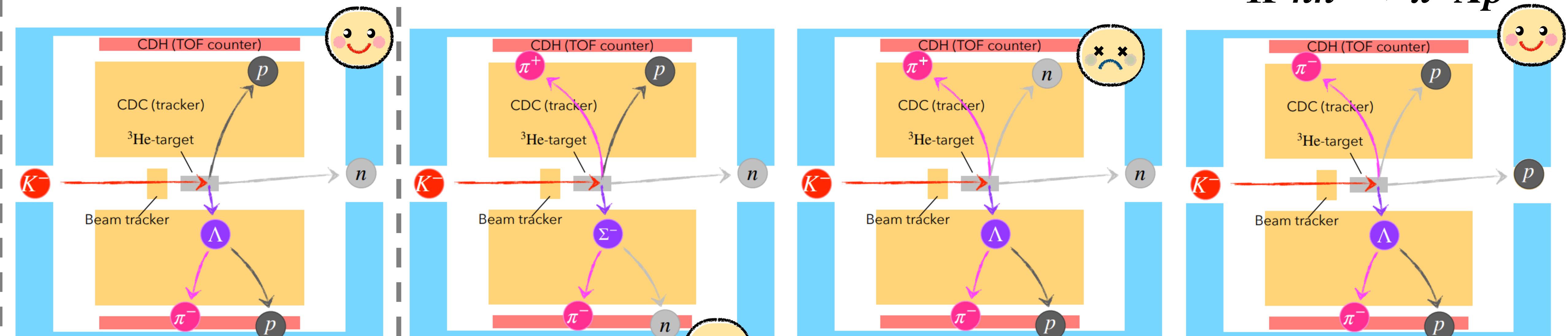
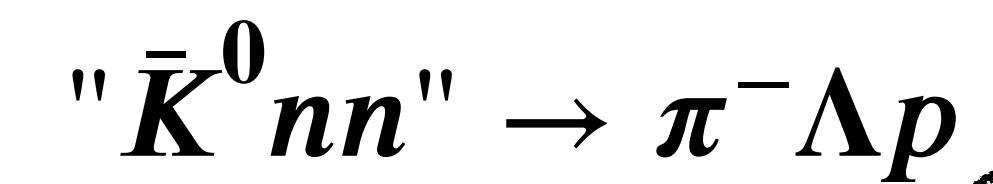
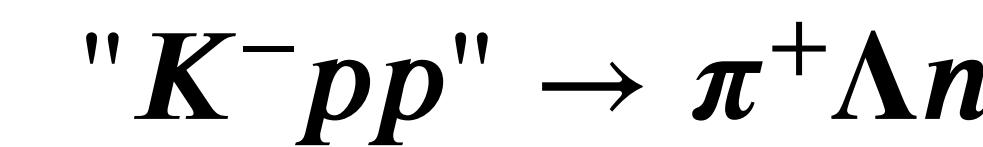
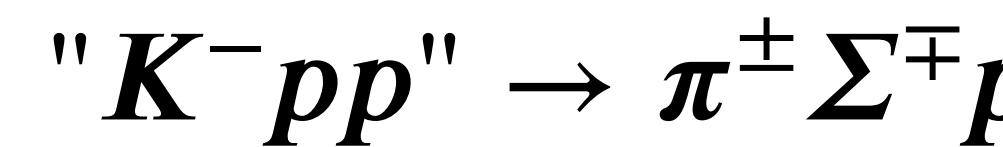
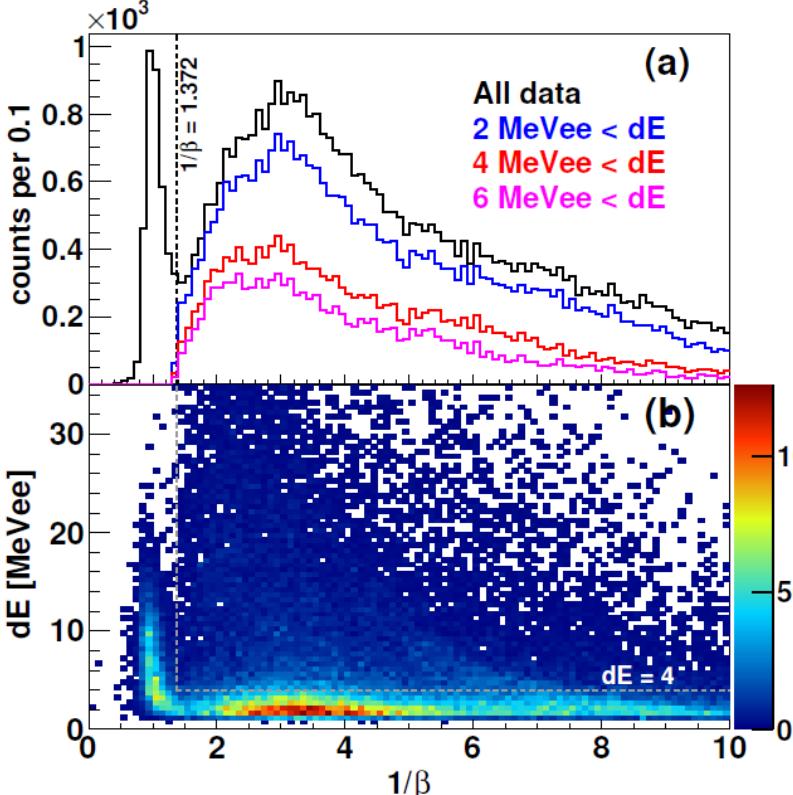
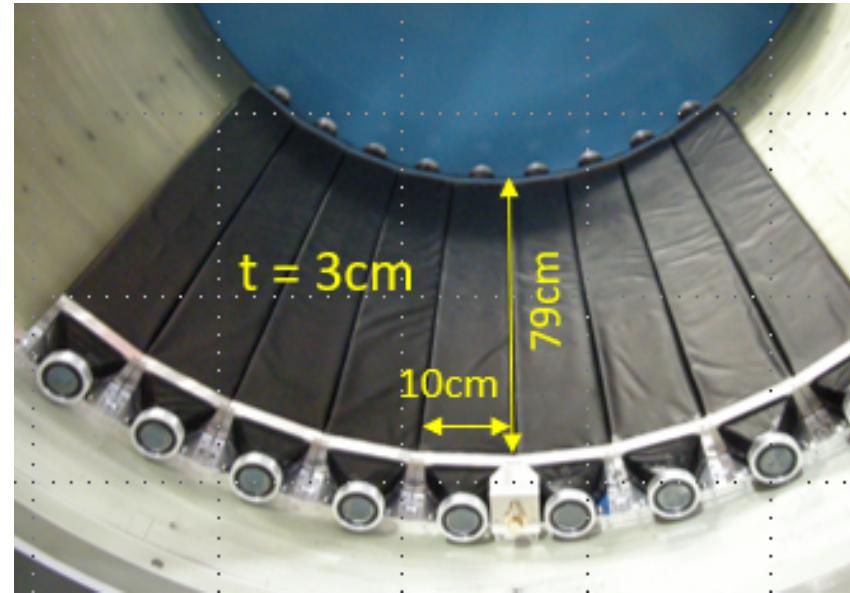
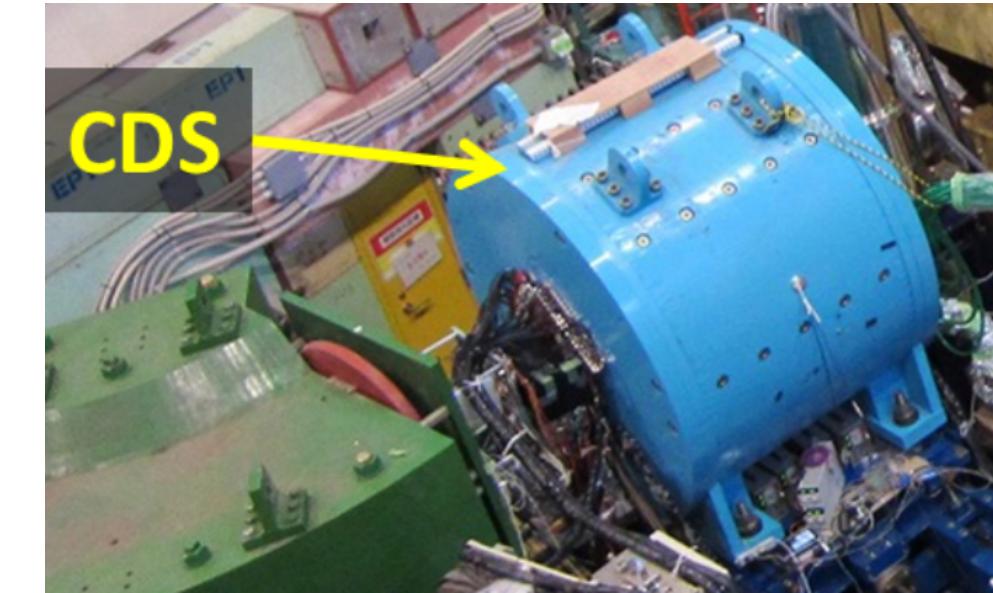
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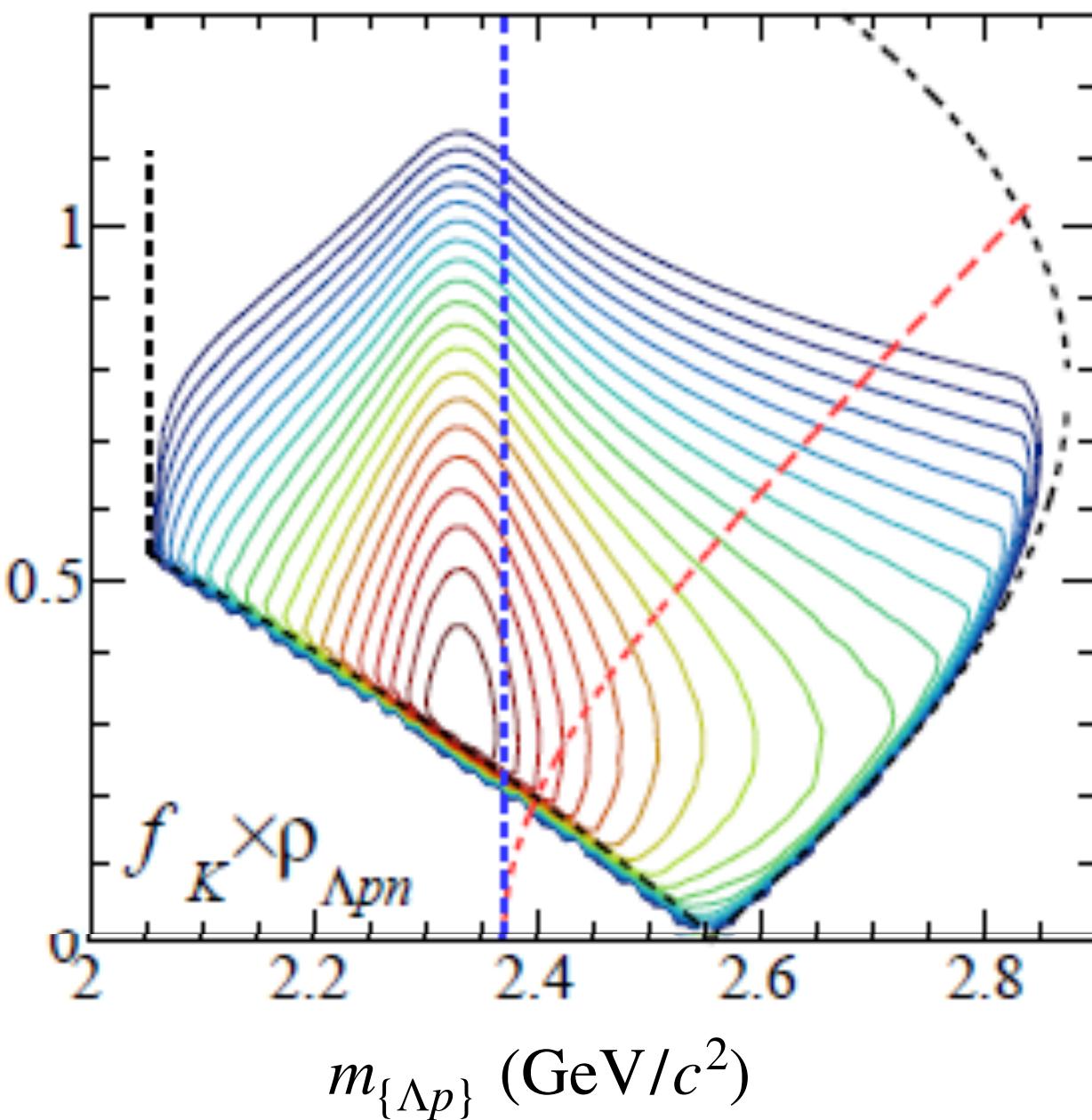


Plane Wave Impulse Approximation
Fit with PWIA

$$\sigma(M, q) \propto \rho(M, q) \times \frac{(\Gamma_{Kpp}/2)^2}{(M - M_{Kpp})^2 + (\Gamma_{Kpp}/2)^2} \times \exp\left(-\frac{q^2}{Q_{Kpp}^2}\right)$$

Phase space Momentum term from spatial integral
Energy term (BW type) from time integral

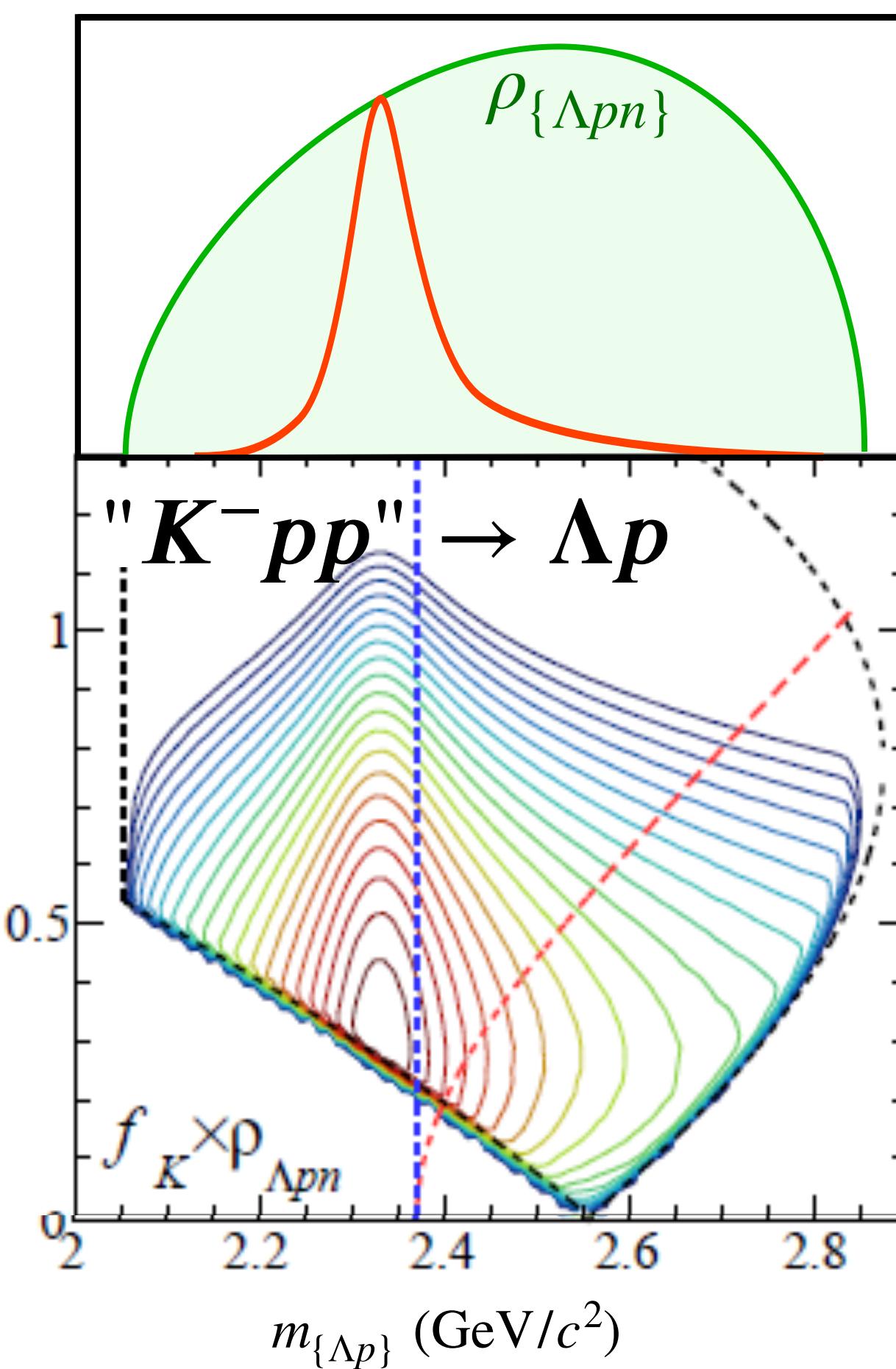
" $K^- pp$ " $\rightarrow \Lambda p$



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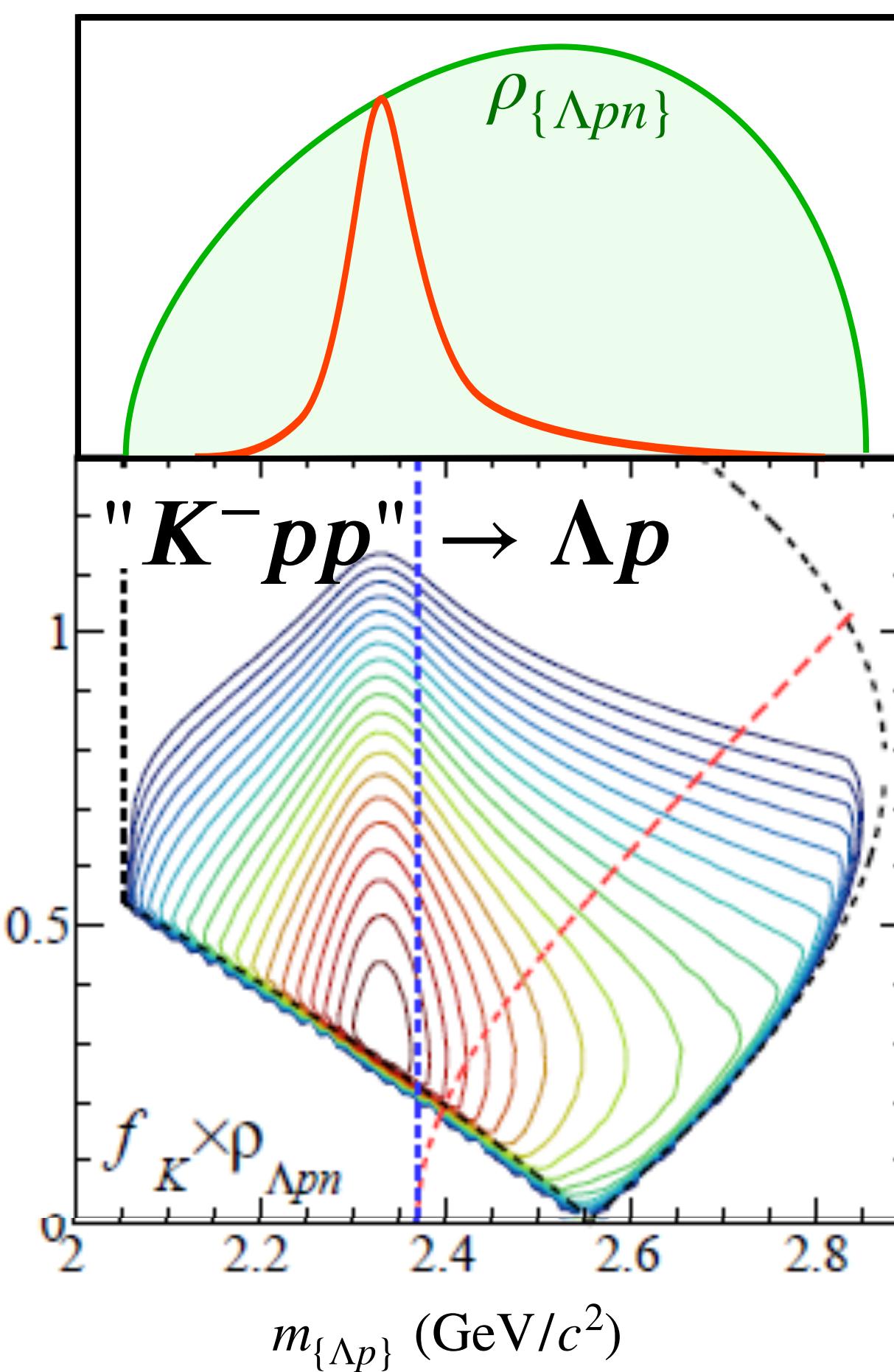


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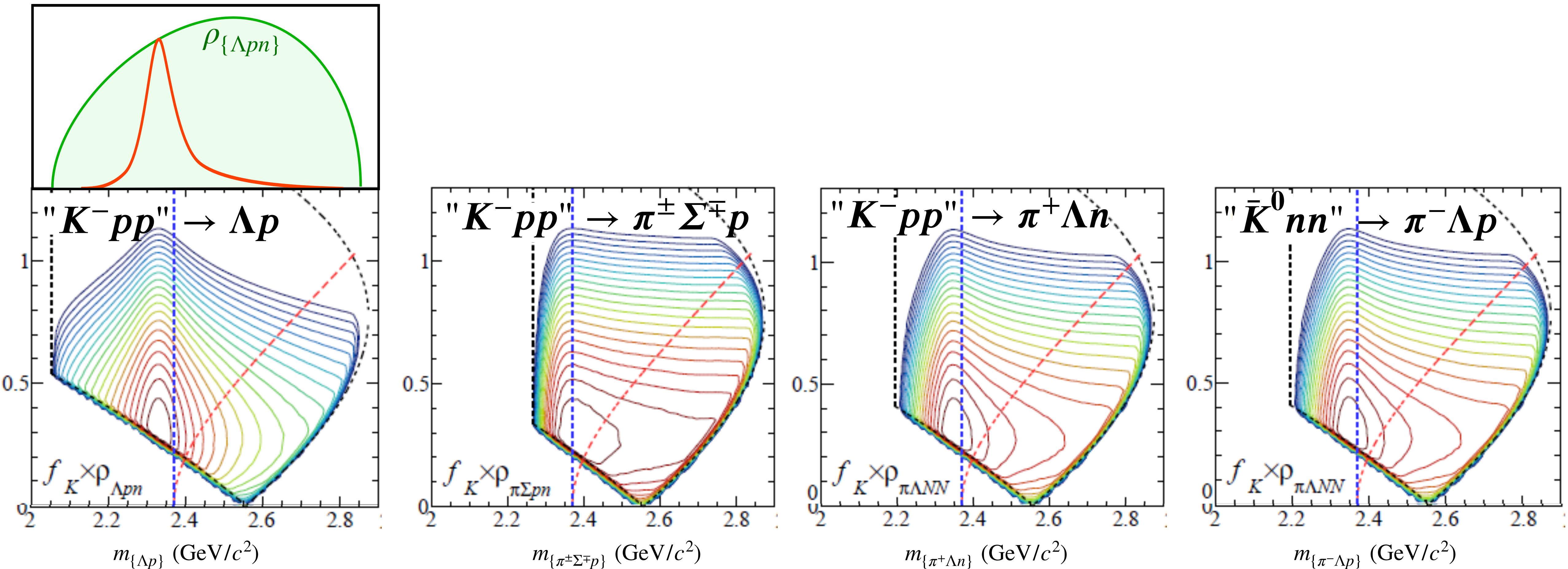
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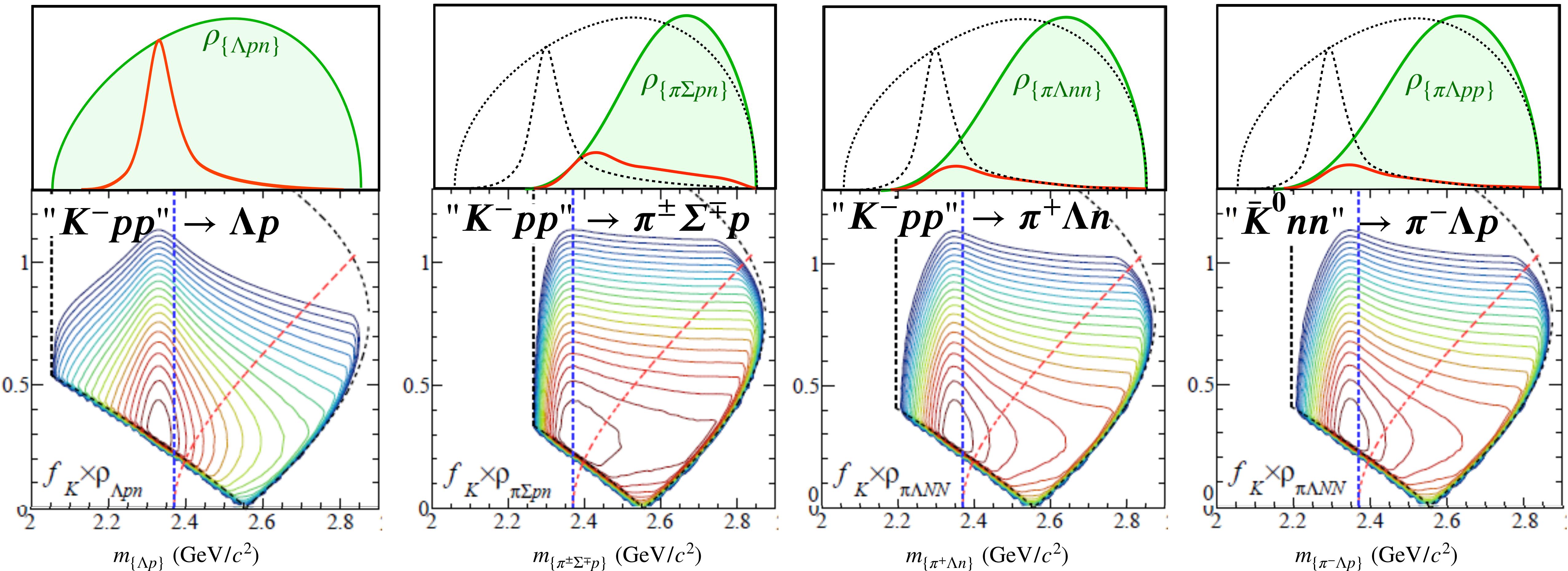
Mesonic decay suppression / spectral modification happens, due to the Phase Space Difference!



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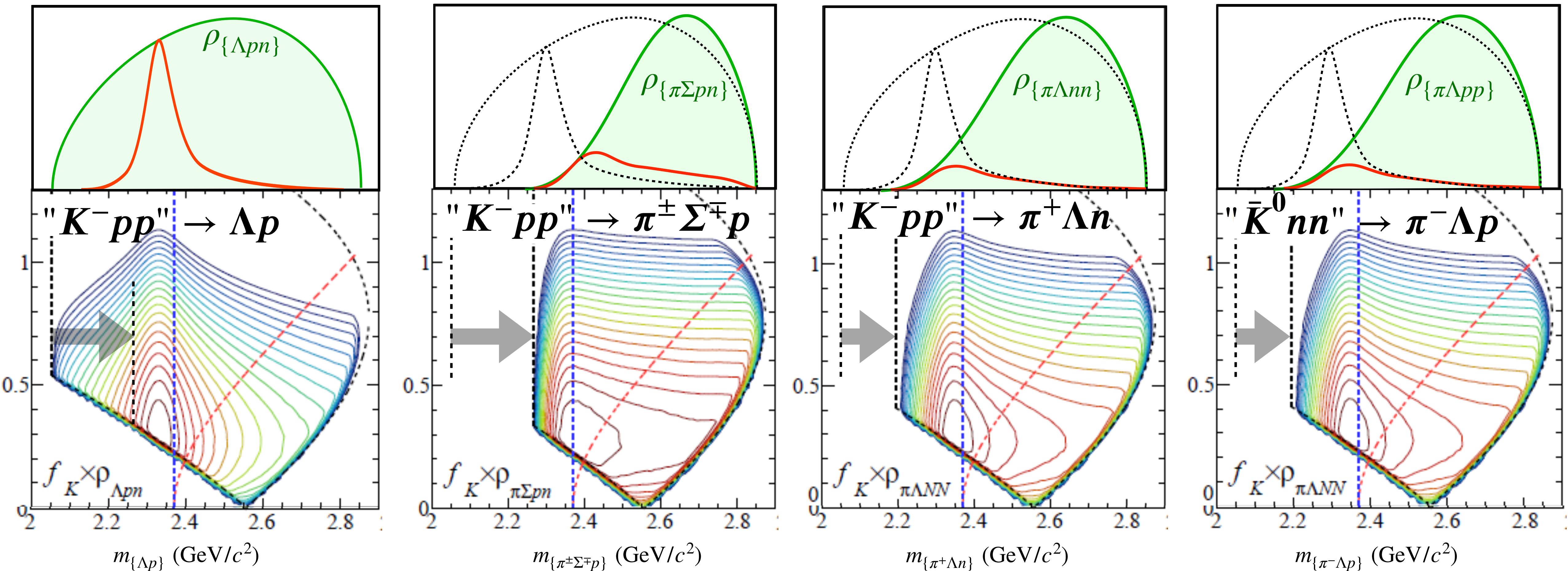


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Mesonic decay suppression / spectral modification happens, due to the **Phase Space Difference!**

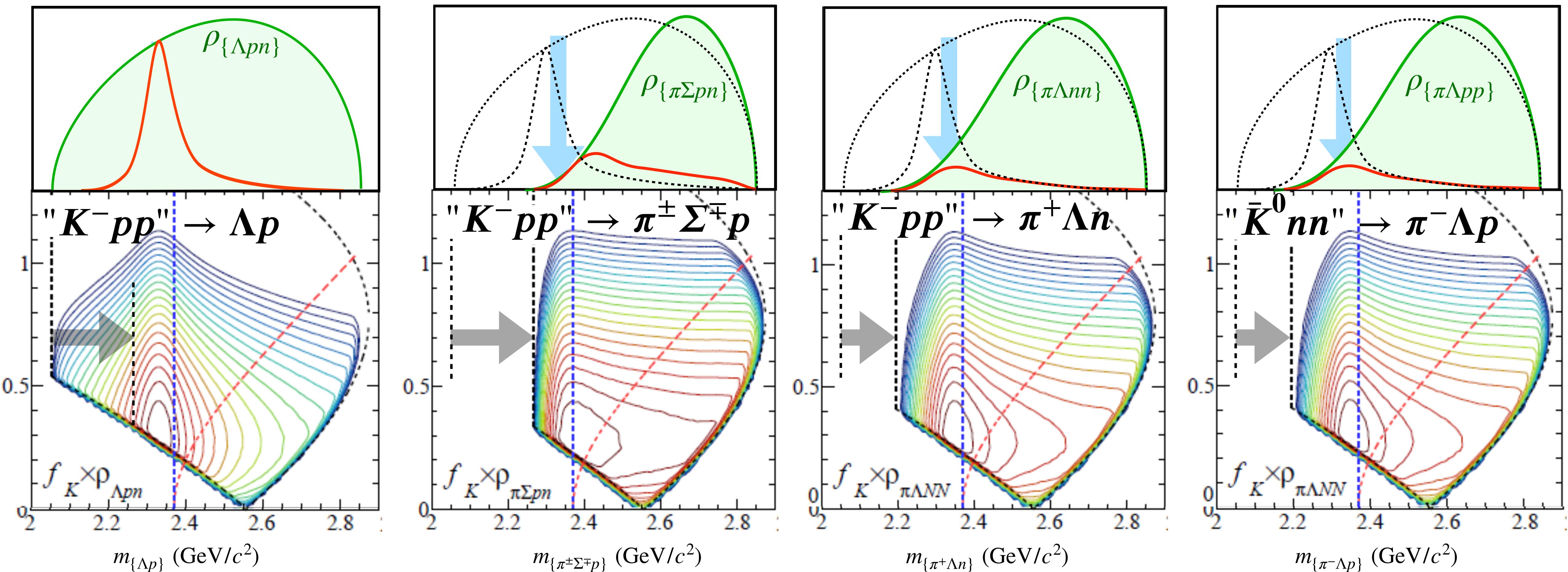
1) **mass threshold** $m_{\{\pi\Lambda p\}} \in [m_\Lambda + m_p, E_{\{K^3\text{He}\}}^* - m_n] \rightarrow [m_\pi + m_Y + m_N, E_{\{K^3\text{He}\}}^* - m_N]$



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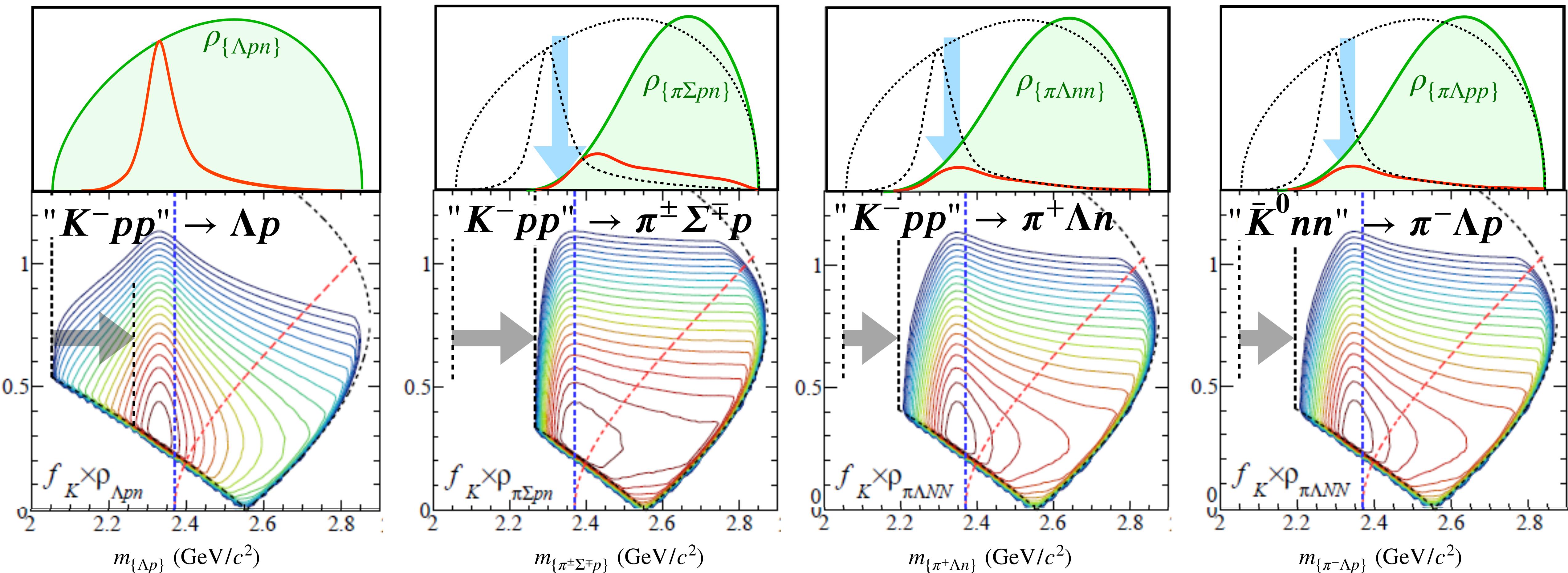


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Thus, the peak structure is difficult to be detected in the $K^- + {}^3\text{He} \rightarrow \{\pi YN\} + N'$ reaction dynamics.

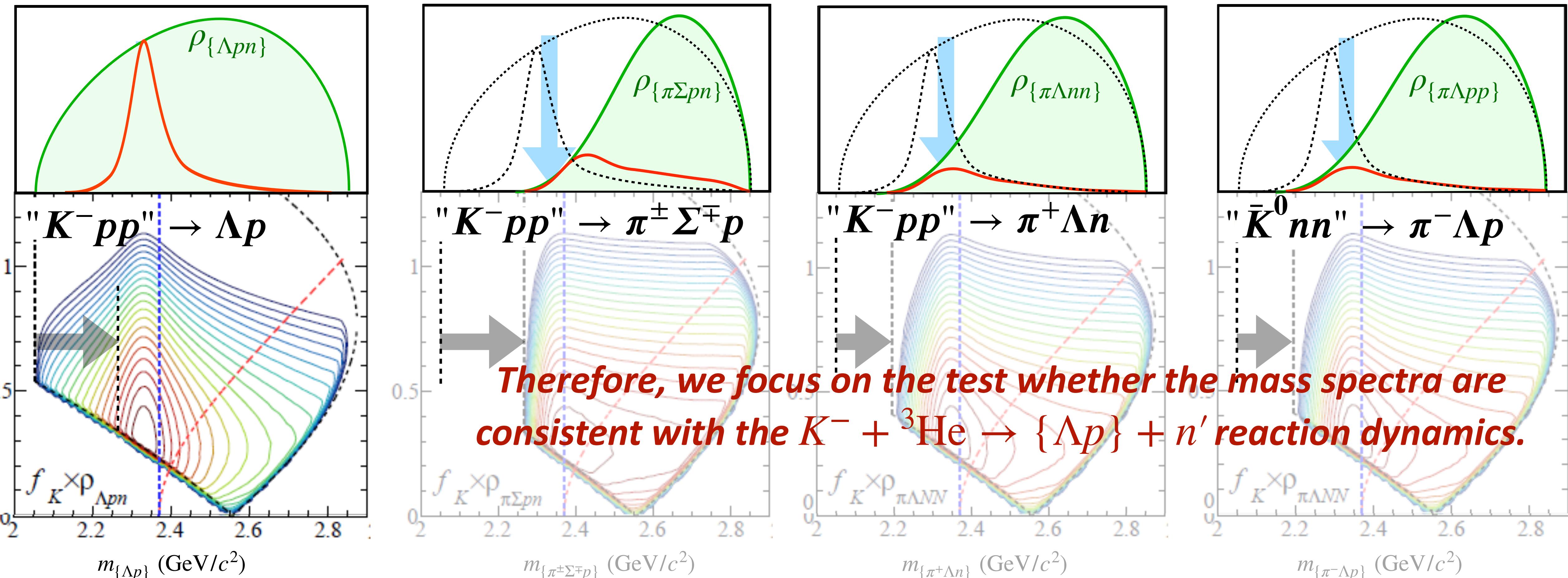


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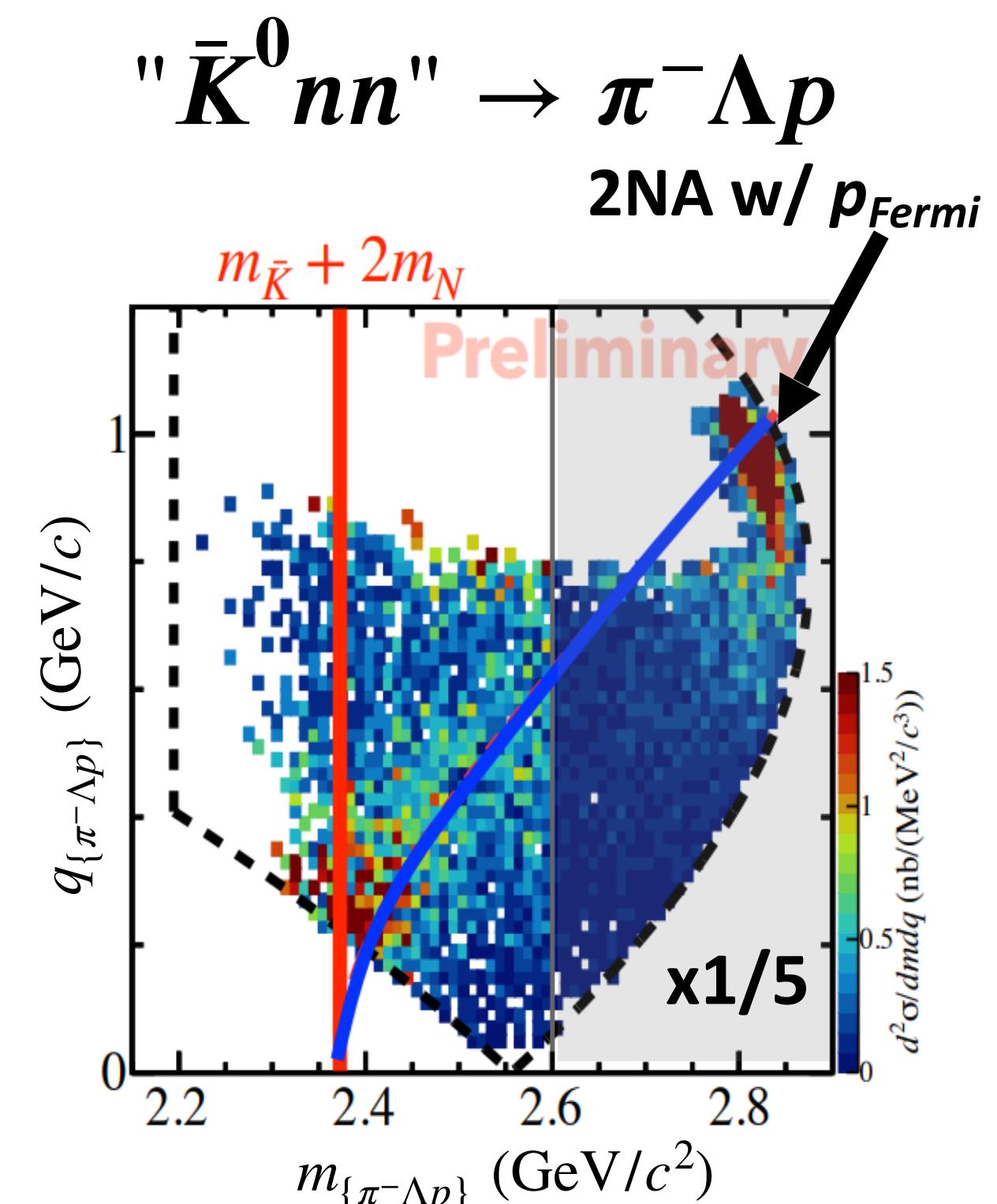
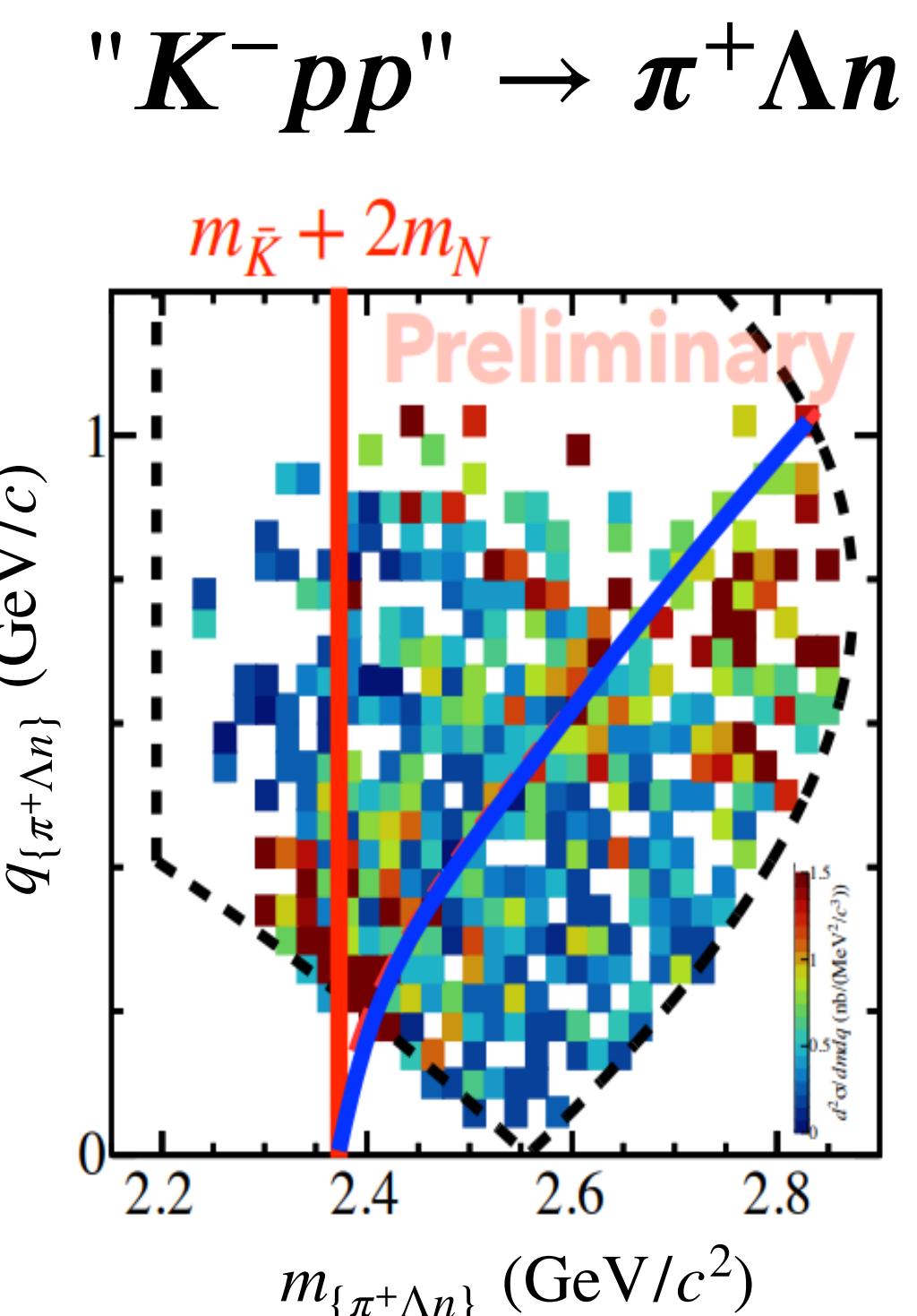
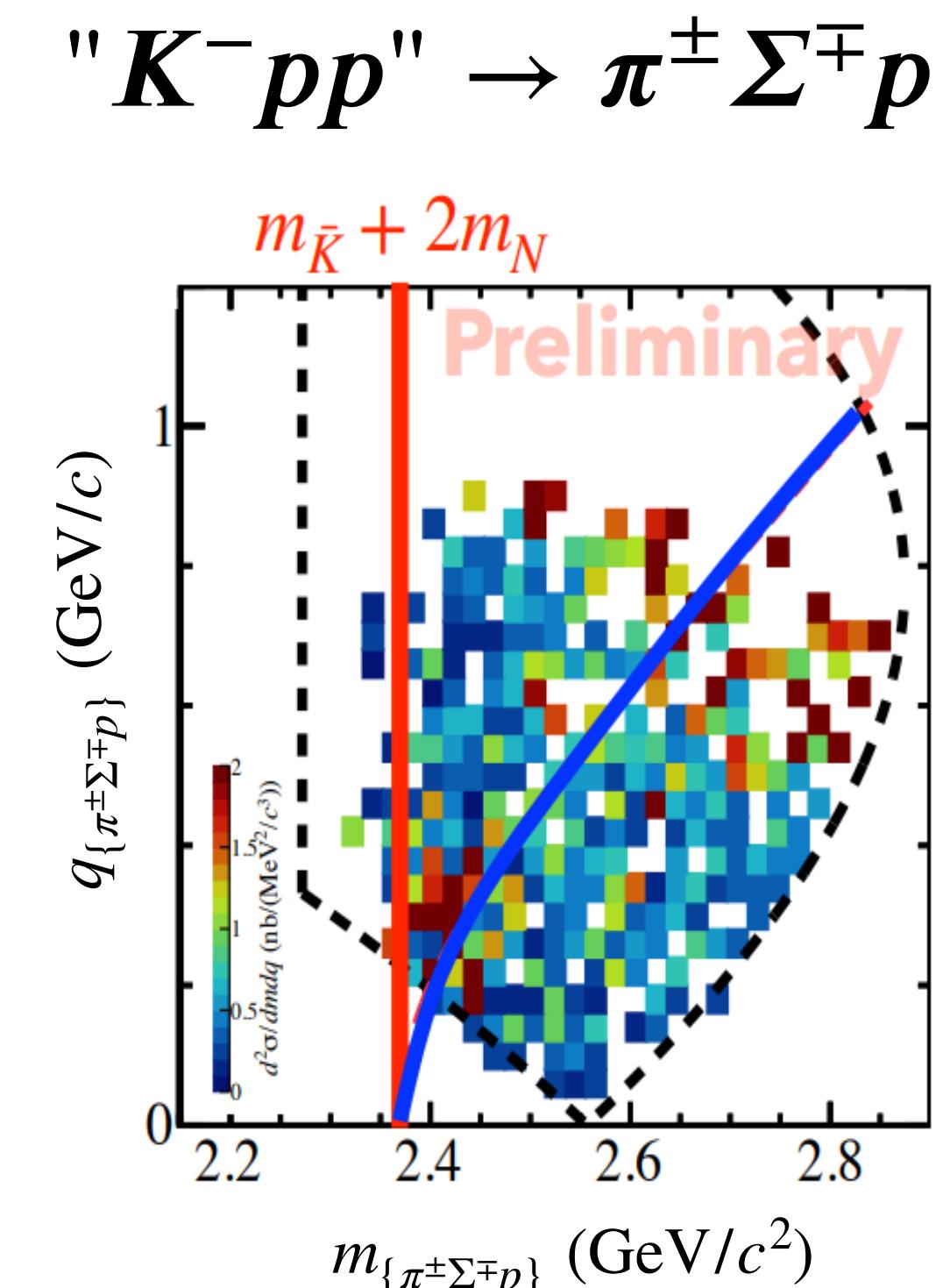
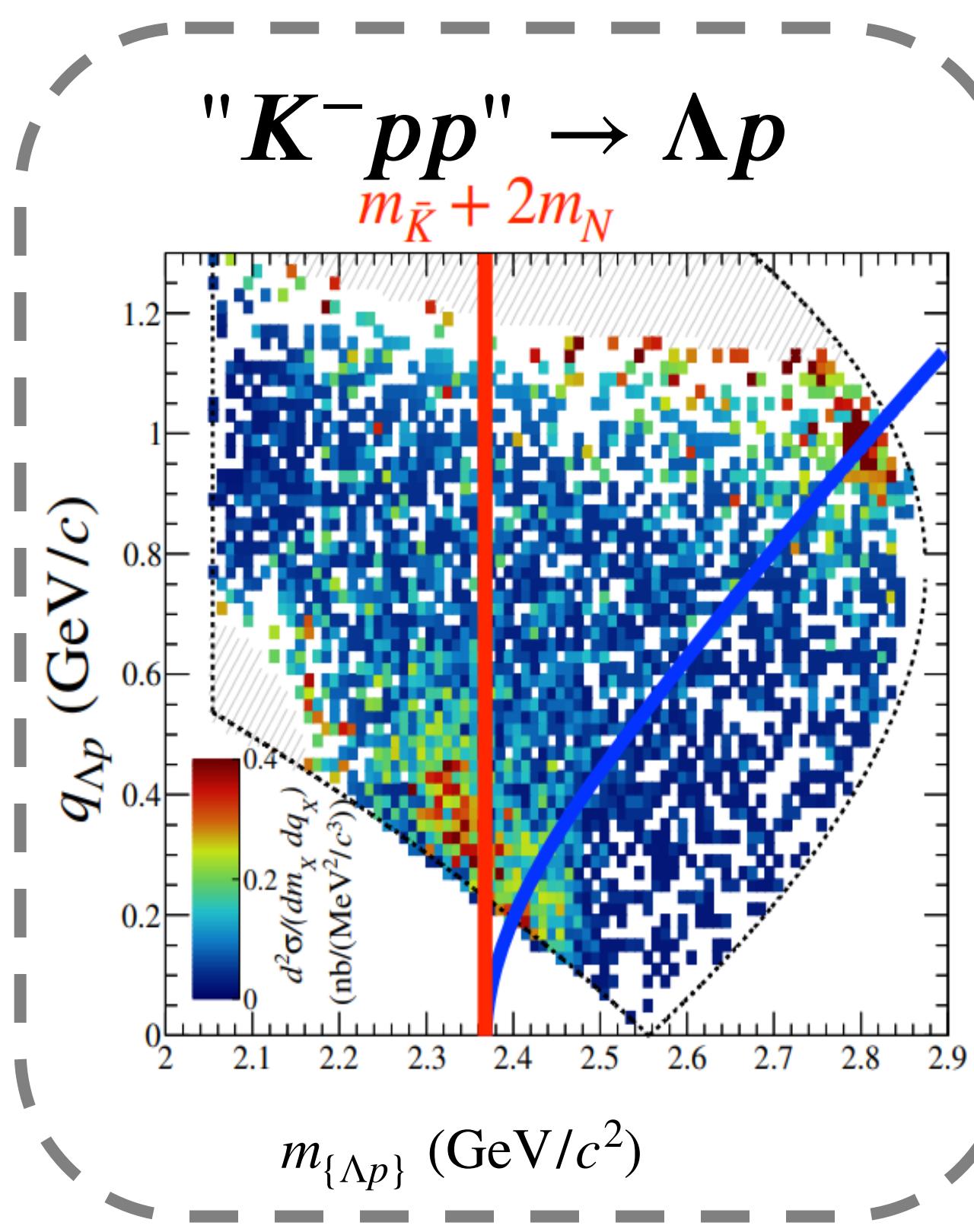
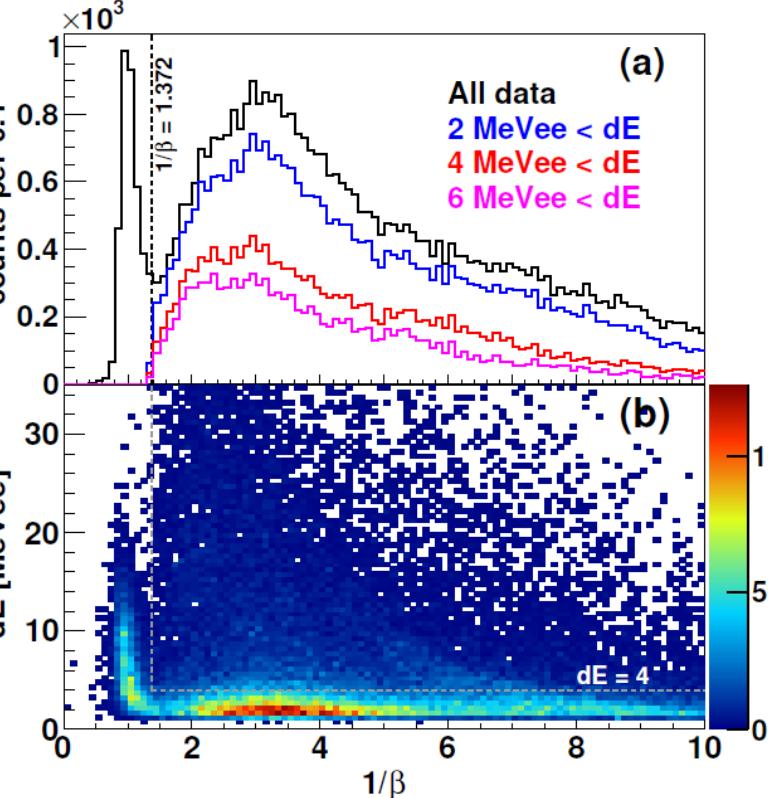
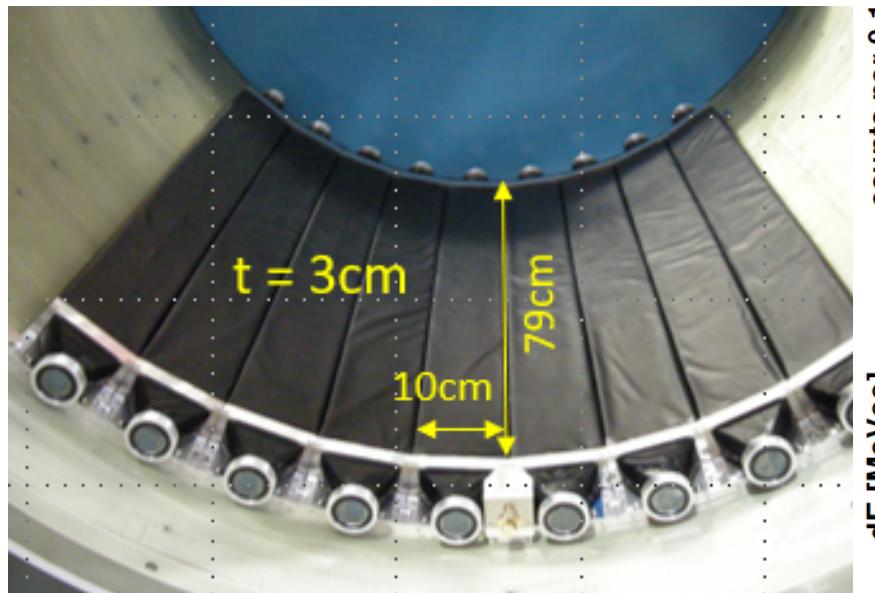
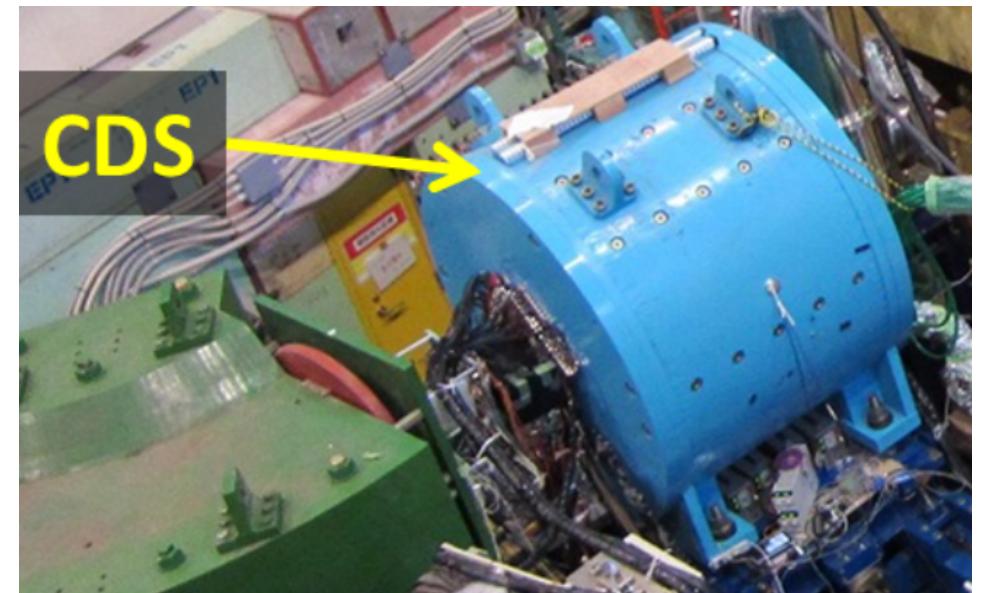
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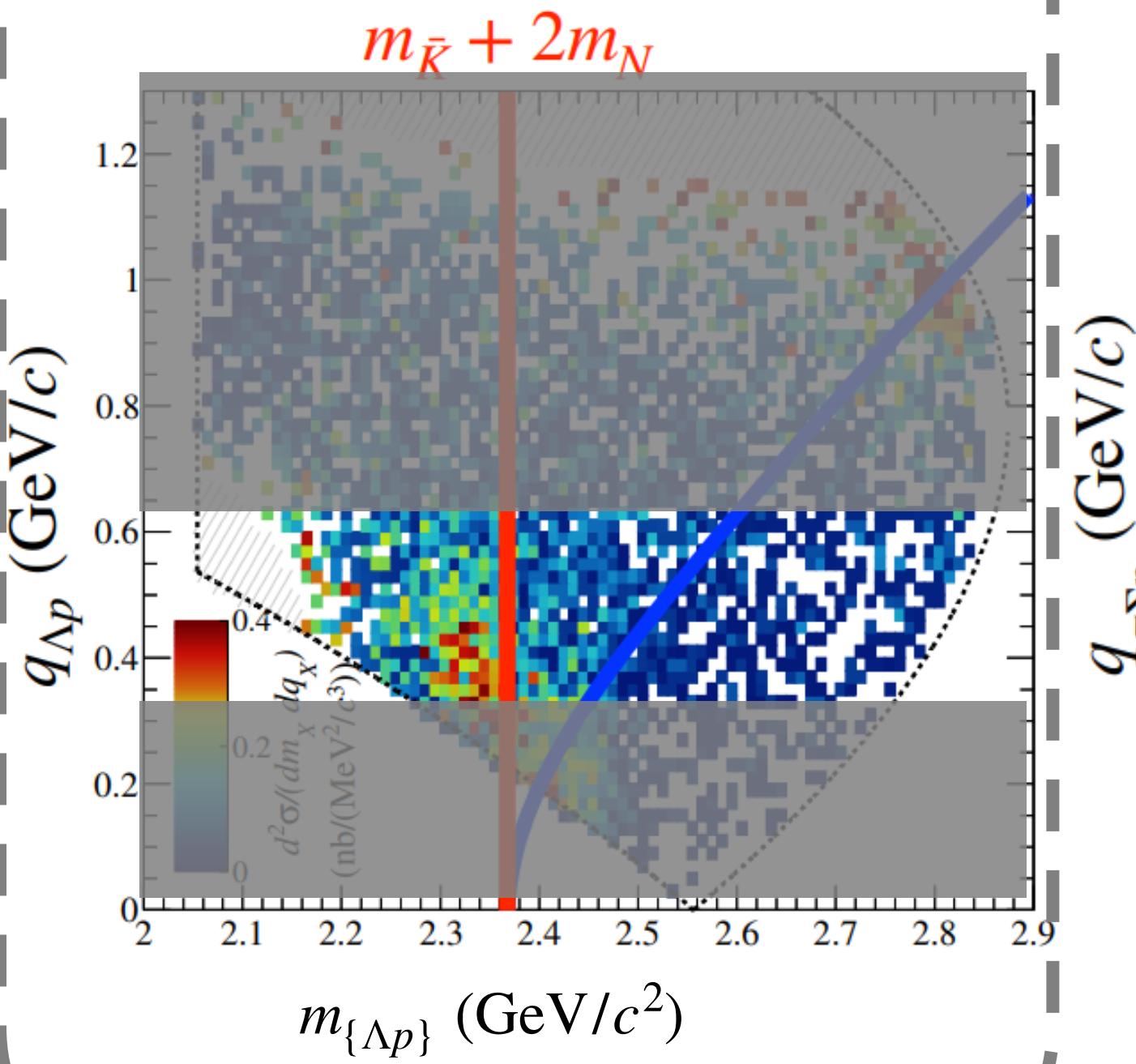
Phase space

Momentum term from spatial integral

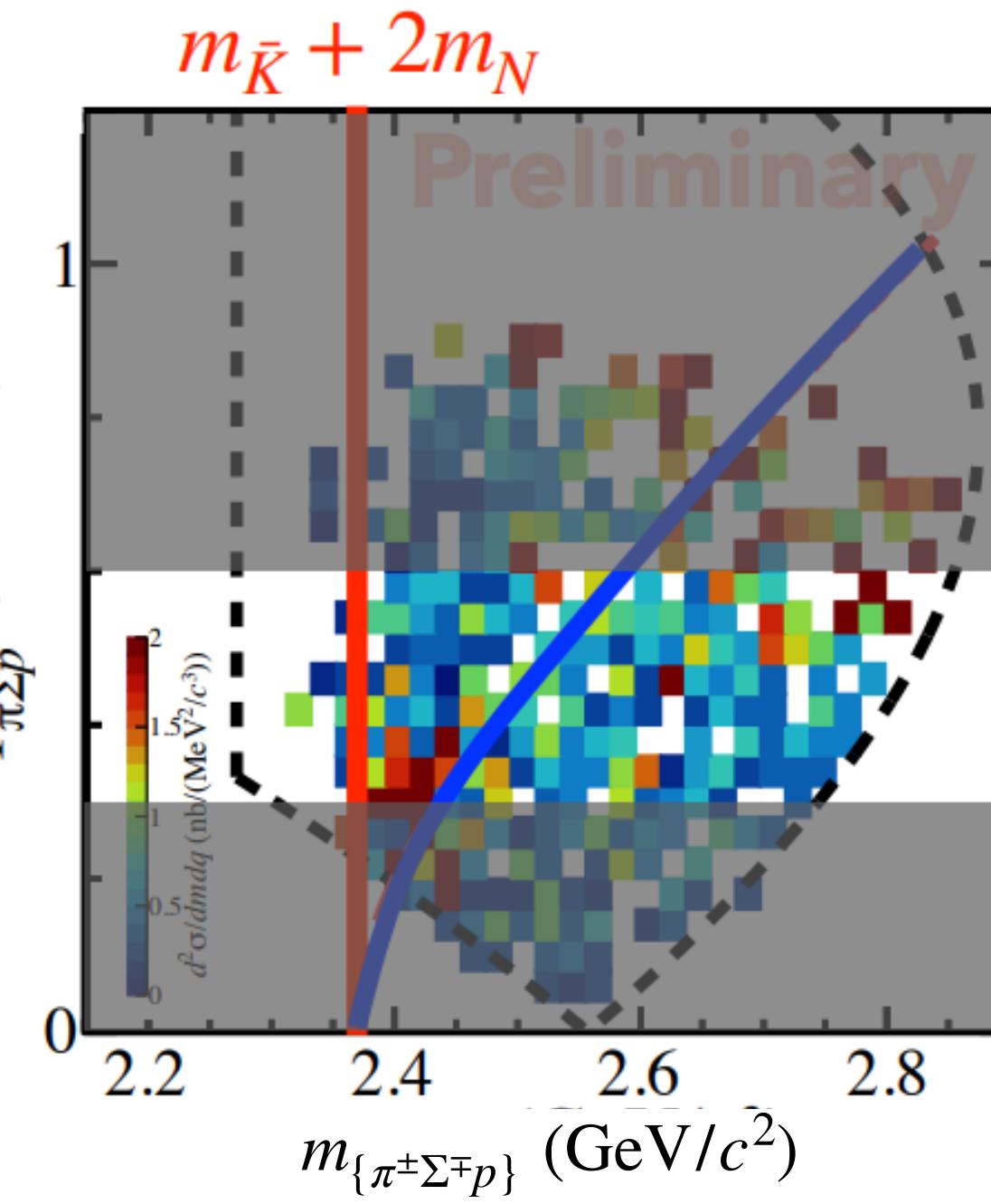
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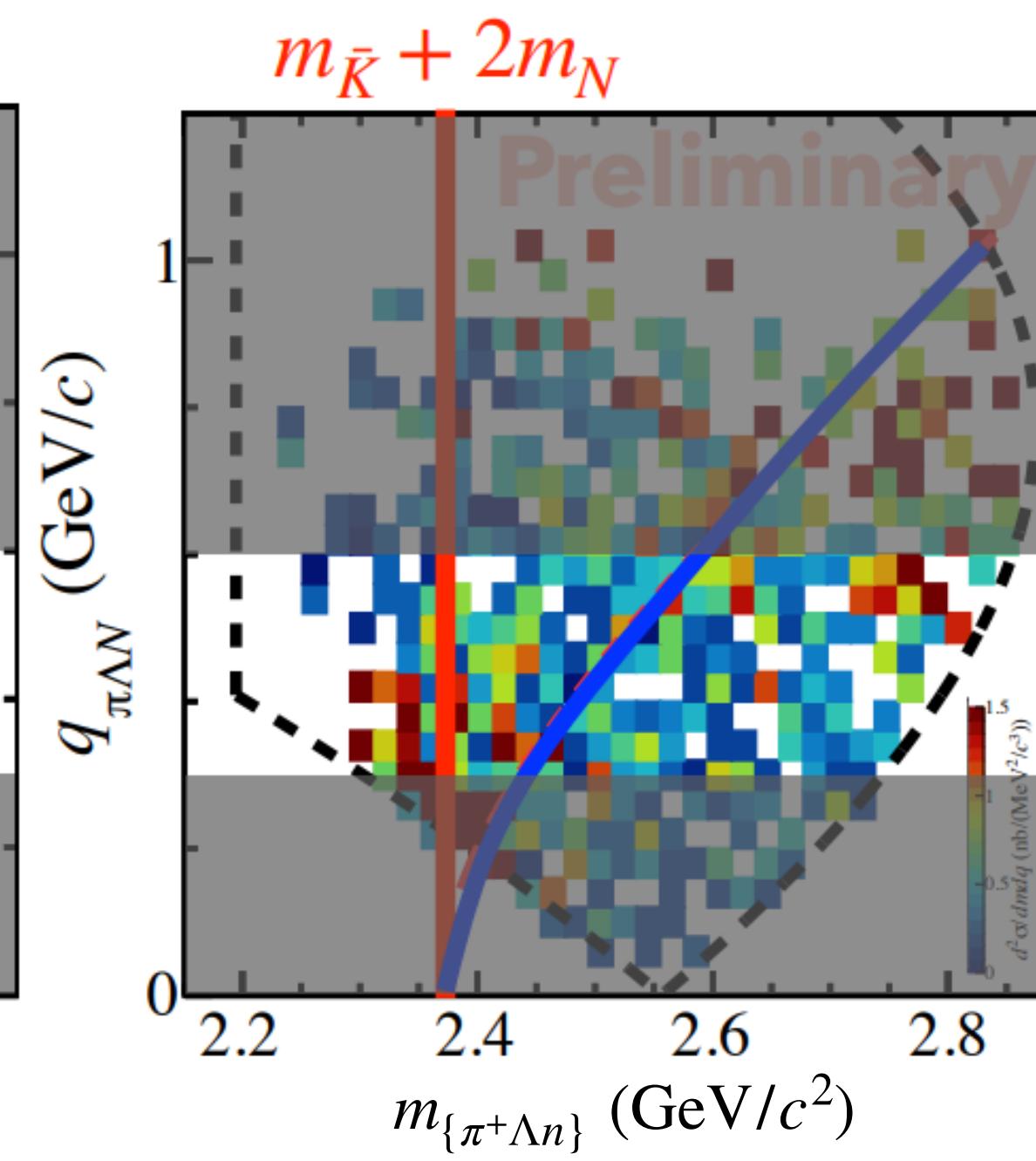
" K^-pp " $\rightarrow \Lambda p$



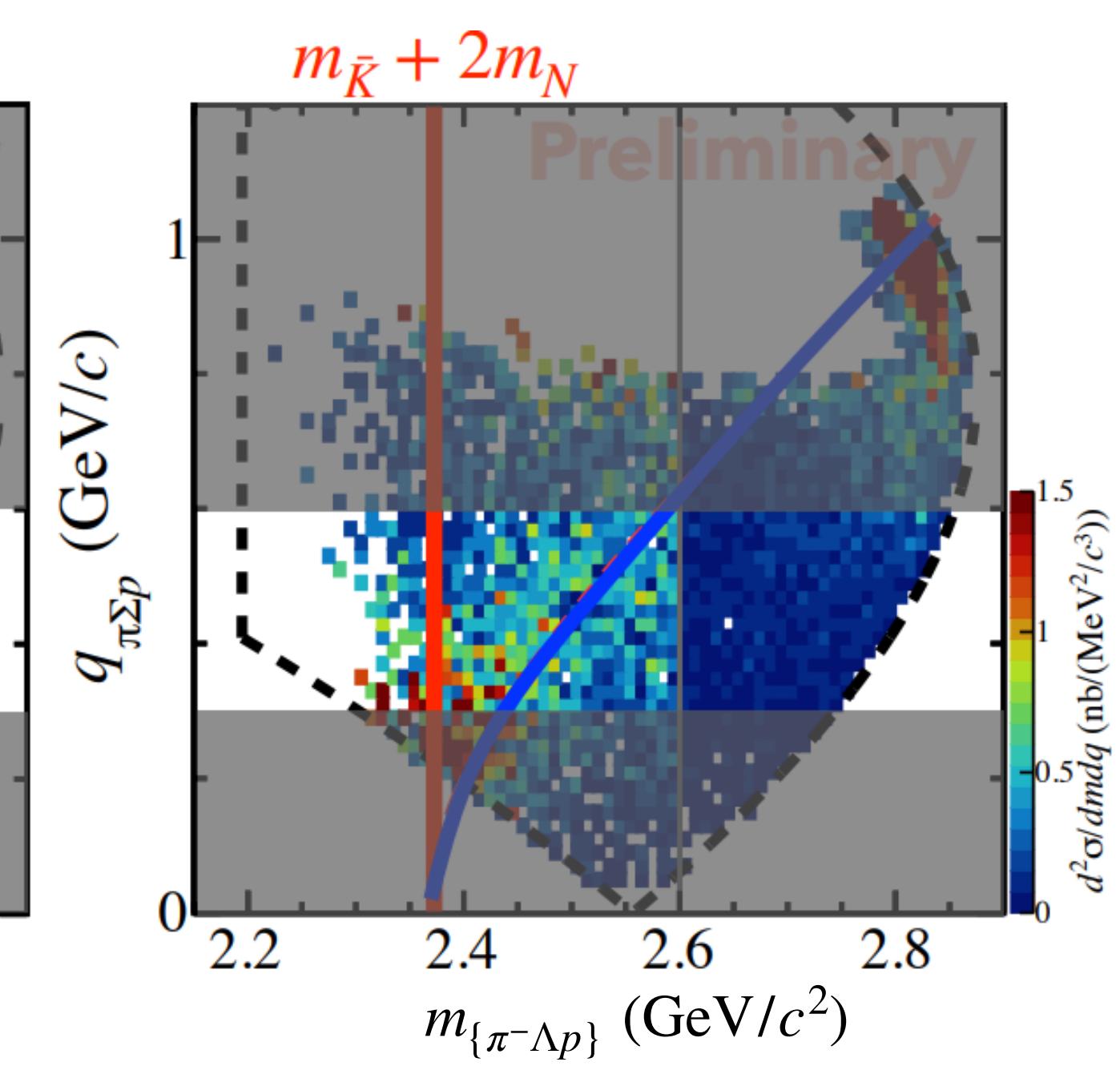
" K^-pp " $\rightarrow \pi^\pm \Sigma^\mp p$



" K^-pp " $\rightarrow \pi^+ \Lambda n$



" $\bar{K}^0 nn$ " $\rightarrow \pi^- \Lambda p$



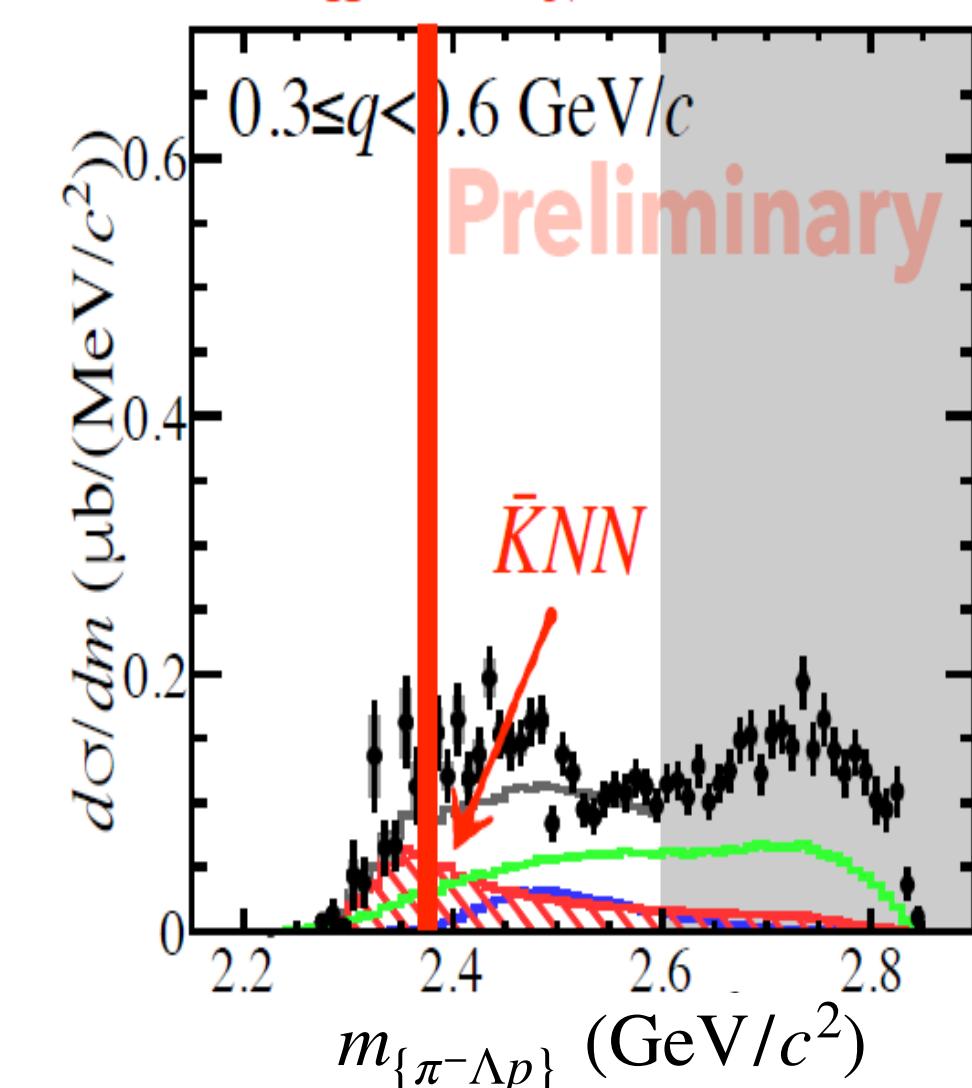
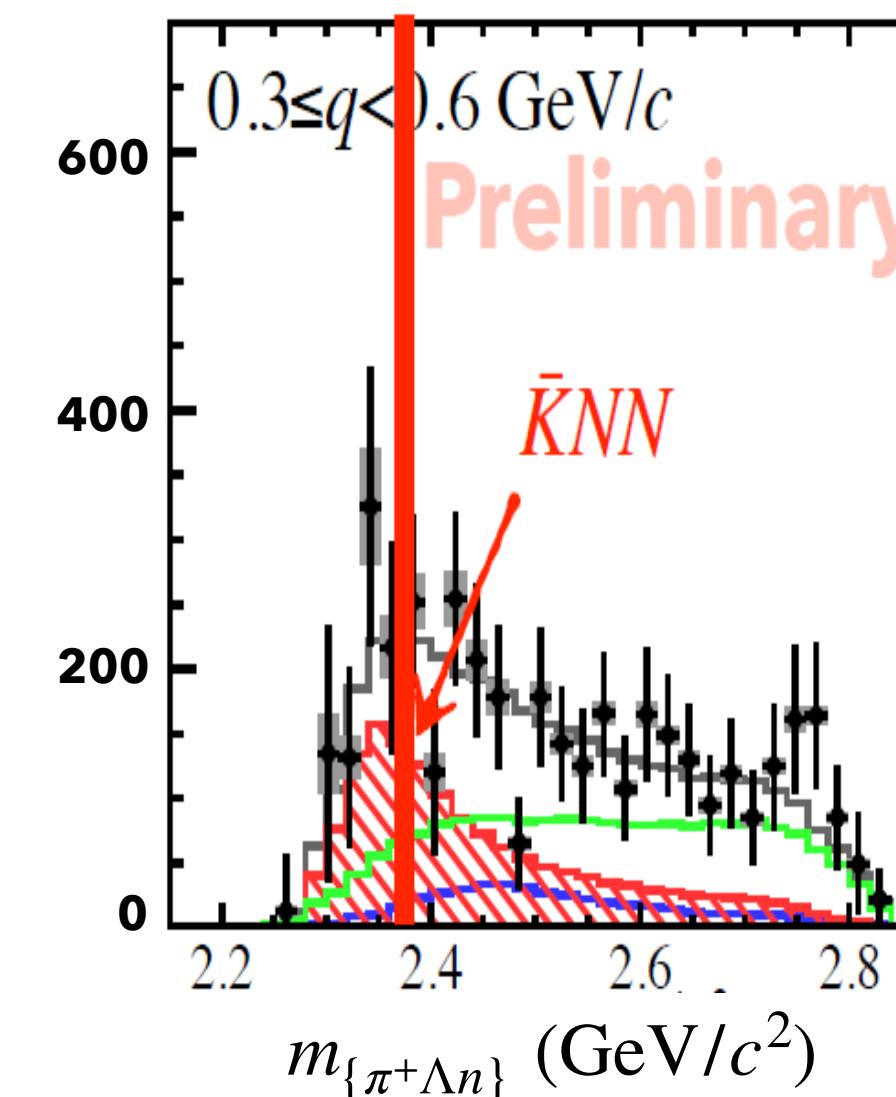
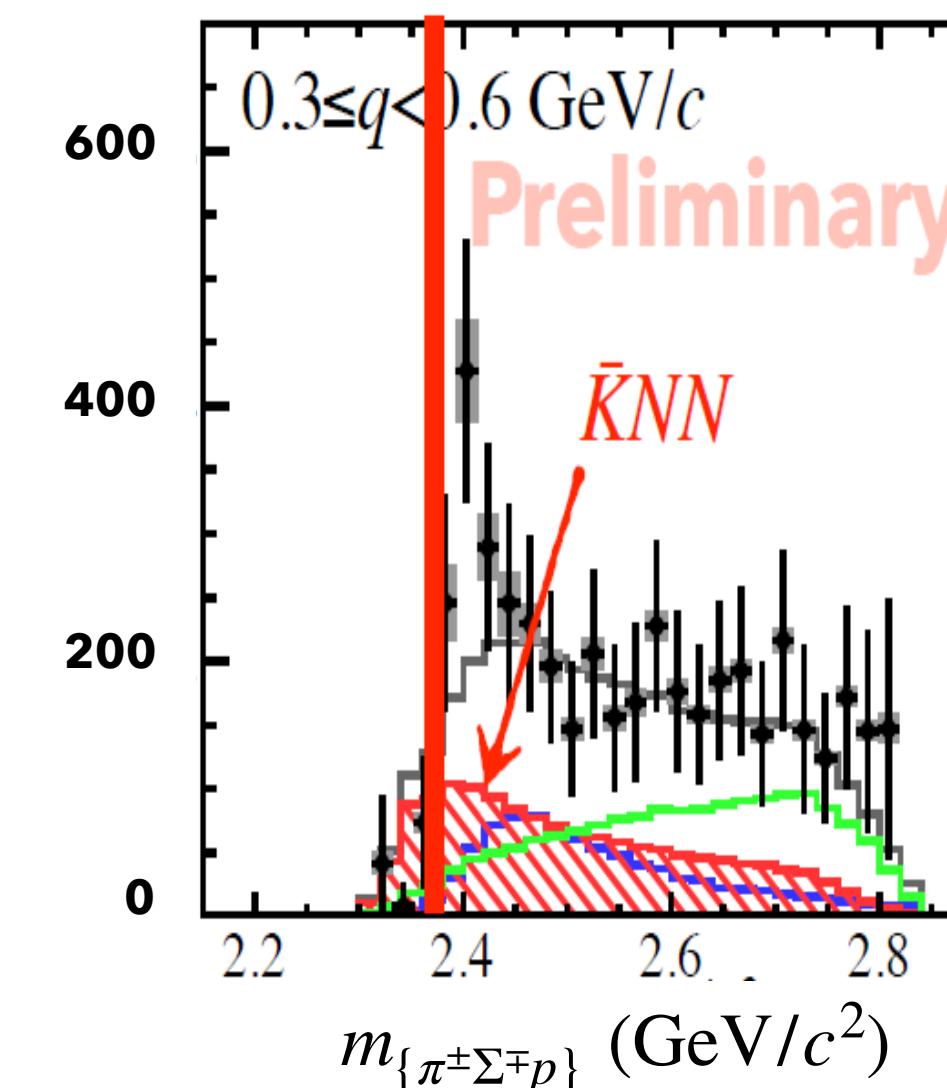
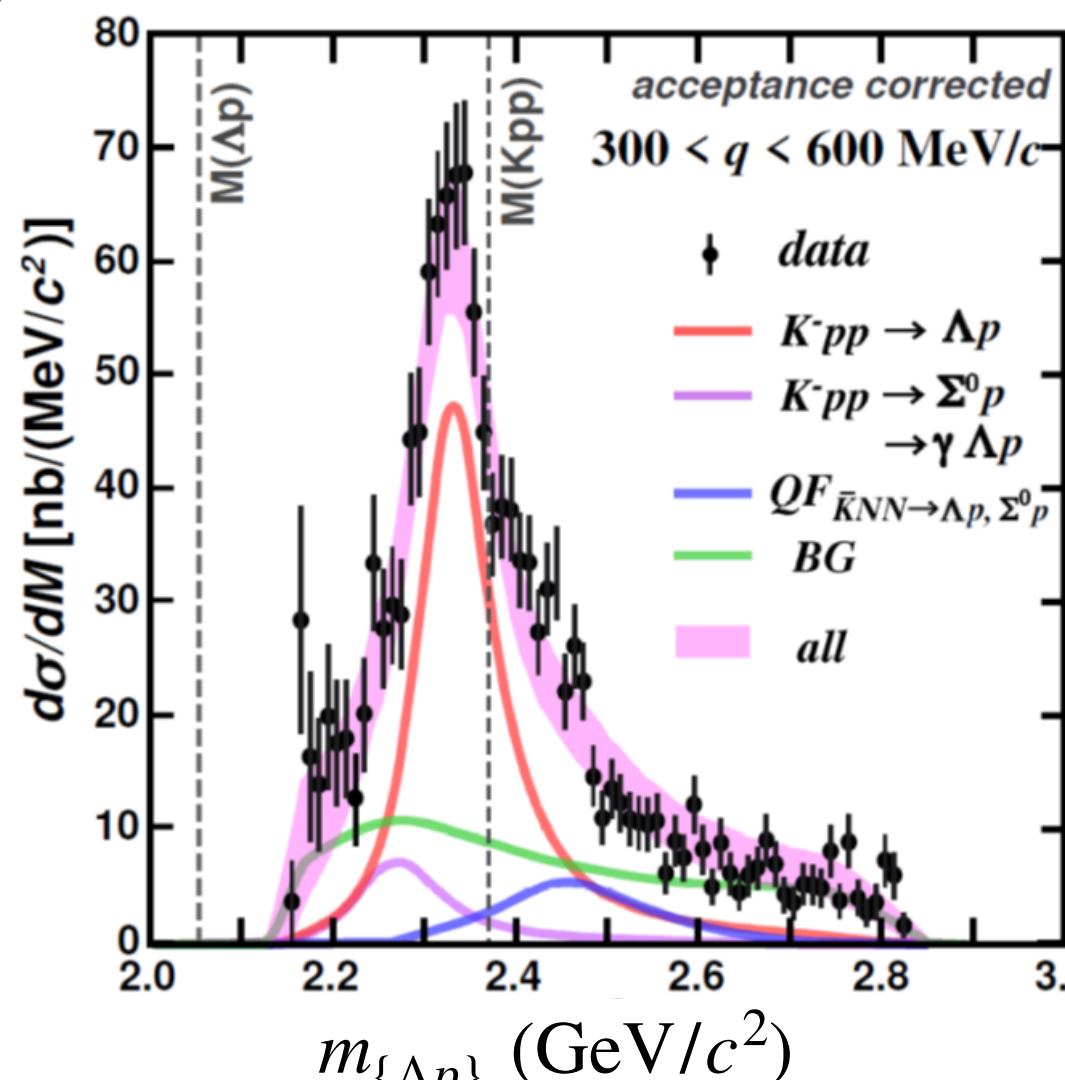
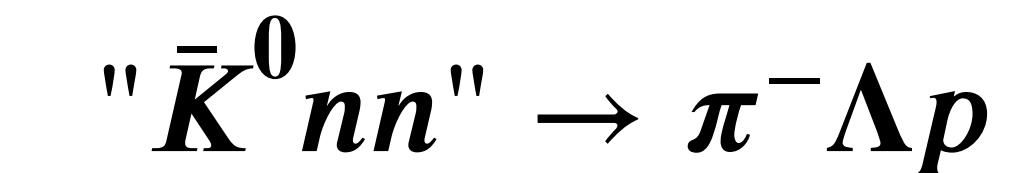
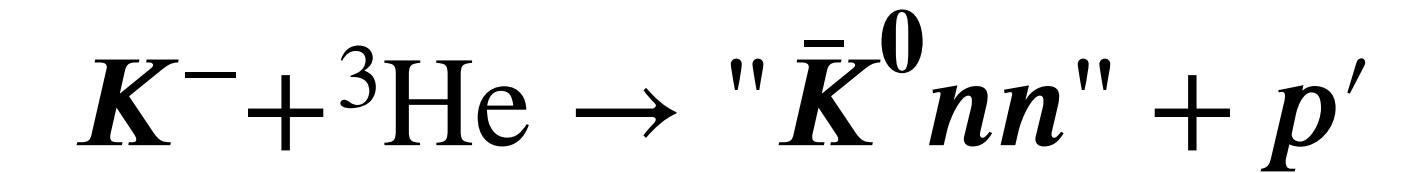
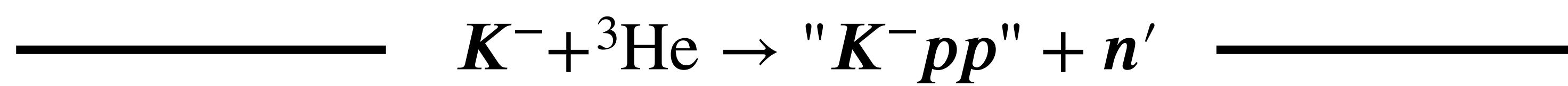
Fit the 1D spectra in $0.3 < q < 0.6$ with the same model func.

Mesonic Decay Analysis (with the E15 Data)

Plane Wave Impulse Approximation
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Phase space Momentum term from spatial integral
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With the model func., the spectra are consistently explained.

Mesonic Decay (suppressed by the 4-body phase space)

" K^-pp " → Λp

$$\sigma_{\bar{K}NN}^{tot} \times Br (\mu b) = \\ 9.3 \pm 0.8^{+1.4}_{-1.0} [\text{all}] \\ 5.5 \pm 0.5^{+0.8}_{-0.6} [<\mathcal{M}(\text{KNN})]$$

" K^-pp " → $\pi^+\Sigma^-p$

$$\sigma_{\bar{K}NN}^{tot} \times Br (\mu b) = \\ 38 \pm 3 \pm 3 [\text{all}] \\ 3.2 \pm 0.2 \pm 0.2 [<\mathcal{M}(\text{KNN})]$$

" K^-pp " → $\pi^+\Lambda n$

$$\sigma_{\bar{K}NN}^{tot} \times Br (\mu b) = \\ 62 \pm 11 \pm 9 [\text{all}] \\ 15.5 \pm 2.7 \pm 2.1 [<\mathcal{M}(\text{KNN})]$$

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$$\sigma_{\bar{K}NN}^{tot} \times Br (\mu b) = \\ 29 \pm 3 \pm 3 [\text{all}] \\ 7.2 \pm 0.6 \pm 0.7 [<\mathcal{M}(\text{KNN})]$$

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" K^-pp " → $\pi^-\Sigma^+p$

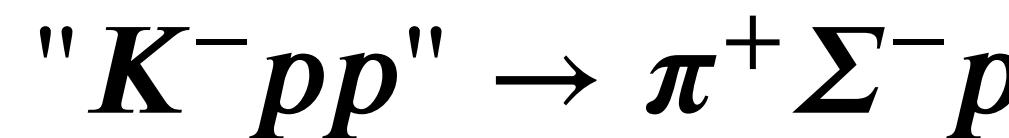
$$\sigma_{\bar{K}NN}^{tot} \times Br (\mu b) = \\ 110 \pm 8 \pm 8 [\text{all}] \\ 9.4 \pm 0.4 \pm 0.7 [<\mathcal{M}(\text{KNN})]$$

Mesonic Decay (suppressed by the 4-body phase space)

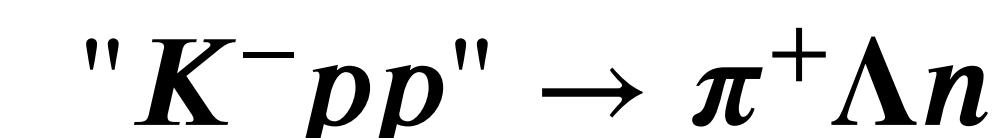
- $\Gamma_{YN} < \Gamma_{\pi YN}$: mesonic decay is dominant (about ~ 10 times in total)



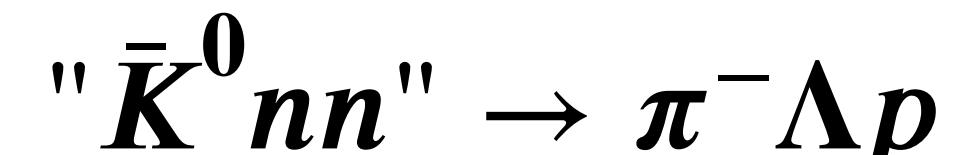
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 $9.3 \pm 0.8^{+1.4}_{-1.0}$ [all]
 $5.5 \pm 0.5^{+0.8}_{-0.6}$ [$< M(KNN)$]



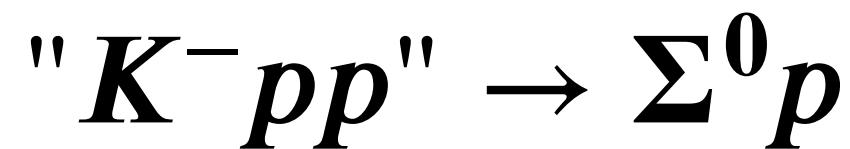
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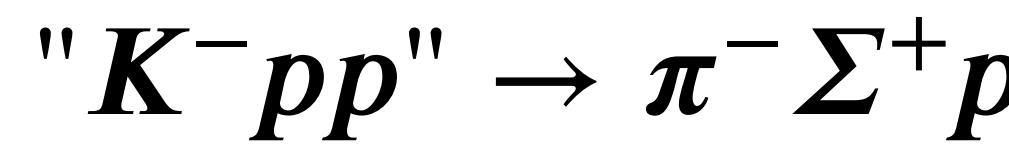
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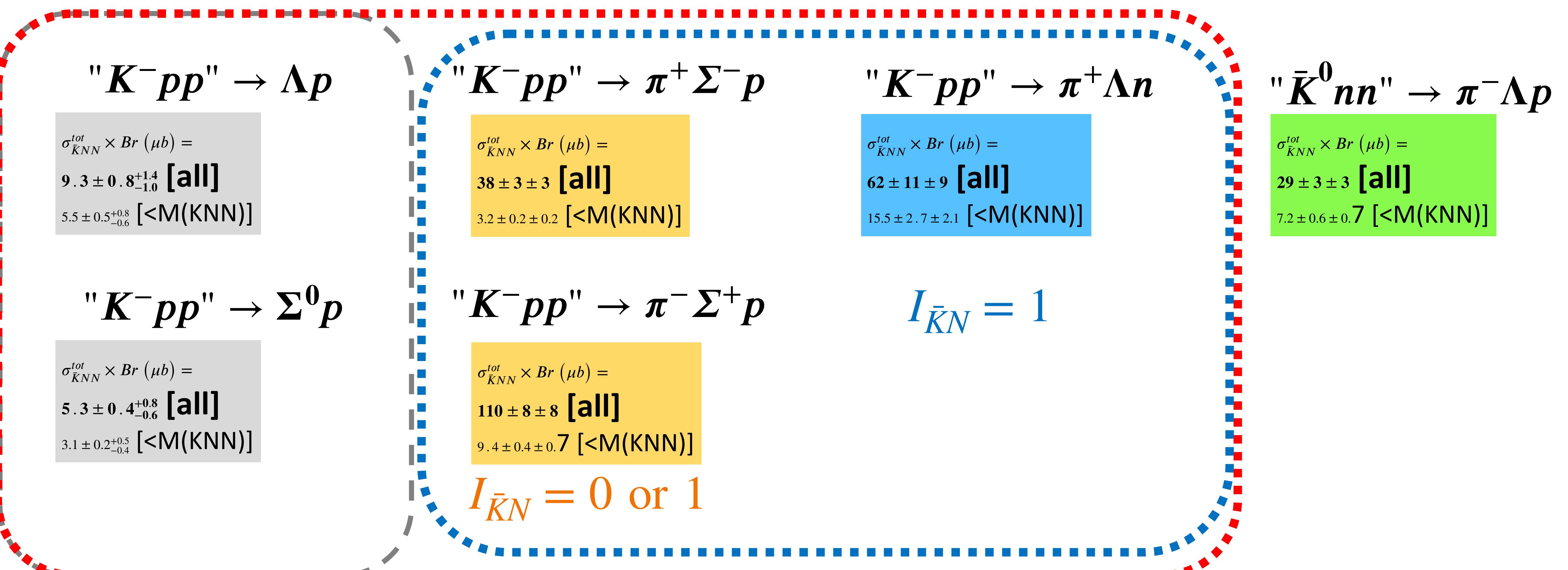
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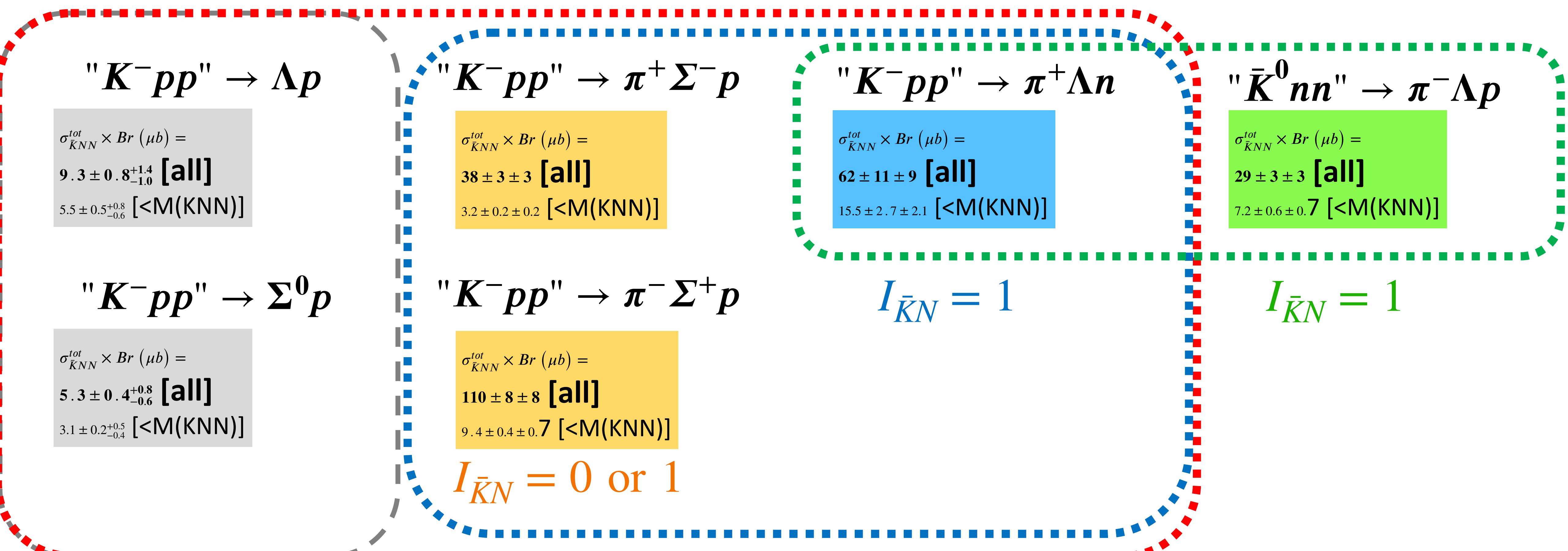
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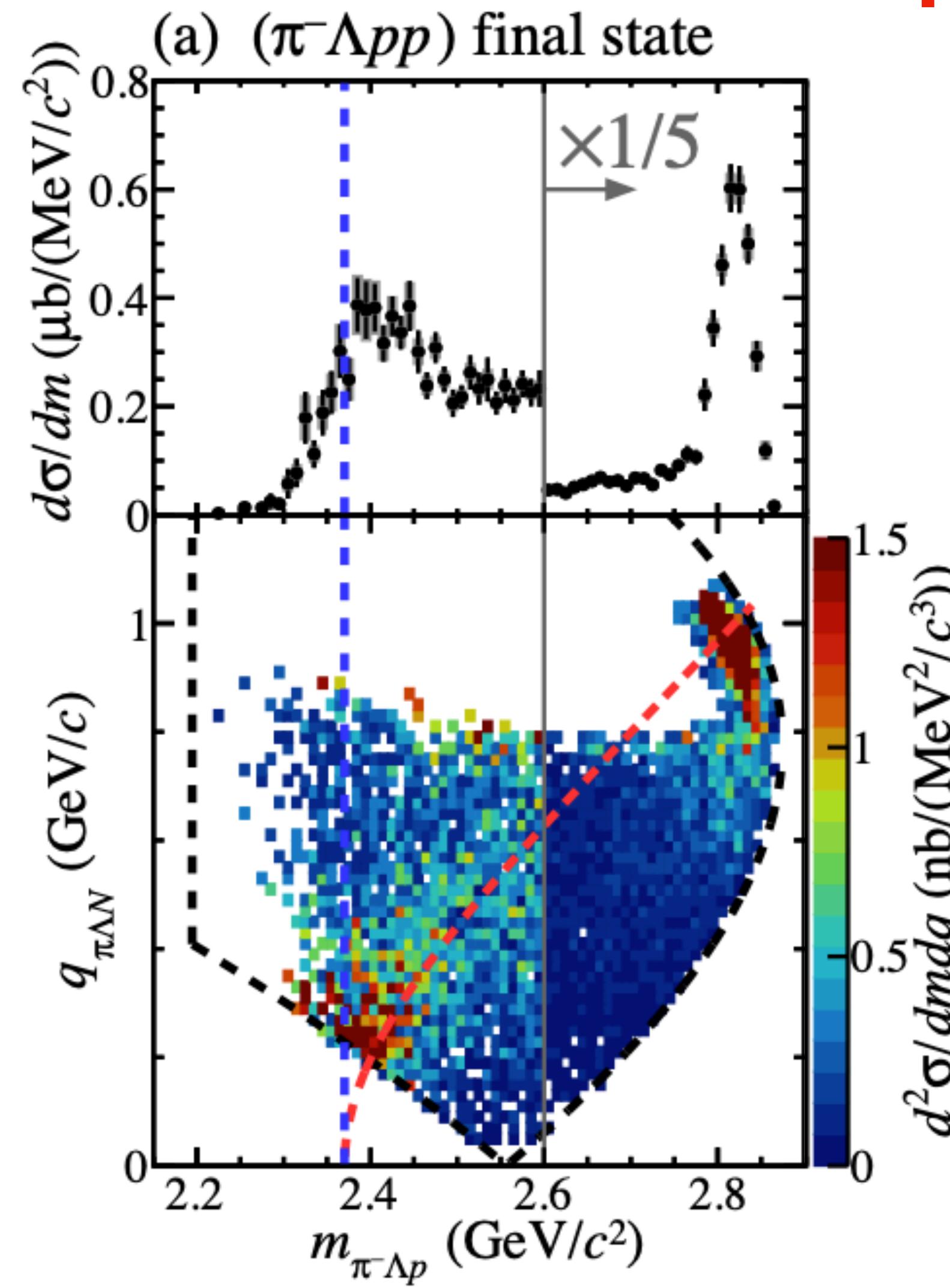
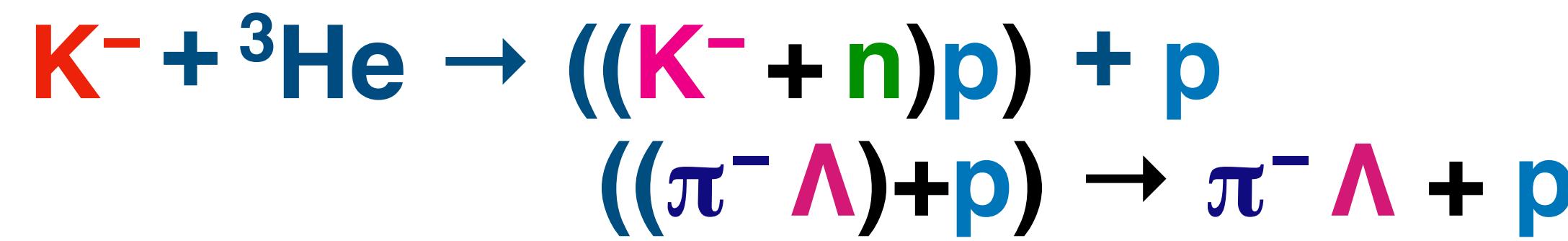


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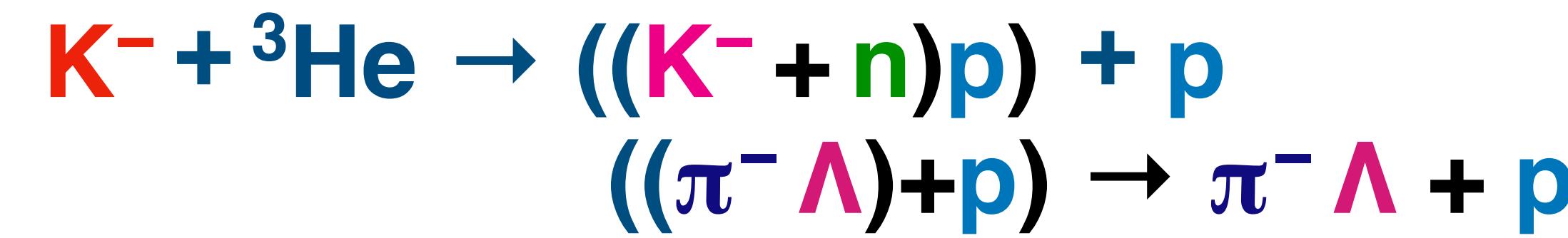
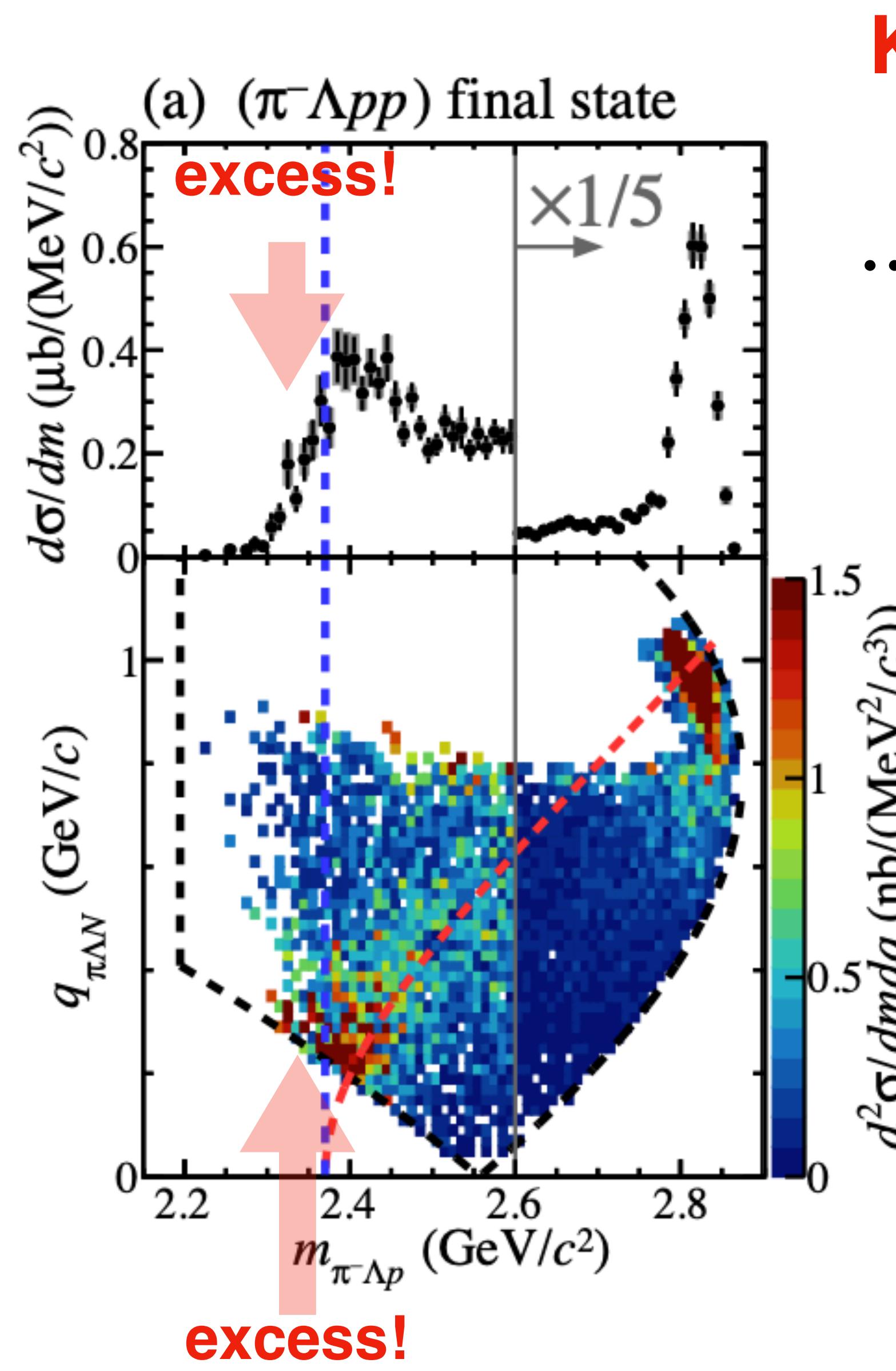
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- $\Gamma_{\pi^+\Lambda n}/\Gamma_{\pi^-\Lambda p} \sim 2$: if we assume $Br_{\{K^-pp\} \rightarrow \pi^+\Lambda p}/Br_{\{\bar{K}^0 nn\} \rightarrow \pi^-\Lambda p} \approx 1$ & $\sigma_{\{K^-pp\}}/\sigma_{\{\bar{K}^0 nn\}} \approx 2$



Mesonic decay branch of $\bar{K}^0 nn$?

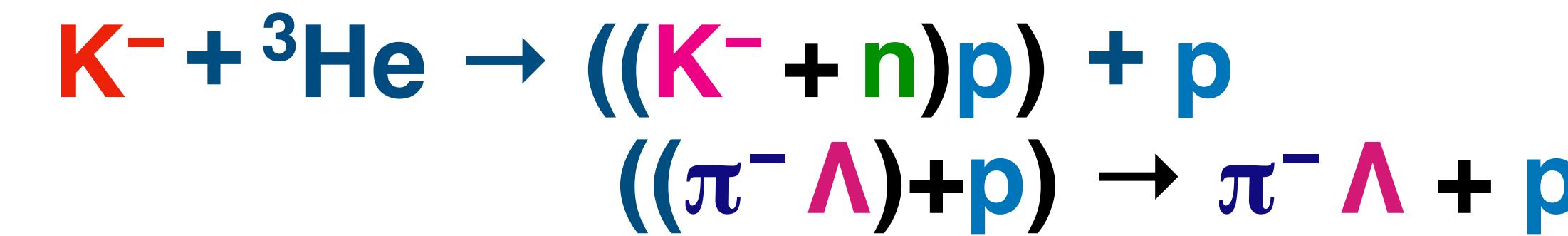
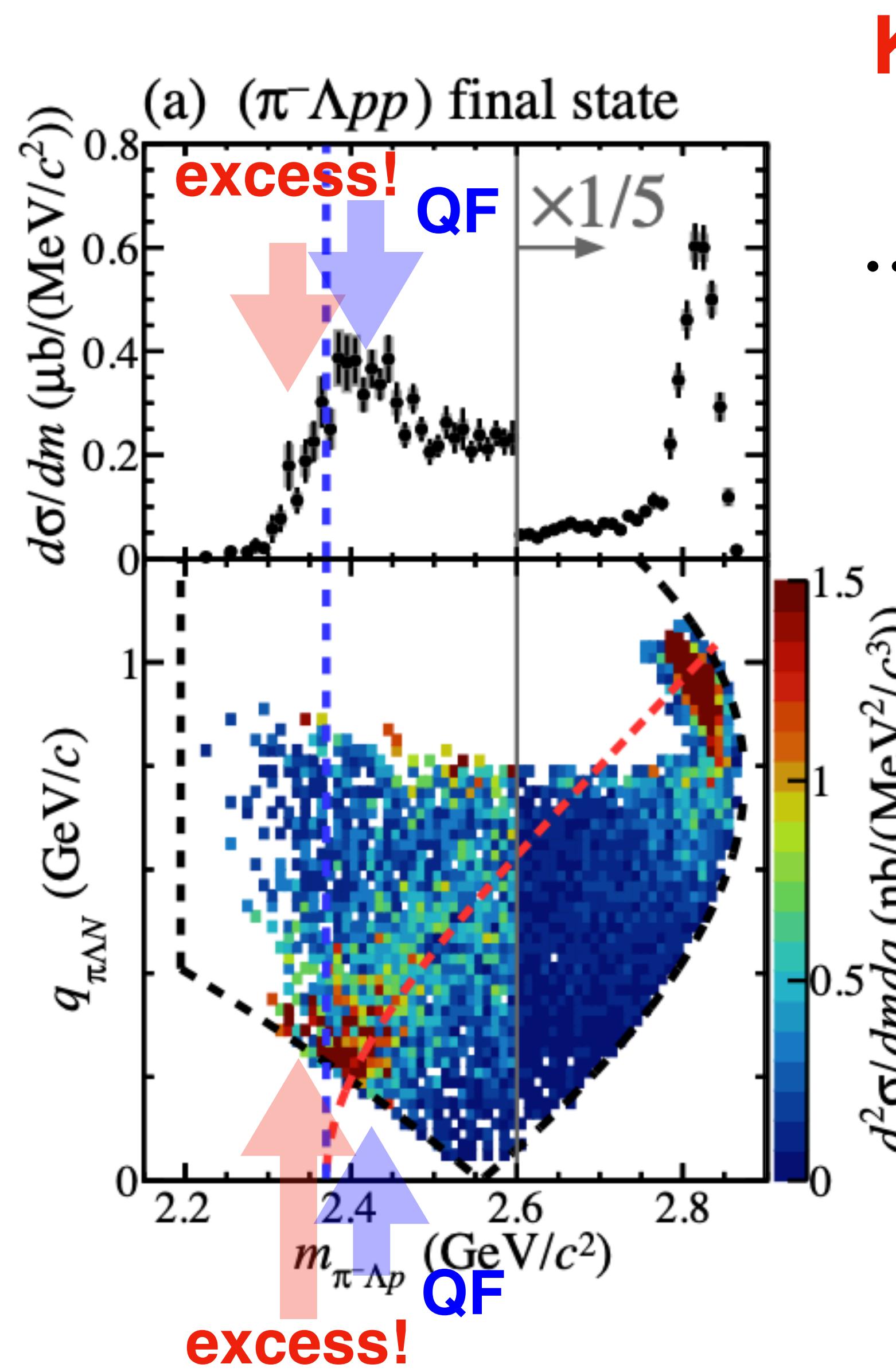


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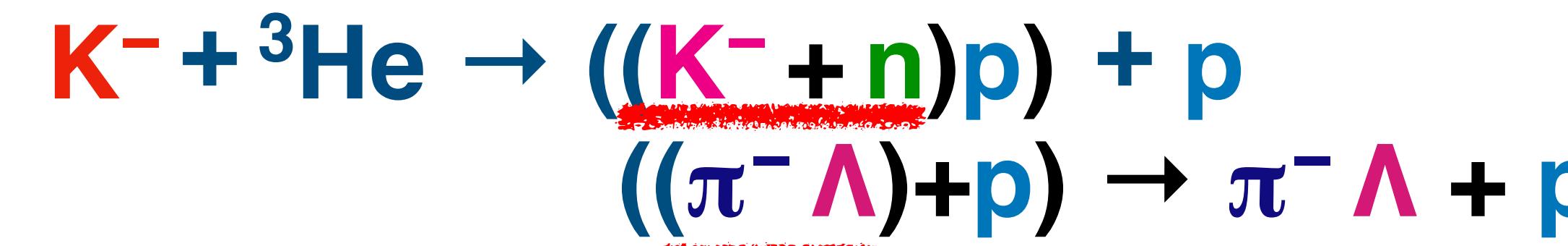
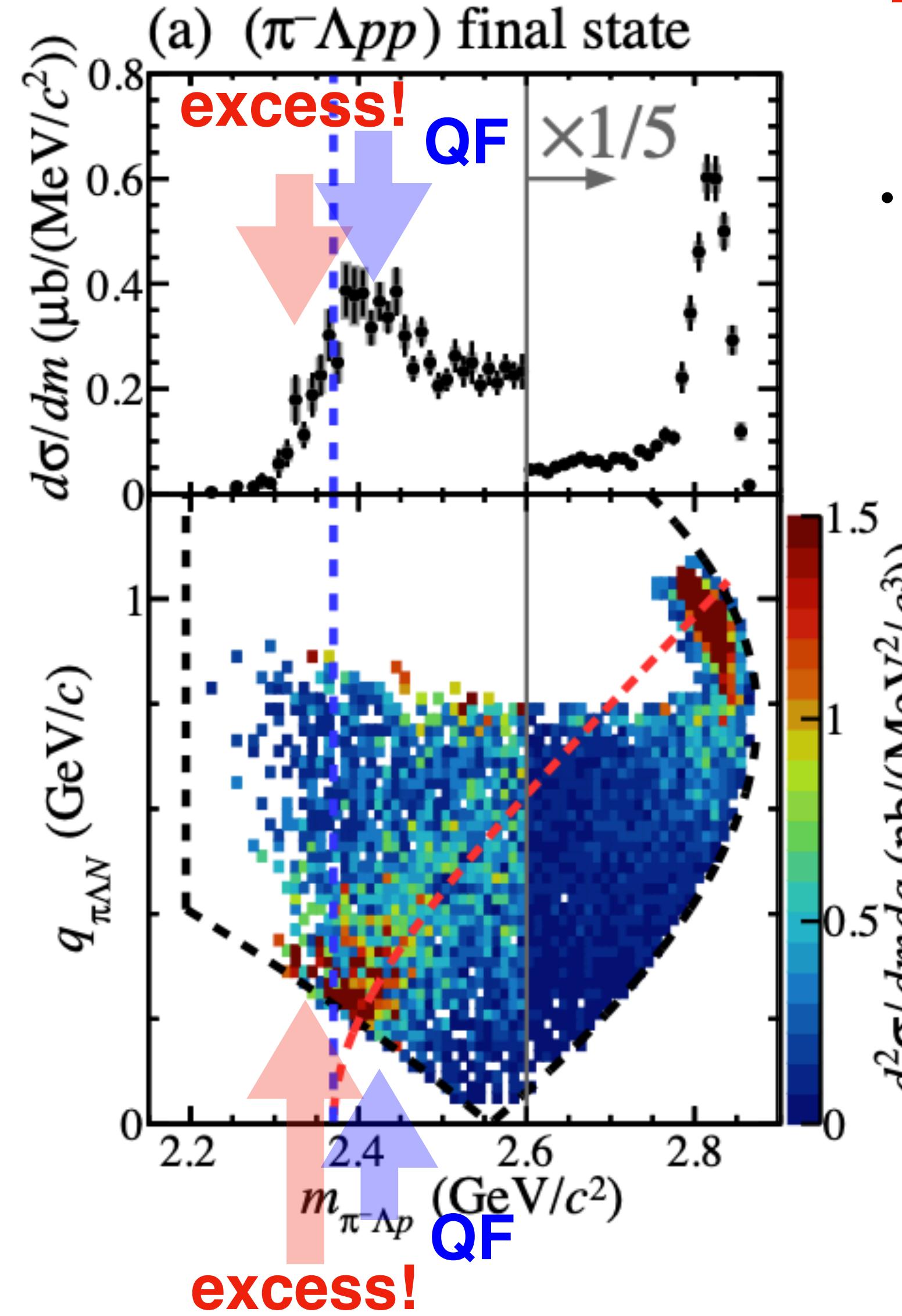
*... but the **excess** is still not easy to see ...*

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... but the excess is still not easy to see ... due to the QF-K

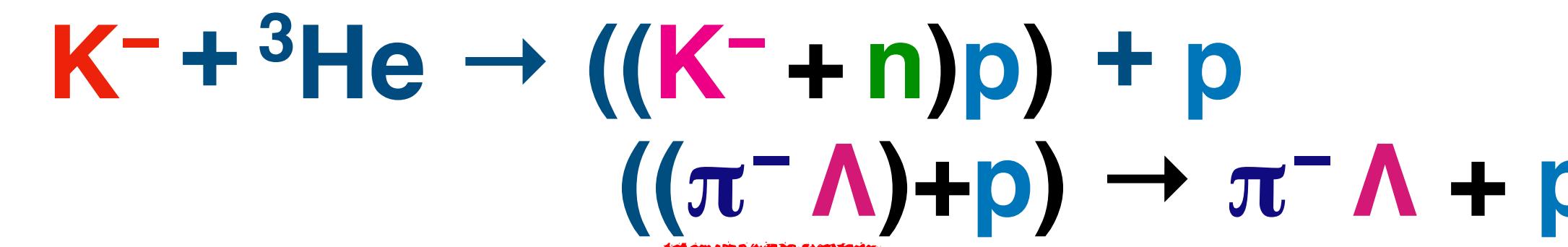
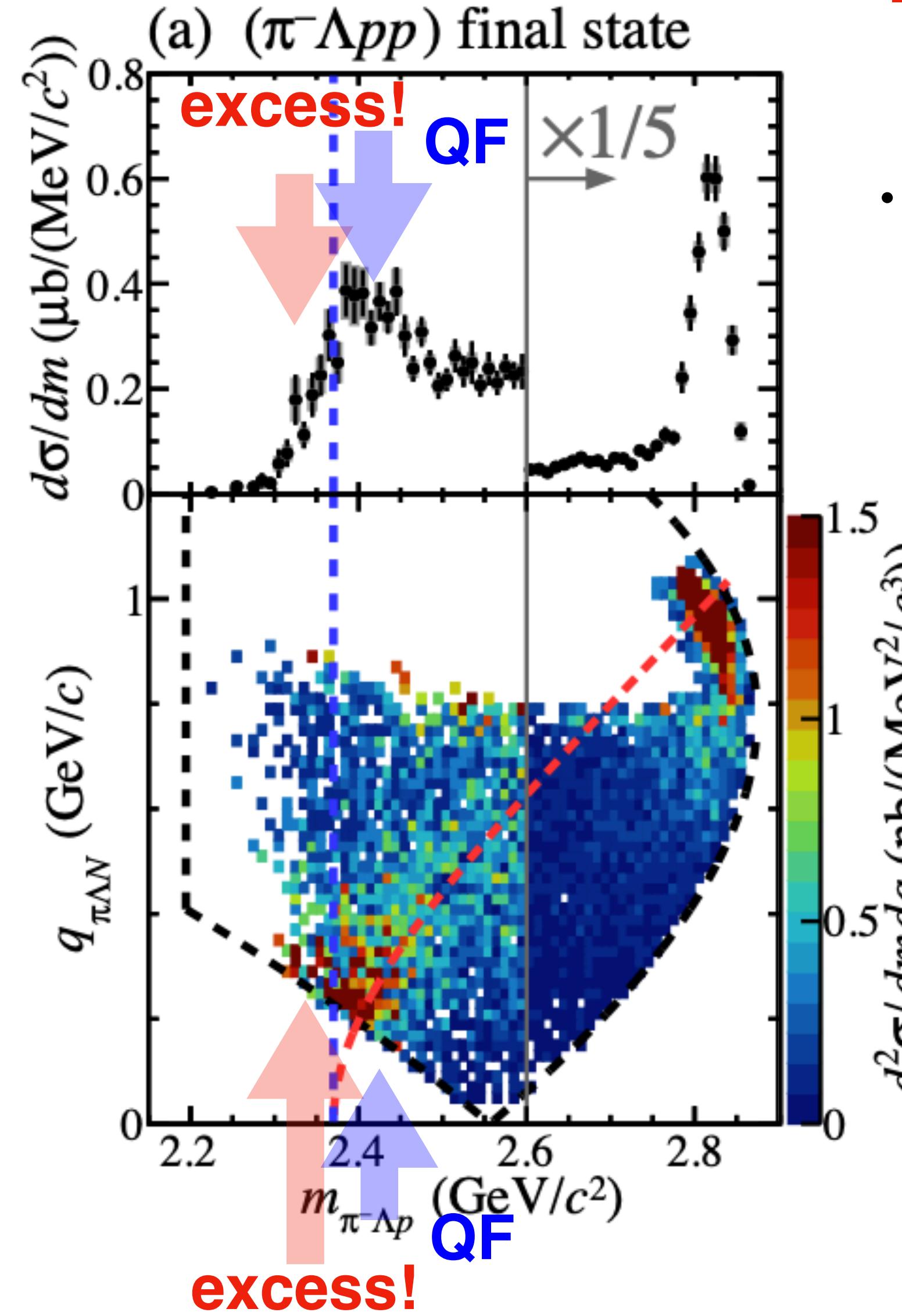
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... but the *excess* is still not easy to see ... due to the *QF-K*

If $m_{\pi^- \Lambda} \geq m_{K^-} + m_n$, then the " $K^- np$ " bound state cannot be formed, but will form a quasi-free background.

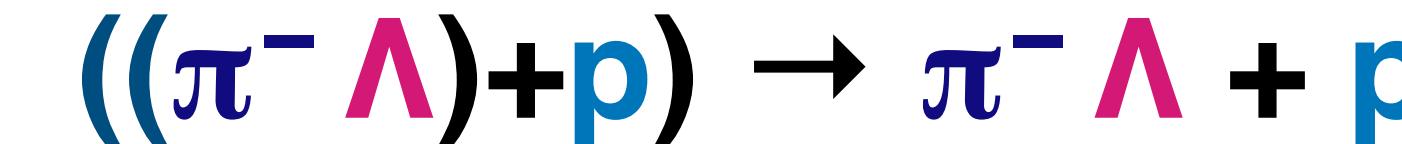
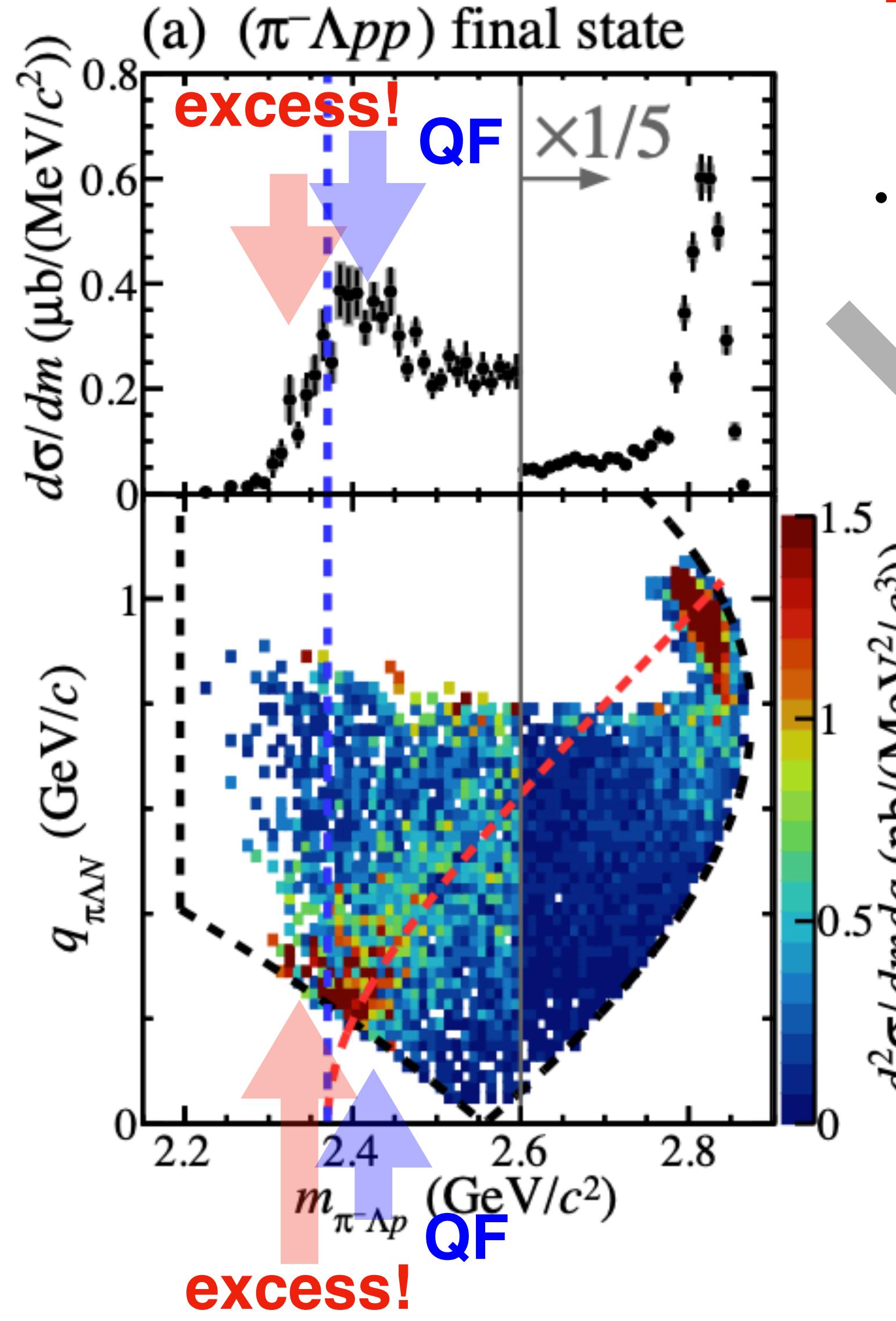
Mesonic decay branch of $\bar{K}^0 nn$?



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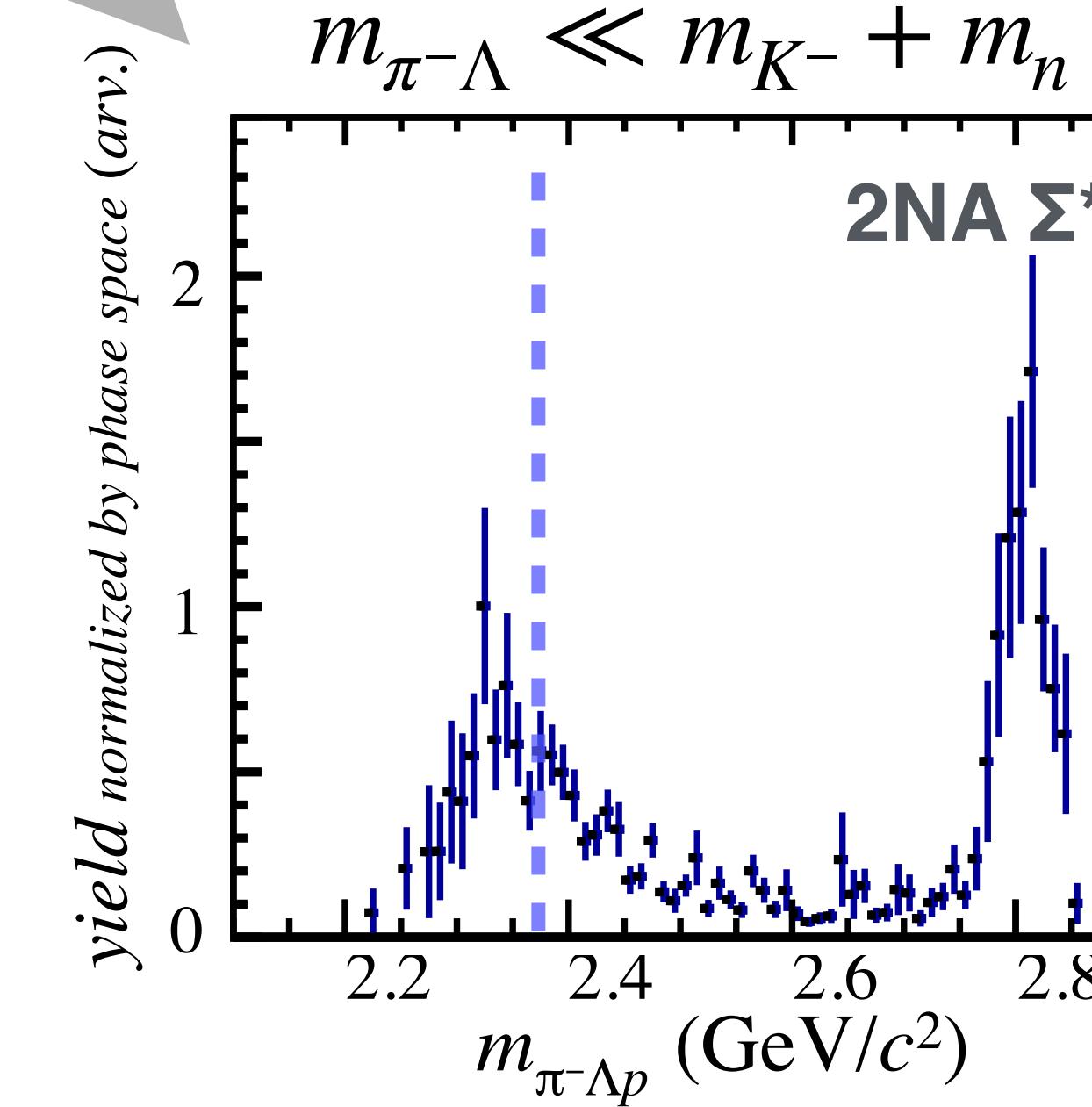
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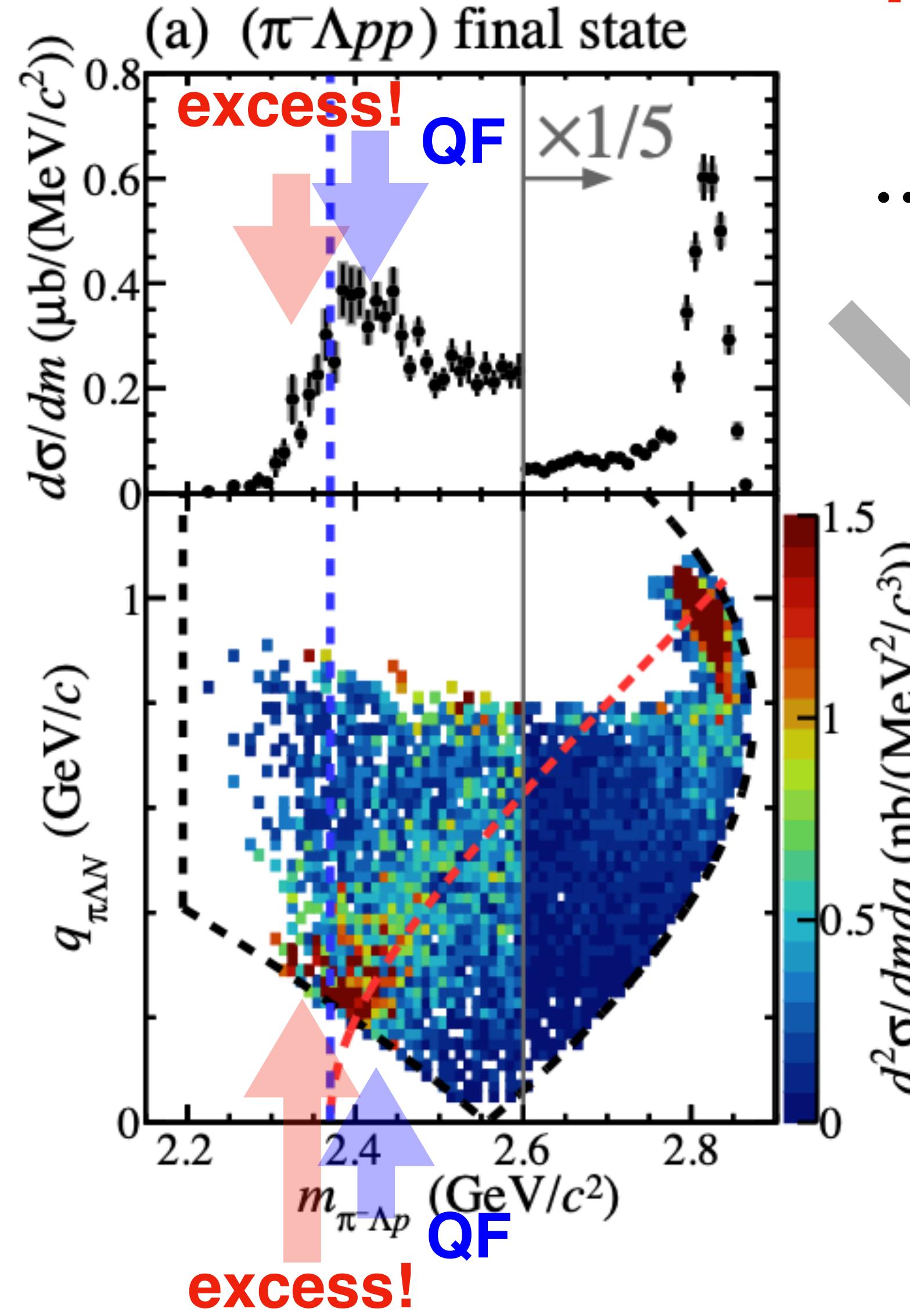


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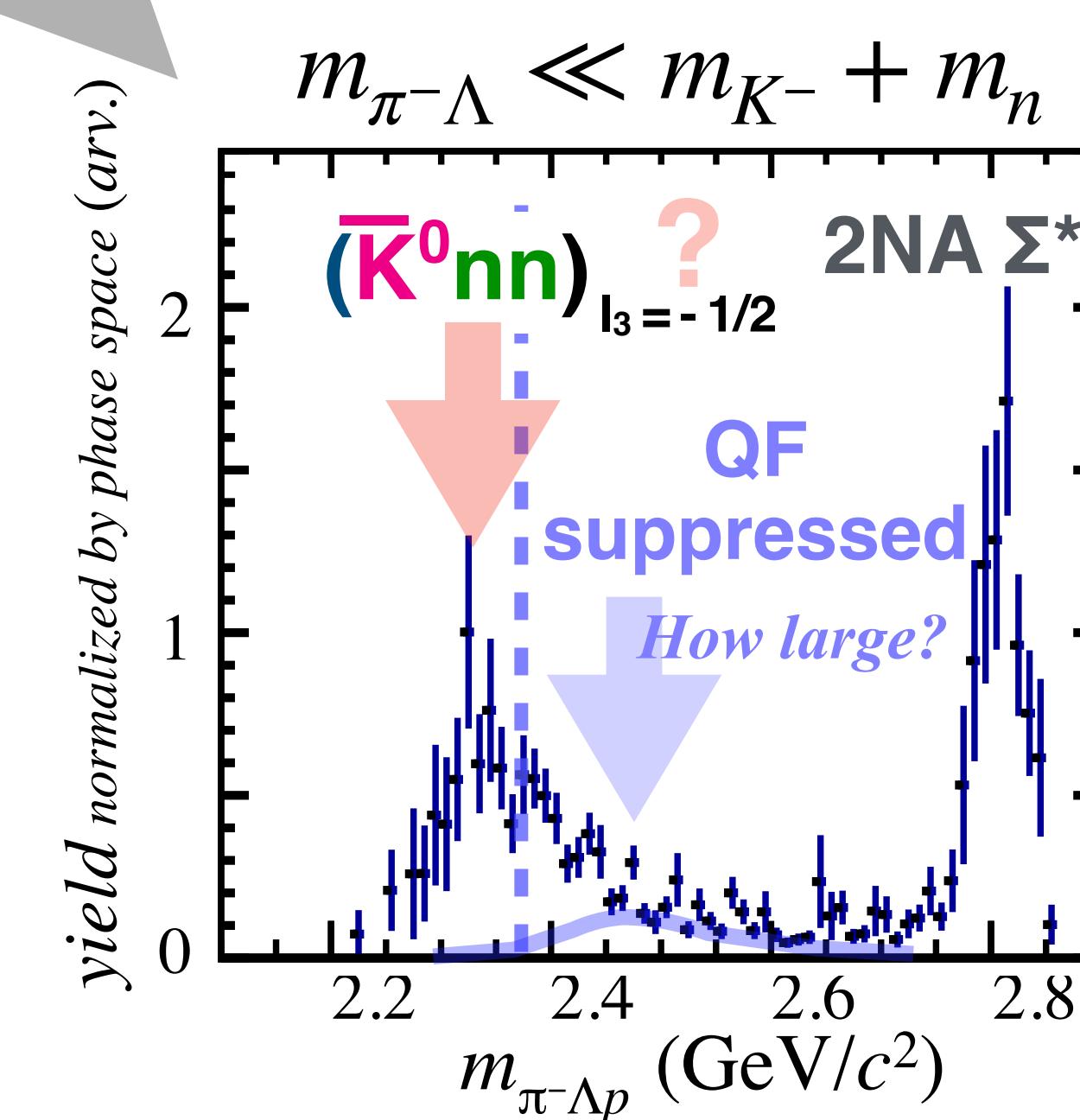


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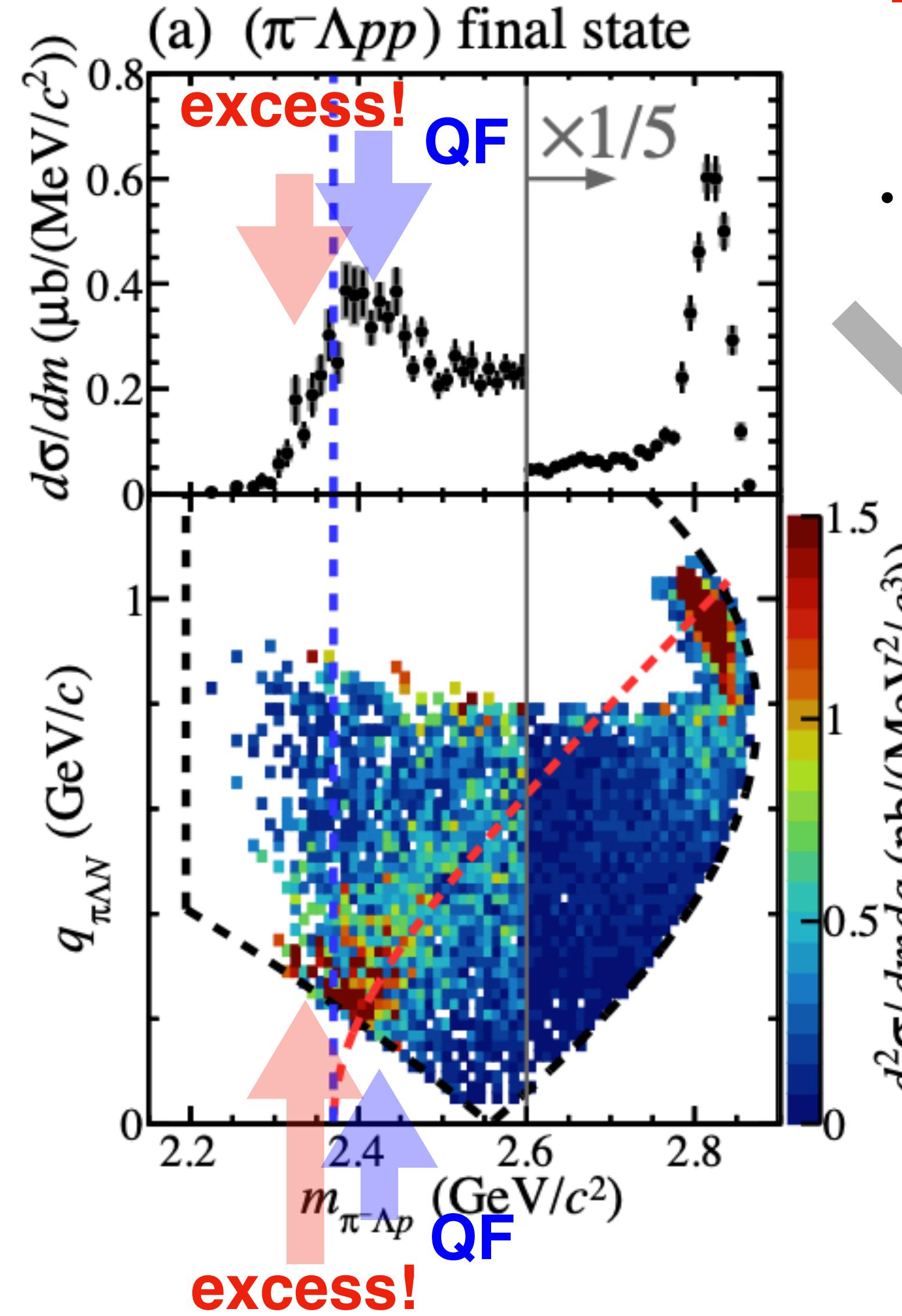
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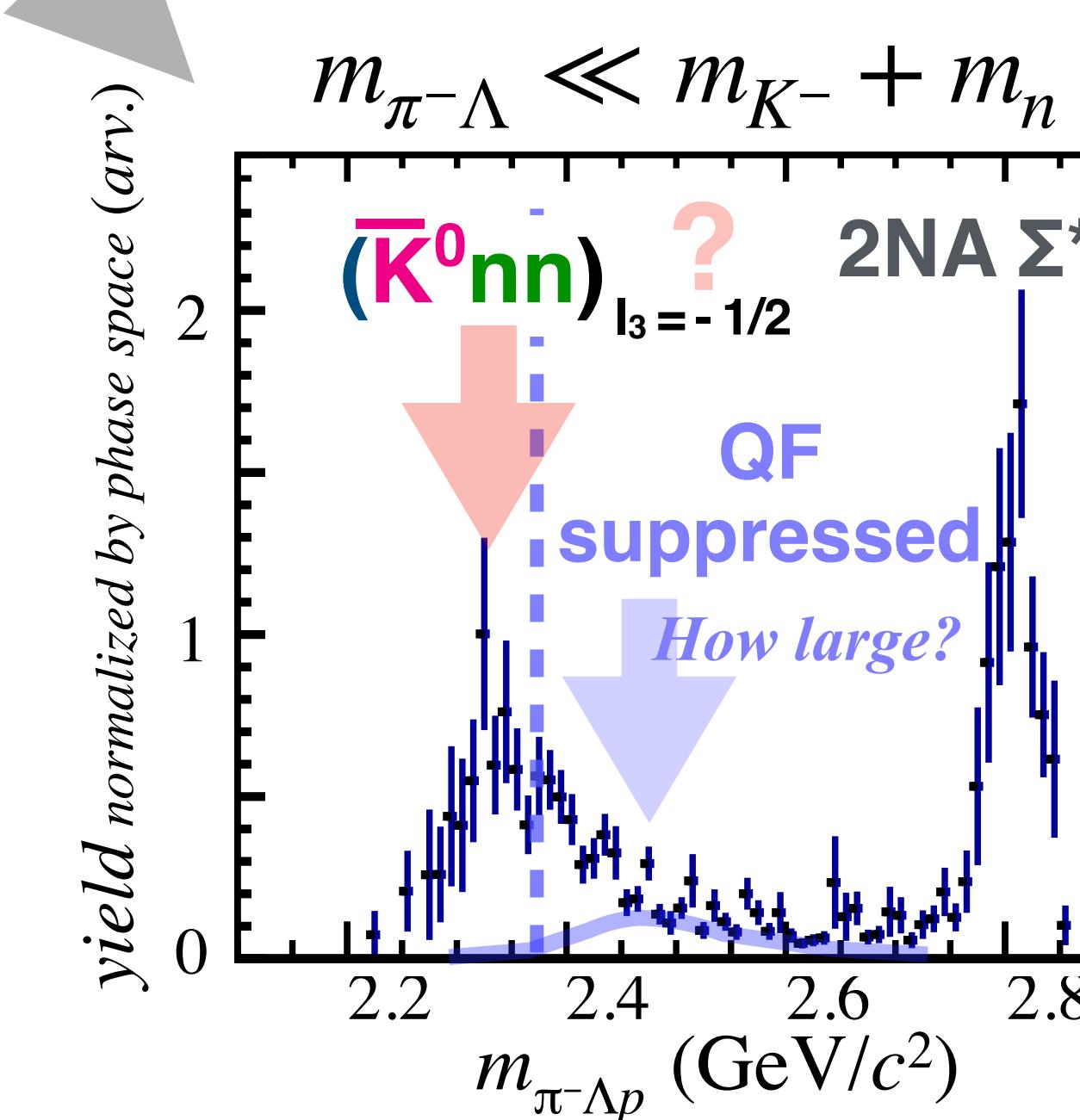
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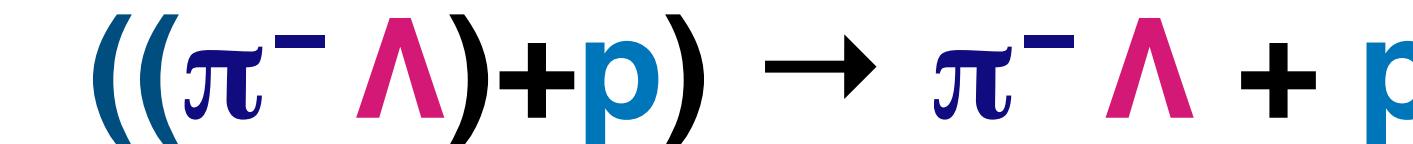
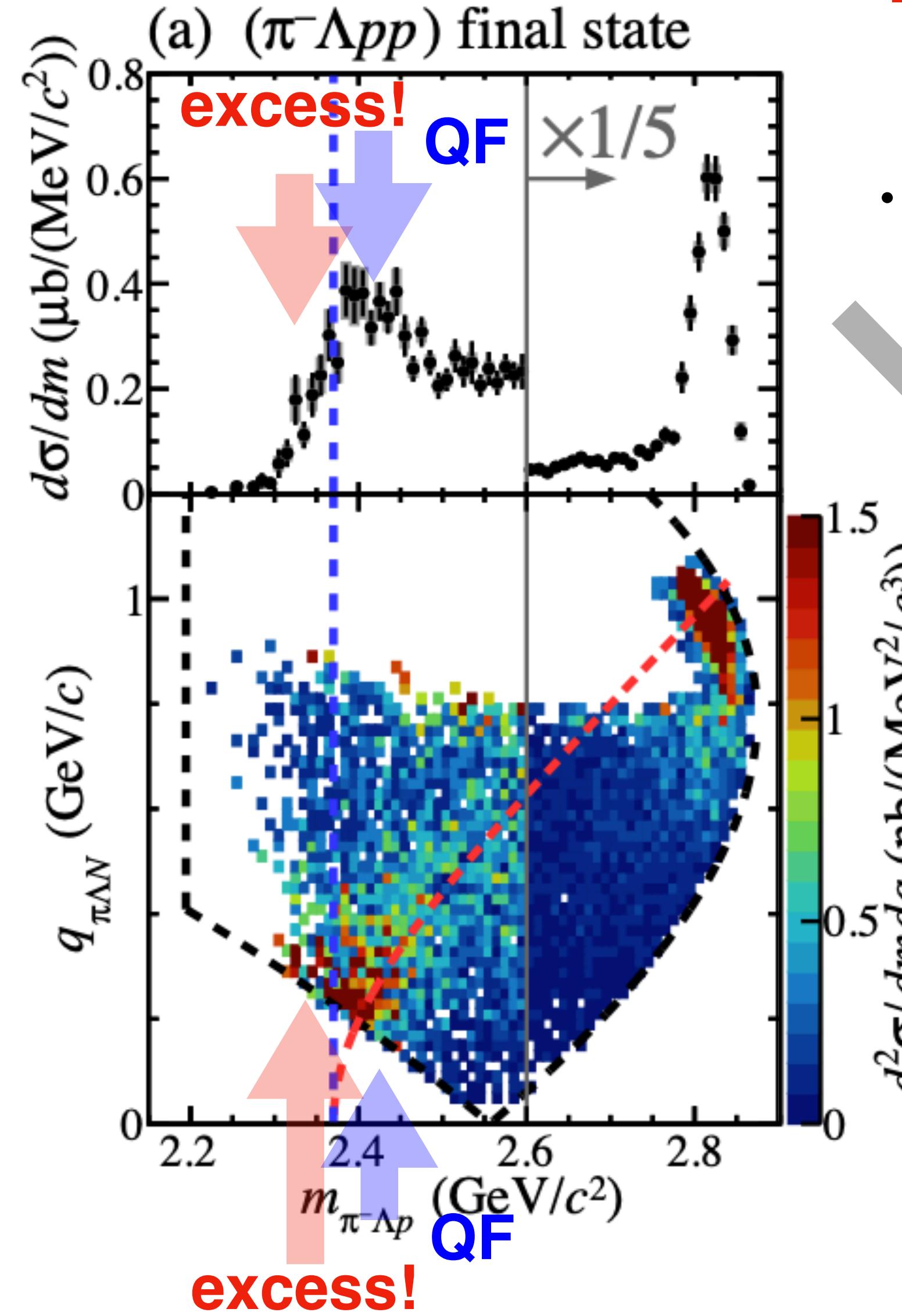
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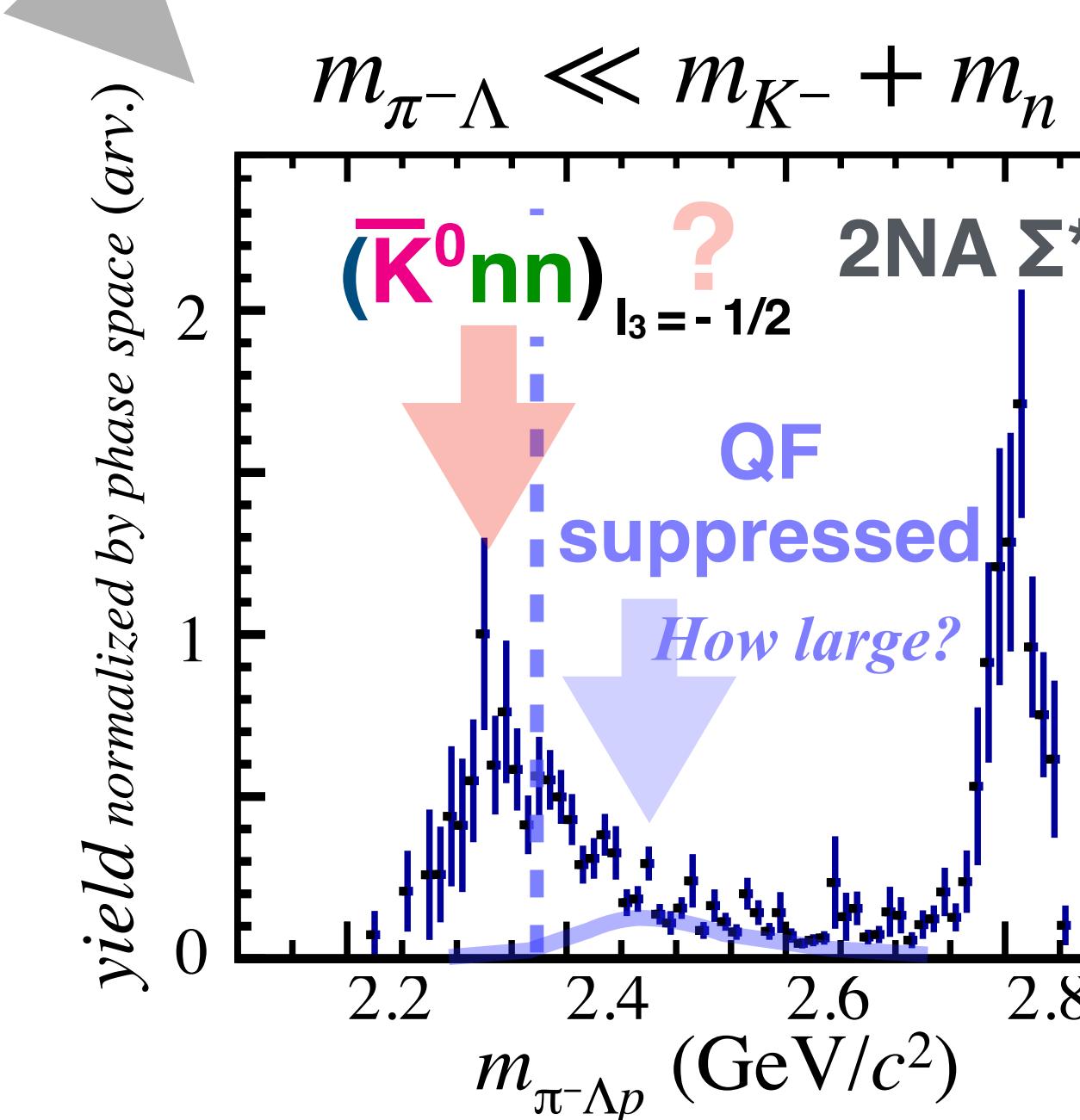


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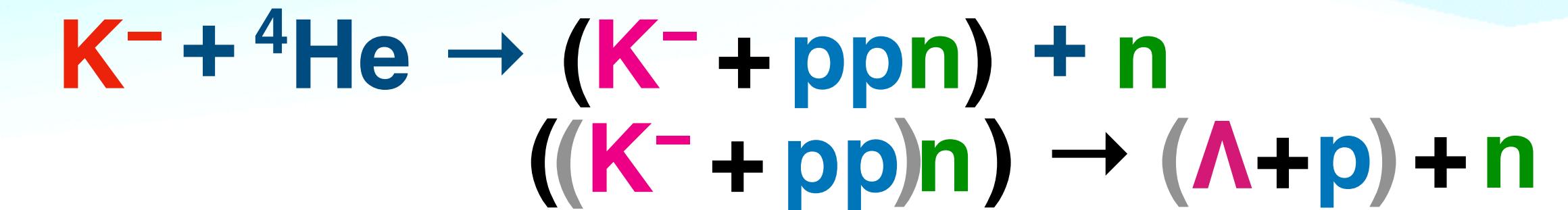
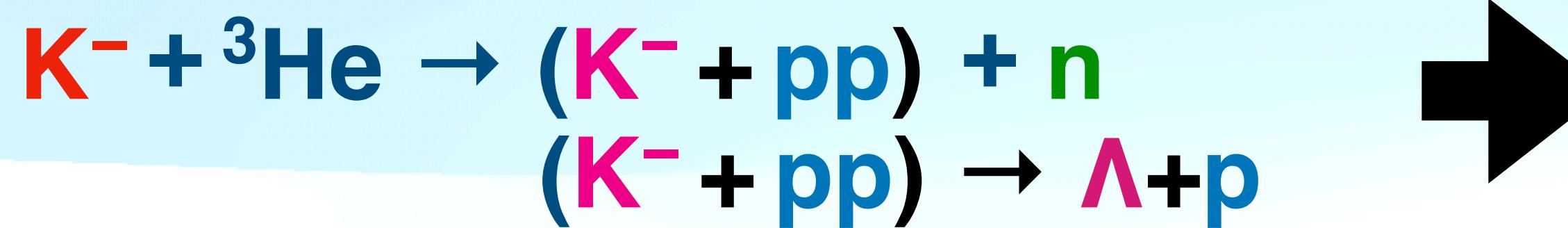


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Further analysis on other data

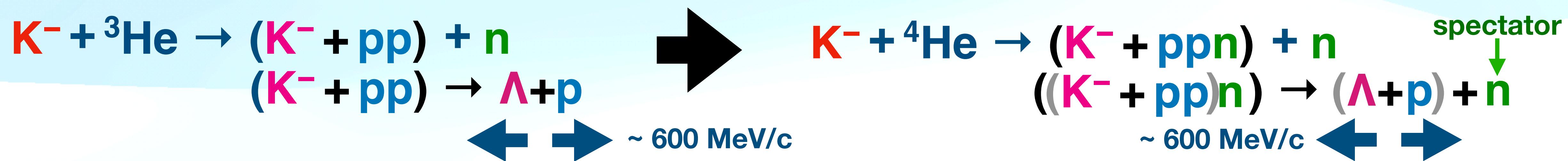
Signal of $\bar{K}NNN$?



Preliminary data analysis for $\bar{K}NNN$ formation study utilizing ${}^4\text{He}$ lifetime measurement via $K^- + {}^4\text{He} \rightarrow \pi^0 + {}^4\text{He}$ reaction giving us a very interesting result

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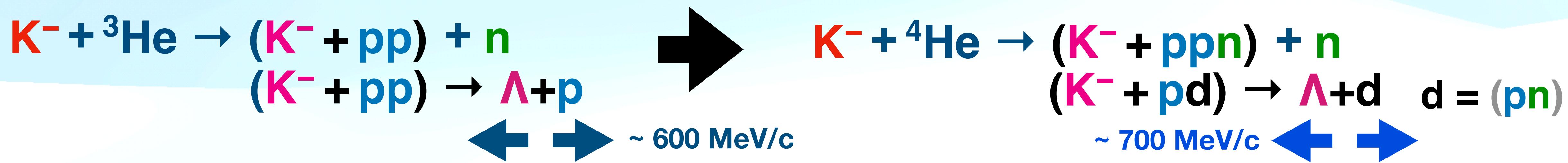
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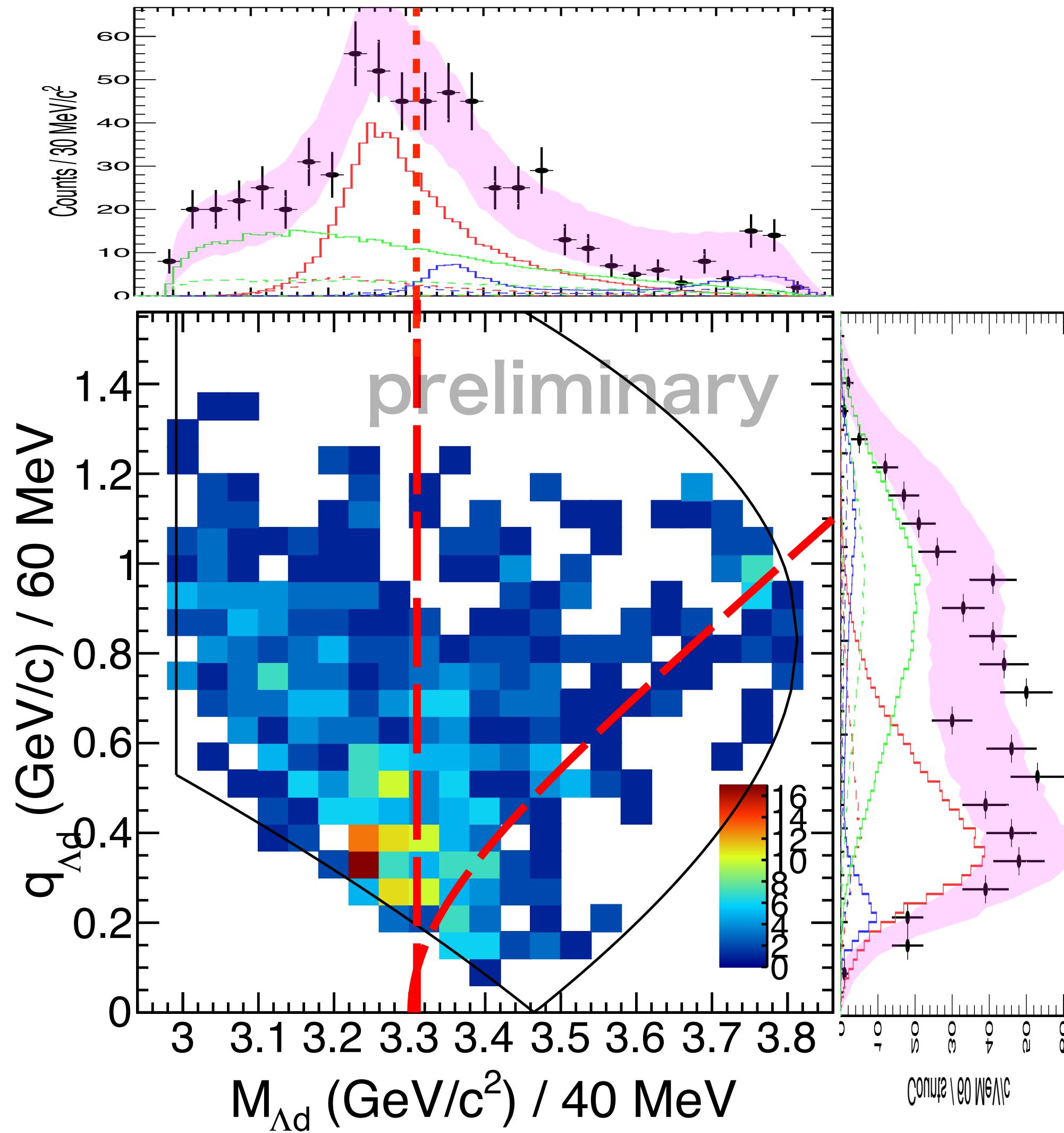
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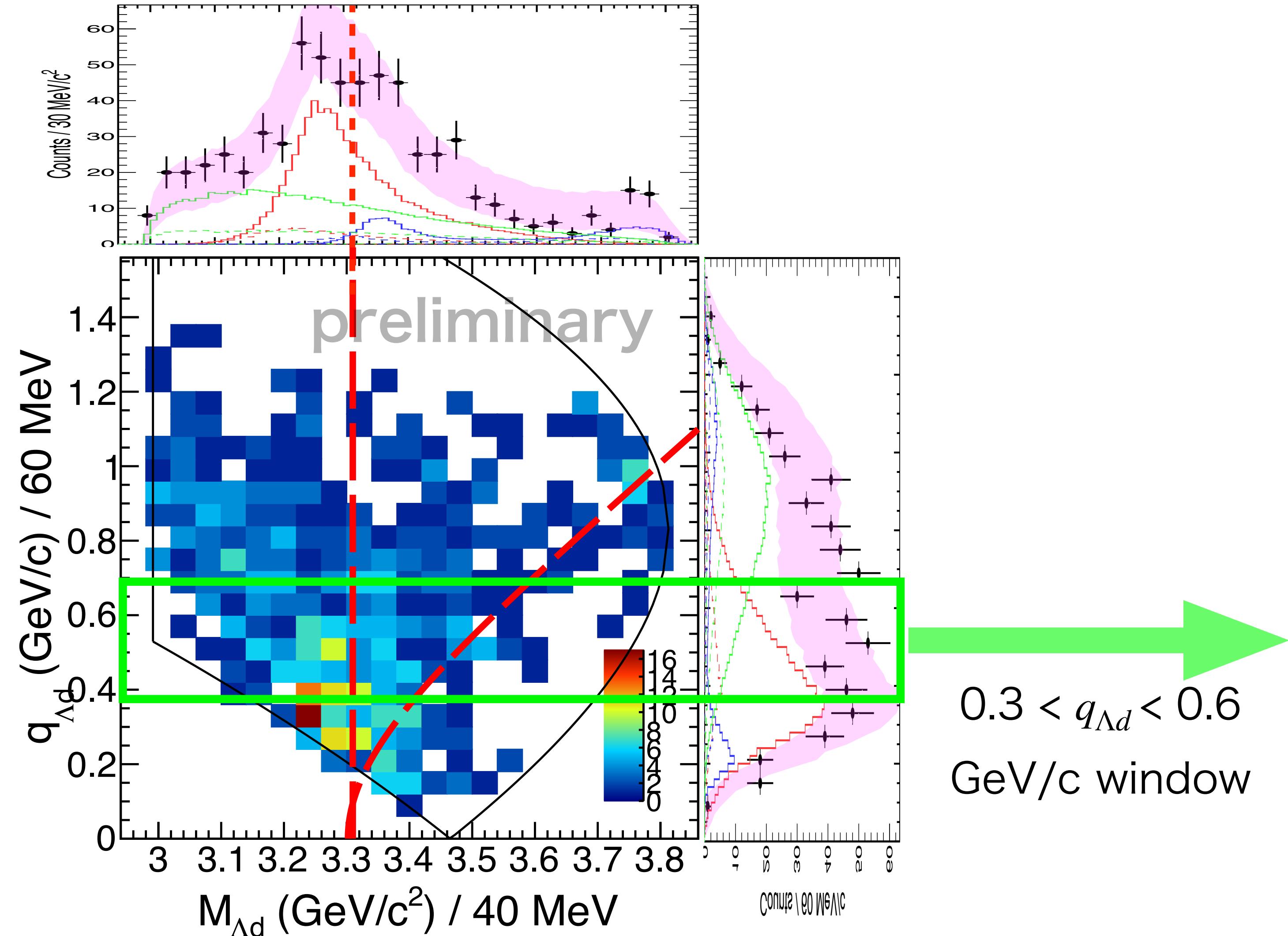
$K^- \bar{4}\text{He} \rightarrow \{\Lambda d\} + n$ Analysis (with the T77 Data)

Promising signal is observed!

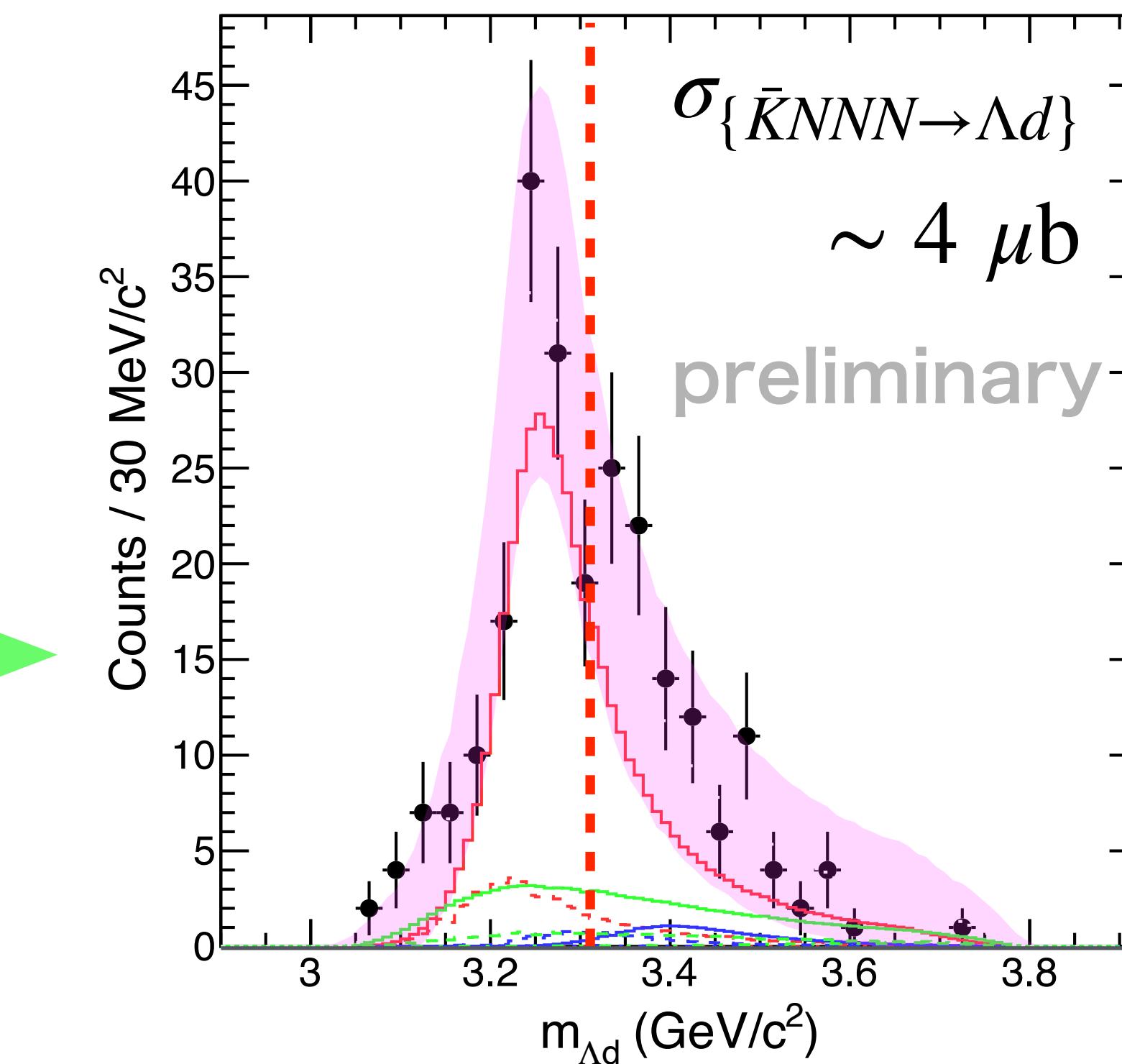


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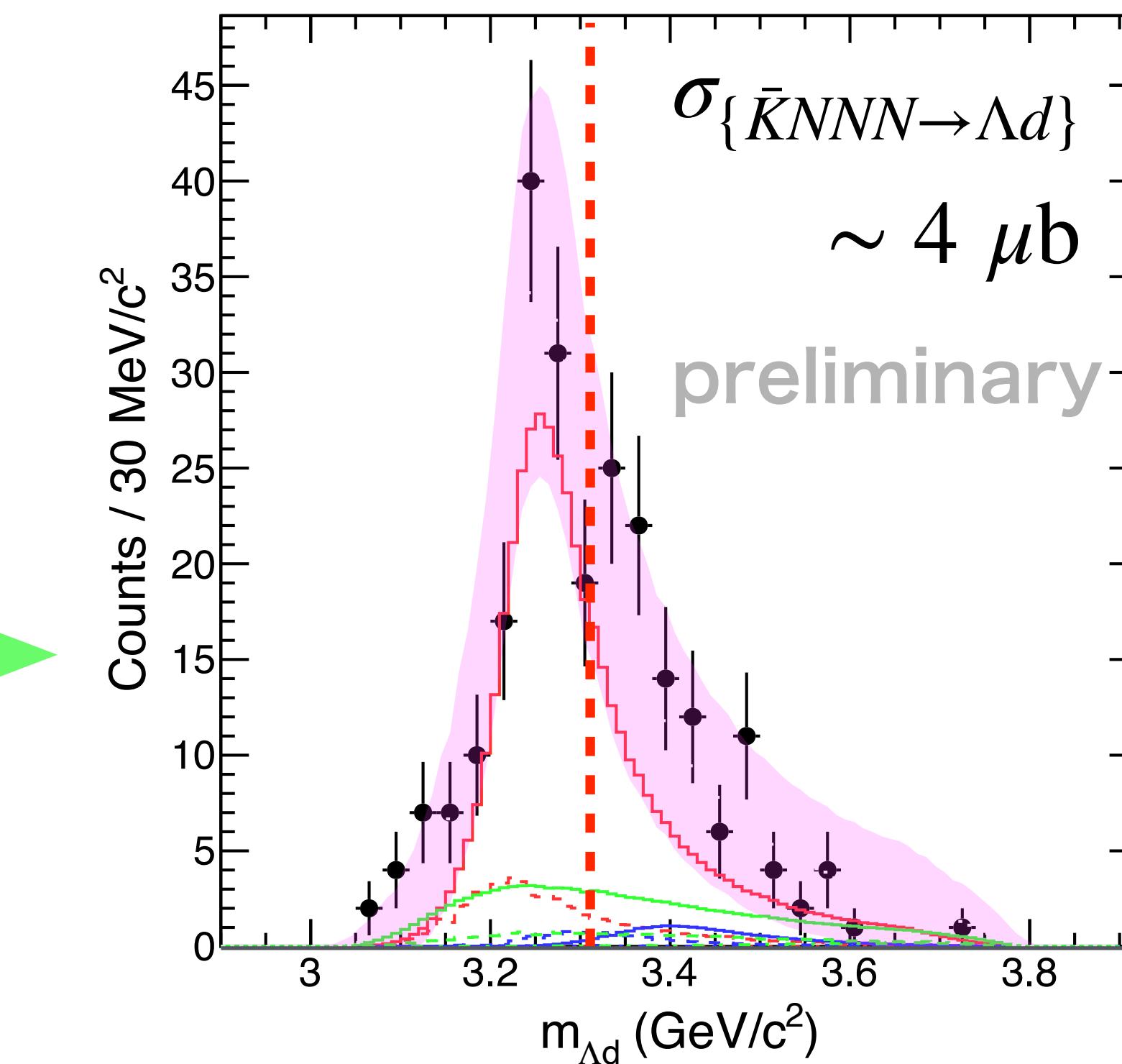
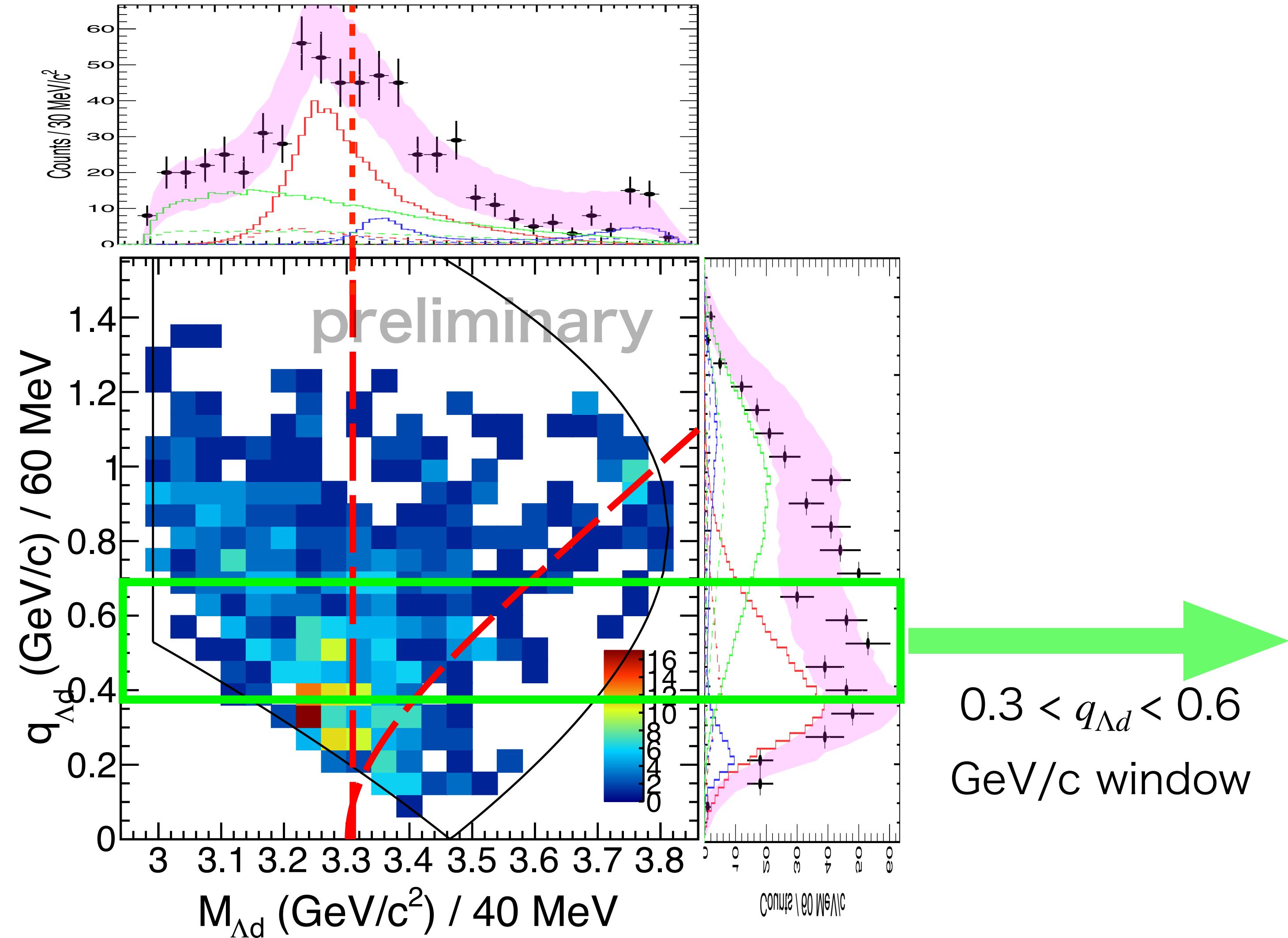


$0.3 < q_{\Lambda d} < 0.6$
GeV/c window



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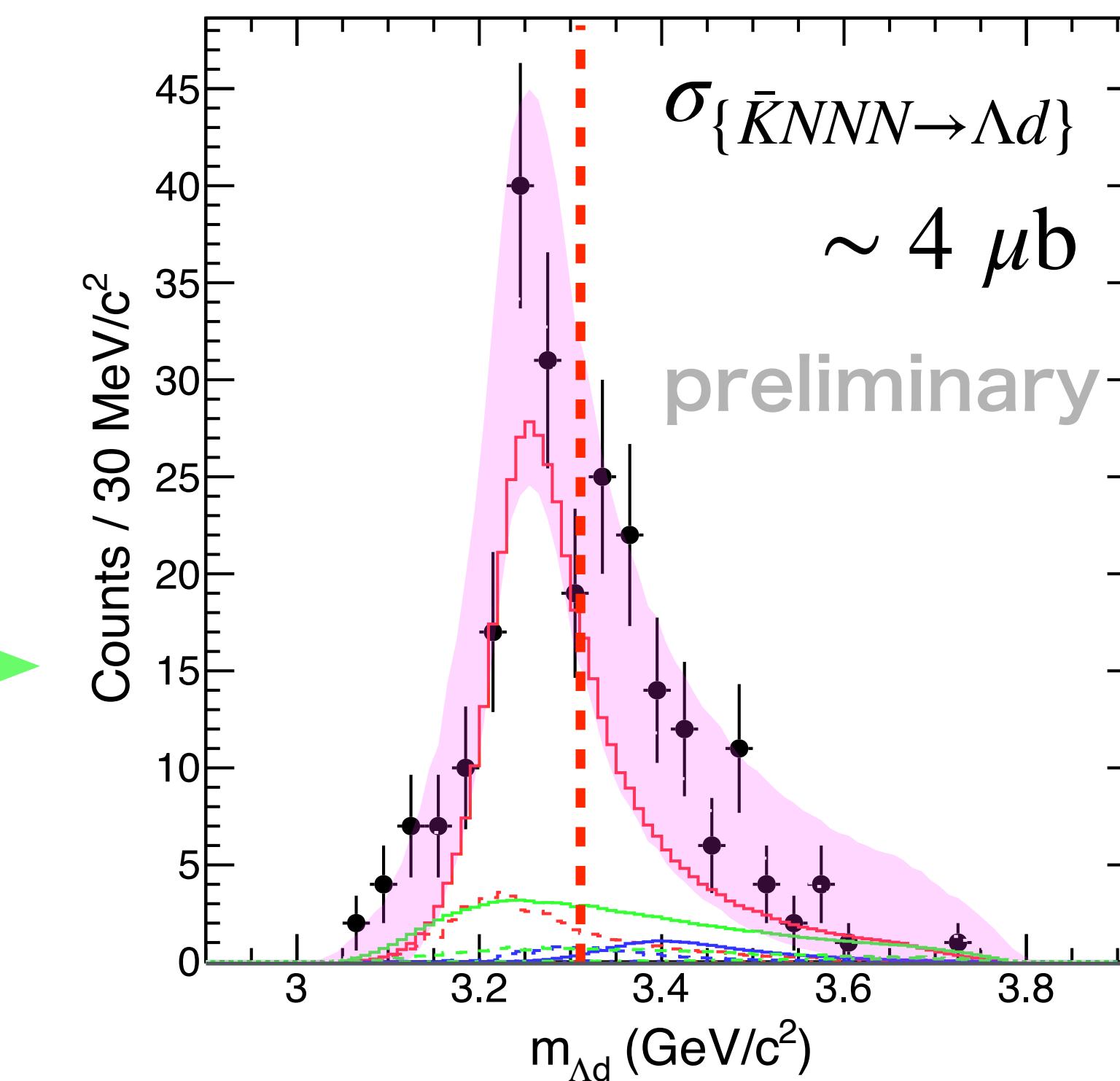
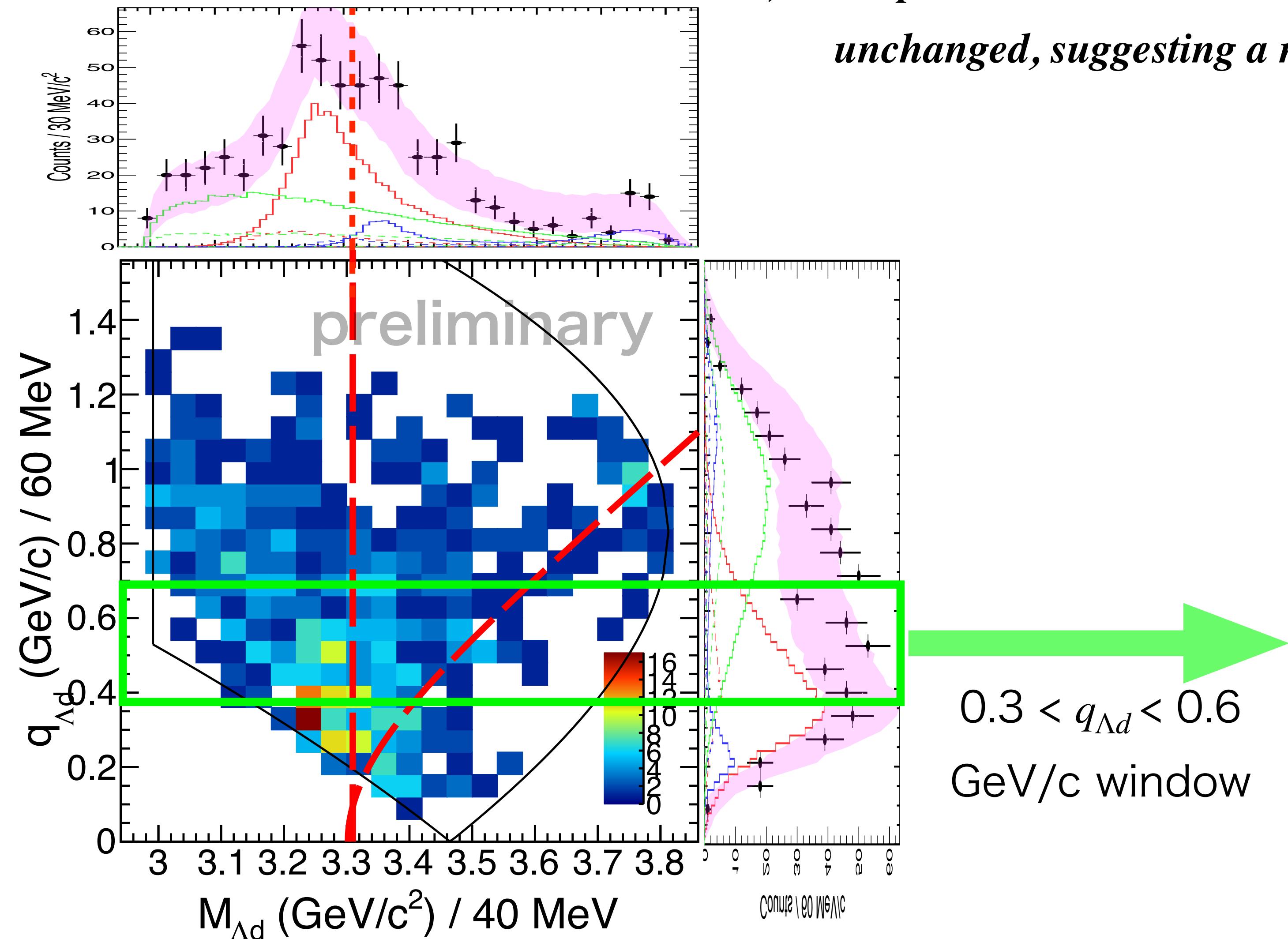
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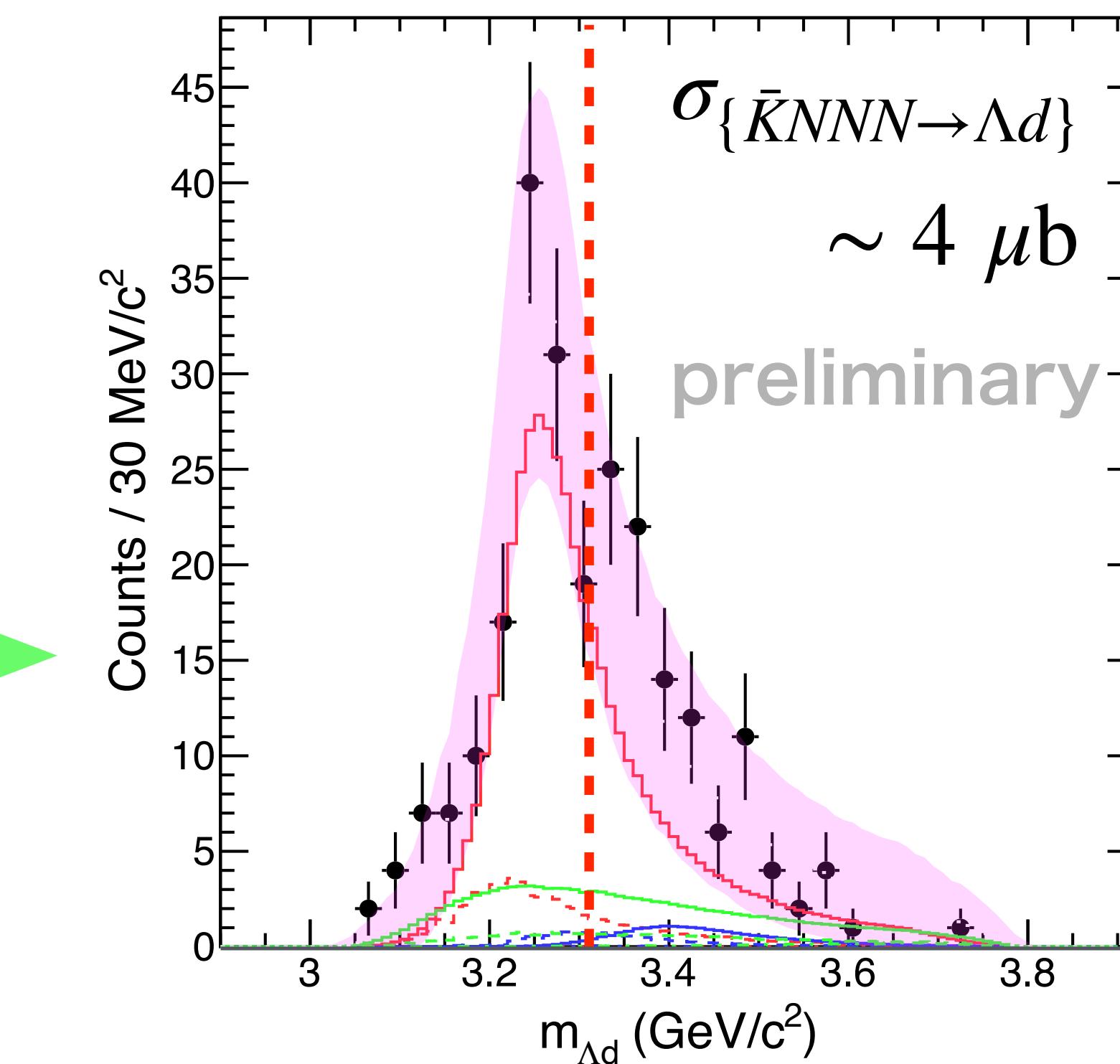
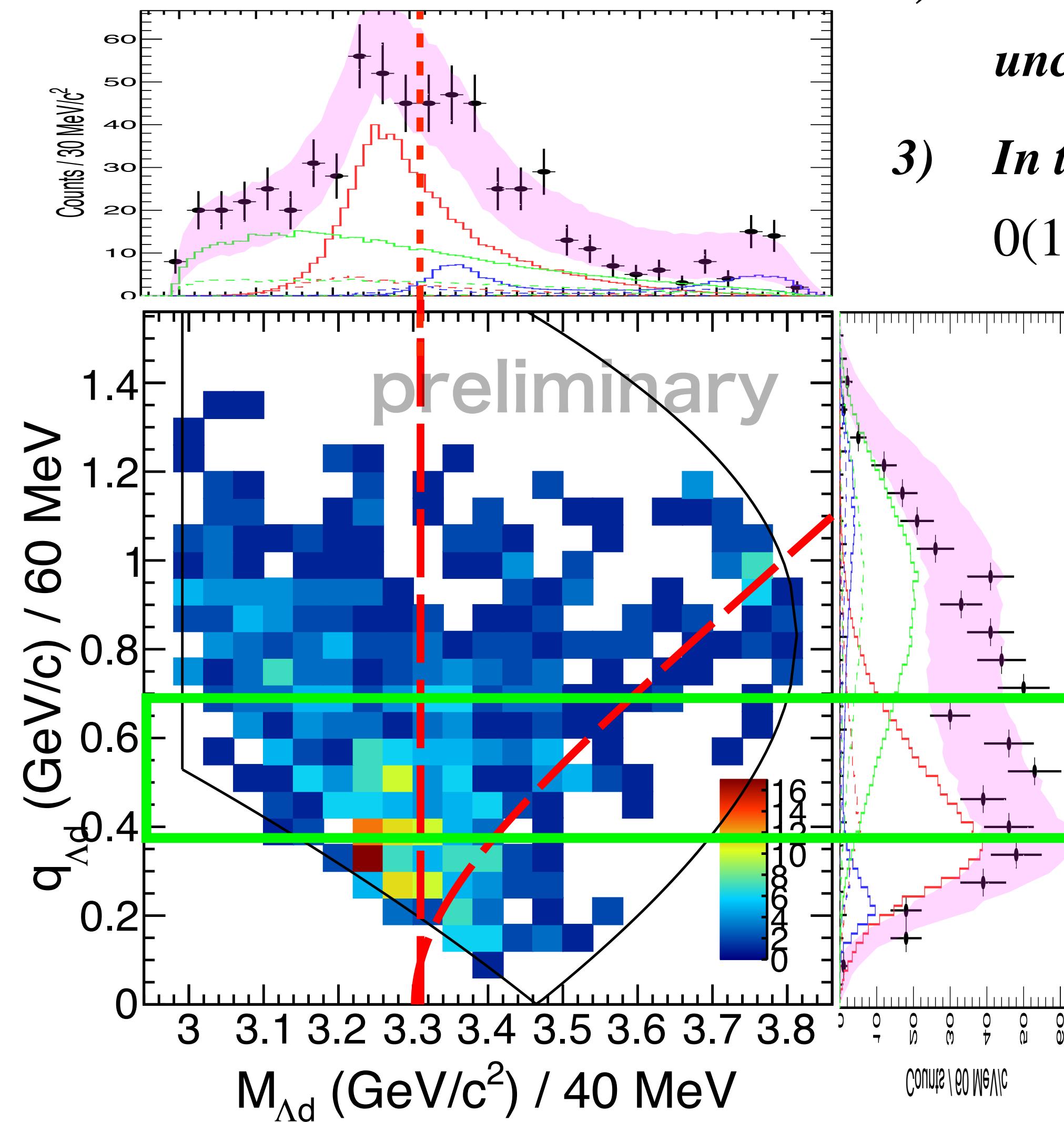
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- 1) *It suggests kaonic nuclei exist more universally, not just in cases like $K^- pp$.*
- 2) *Despite more nucleons absorbing the K meson, the width remains unchanged, suggesting a near plateau.*
- 3) *In the case of $K^- ppn$, isospin and spin-parity $I(J^P)$, has been fixed to be $0(1/2^-)$ with high certainty through Λp decay.*



$I(J^P)$ of X in $K^- {}^4He \rightarrow X + n$

What is the $I(J^P)$ of the observed state? := 0(1/2⁻)

- Λpn would be the major decay mode (relatively easily identified in this channel)

1. “X” $\rightarrow \Lambda d$ decay mode is unique evidence of $I(X) = 0$

- $I(J^P) : \Lambda = 0(1/2^+), d = 0(1^+), K^- = 1/2(0^-), {}^3He = 1/2(1/2^+), {}^4He = 0(0^+)$

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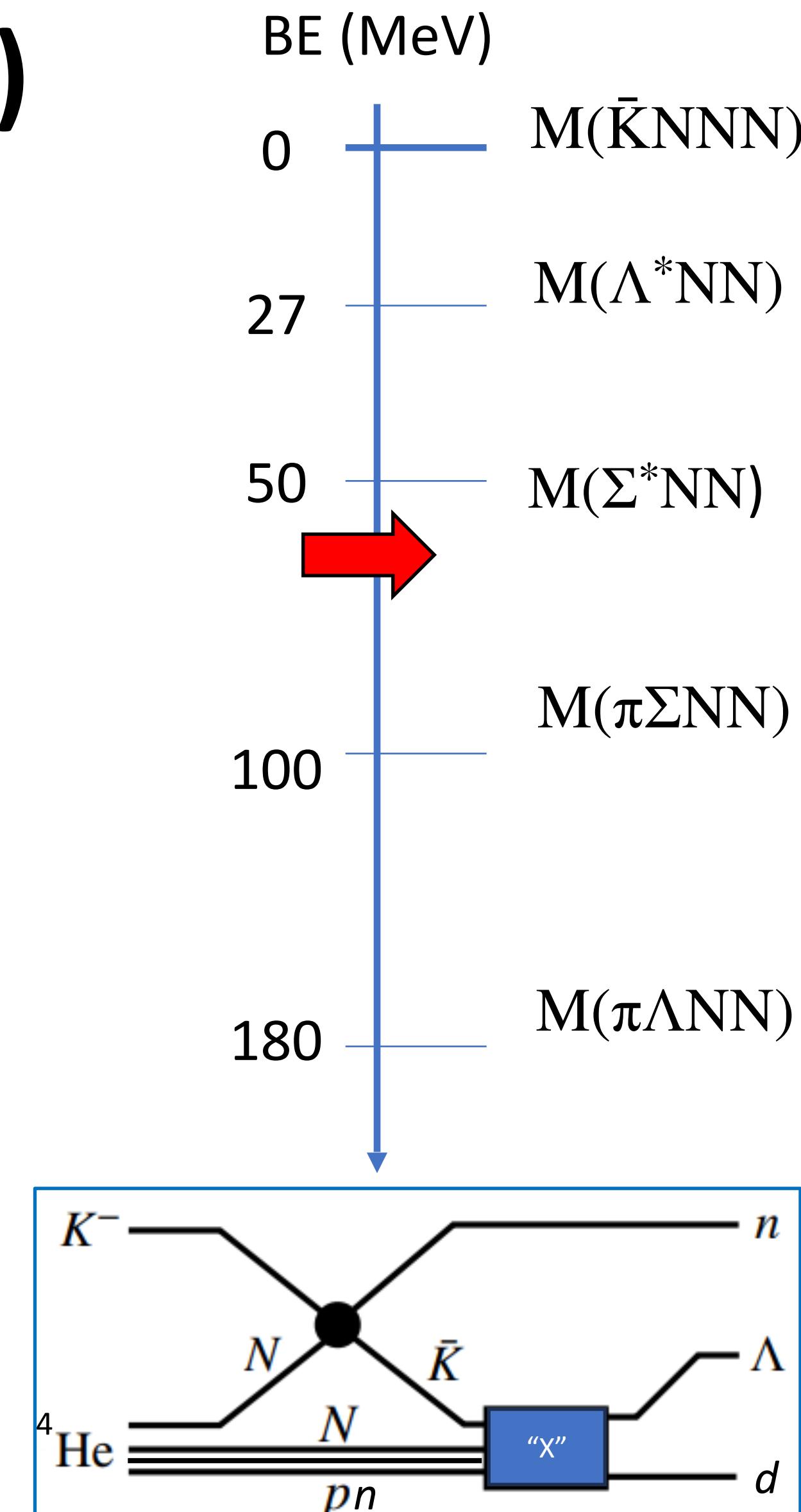
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2. If "X" = " $\bar{K}\text{-ppn}$ ", then J would be $J_{\bar{K}\text{-ppn}} = 1/2$, because J_{NNN}

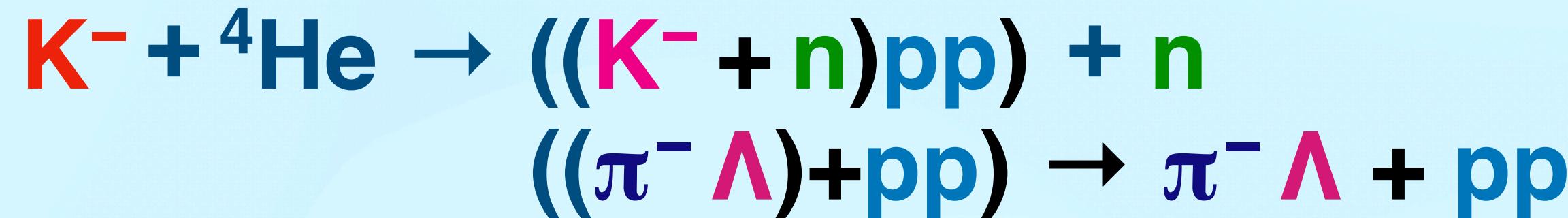
must be $1/2$ (two nucleons out of three are identical).

- **Most likely, it is "X" = " $\bar{K}\text{-ppn}$ ", in which \bar{K} 's isospin anti-parallel coupling with NNN [$I(J) = 1/2(1/2)$] in S-wave**
- **Difficult to interpret as $\Sigma^* NN$: 1) spin/isospin of NN [$I(J) = 1(0)$] part must flip to form deuteron [$I(J) = 0(1)$] without breaking the NN pair, while 2) the energy release of $\Sigma^* N \rightarrow \Lambda N$ is much bigger than the binding energy of d .**

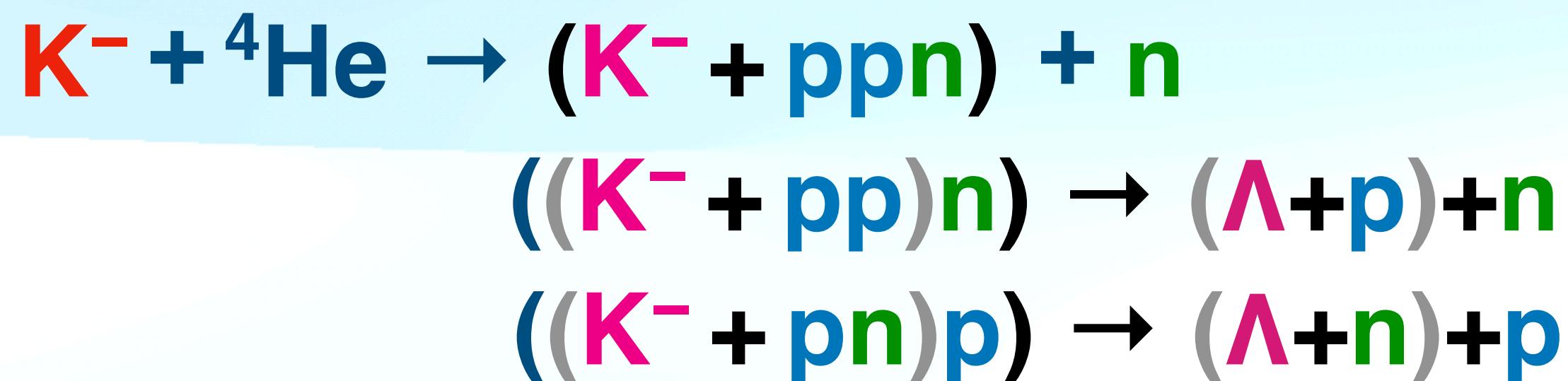


Size of $\bar{K}NNN$ via Data?

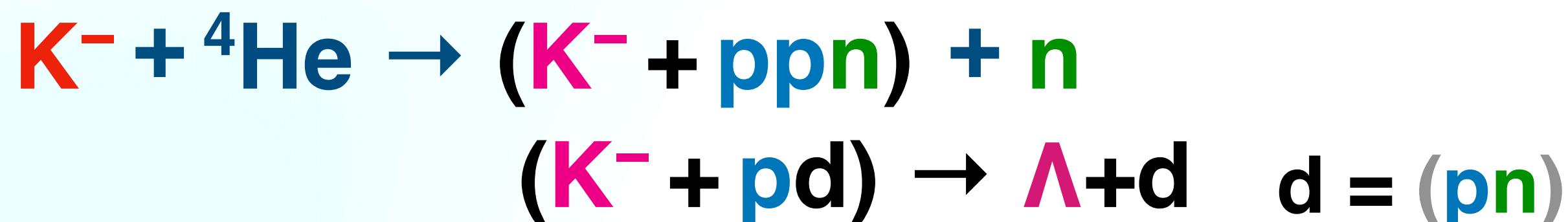
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The ratio of the three nucleon processes would be sensitive to the core size.

$$1N_{\bar{K}A} : 2N_{\bar{K}A} : 3N_{\bar{K}A} \sim \rho_N : \rho_N^2 : \rho_N^3 ?$$

Theoretical inputs are needed.

Two nucleon \bar{K} absorption ($2N_{\bar{K}A}$):

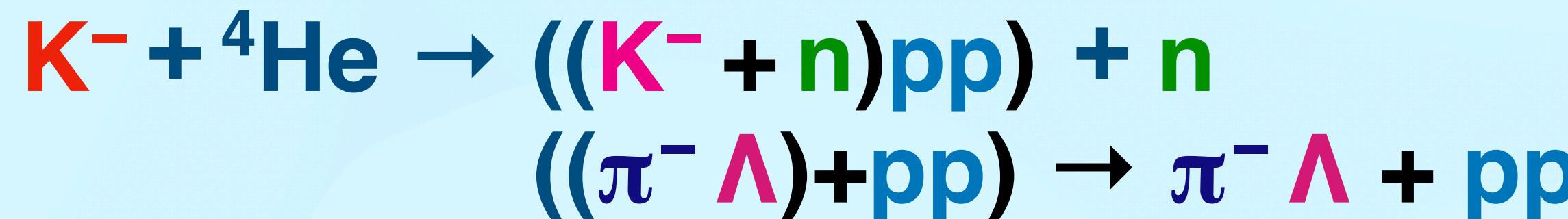


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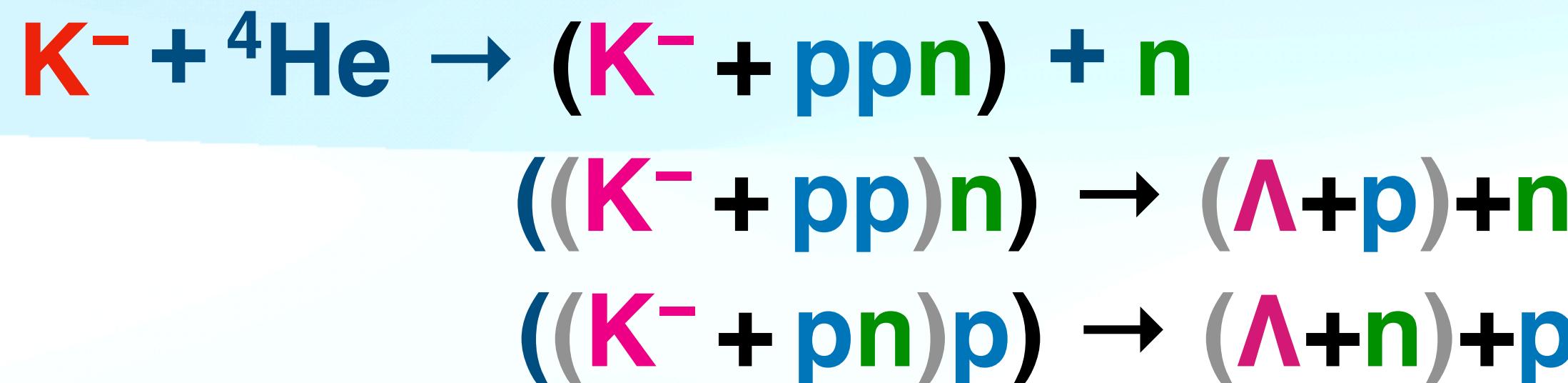


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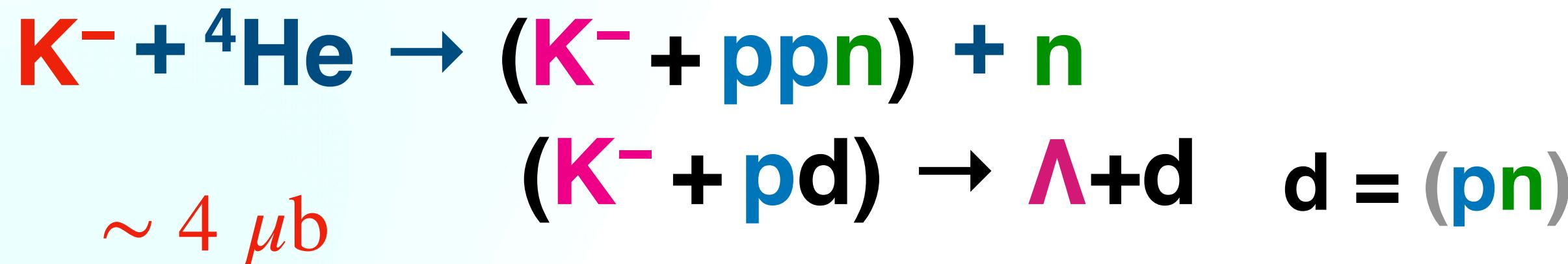
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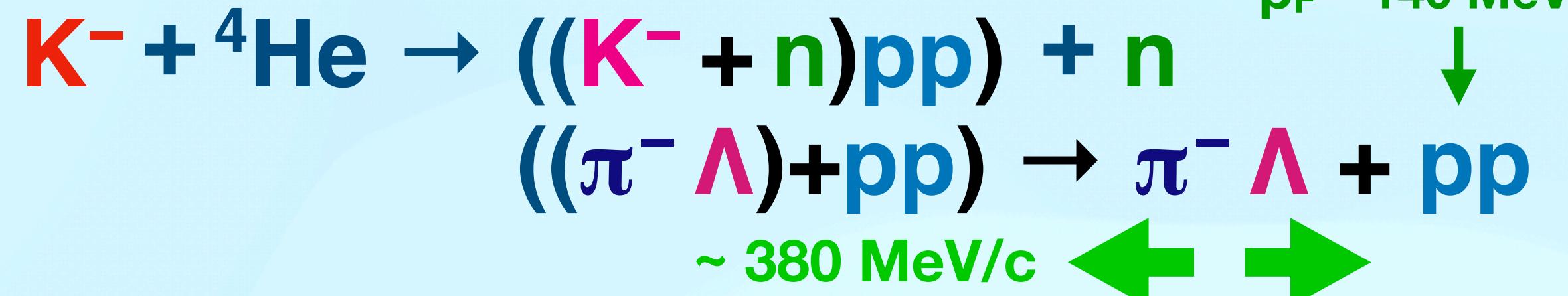


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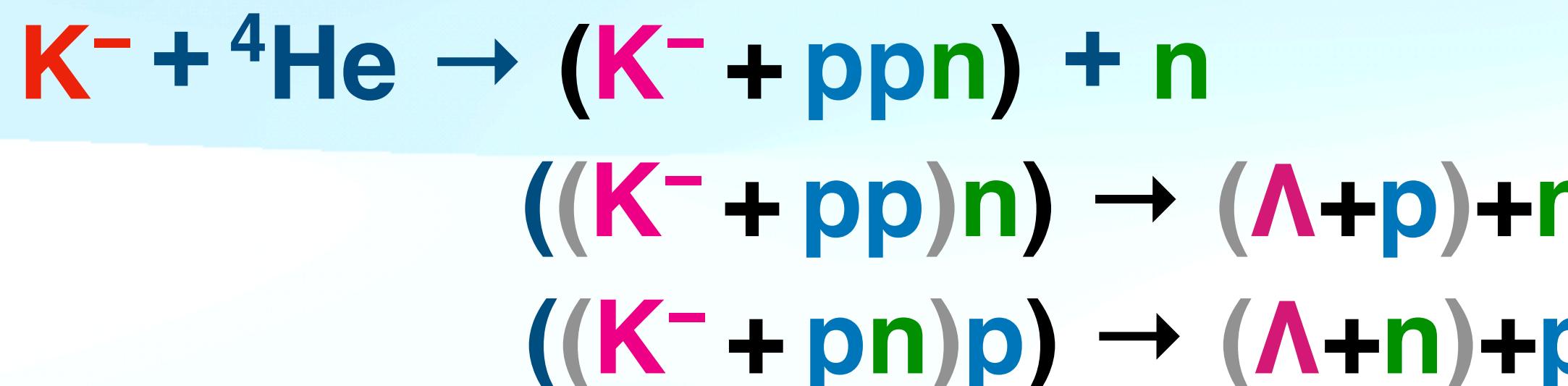


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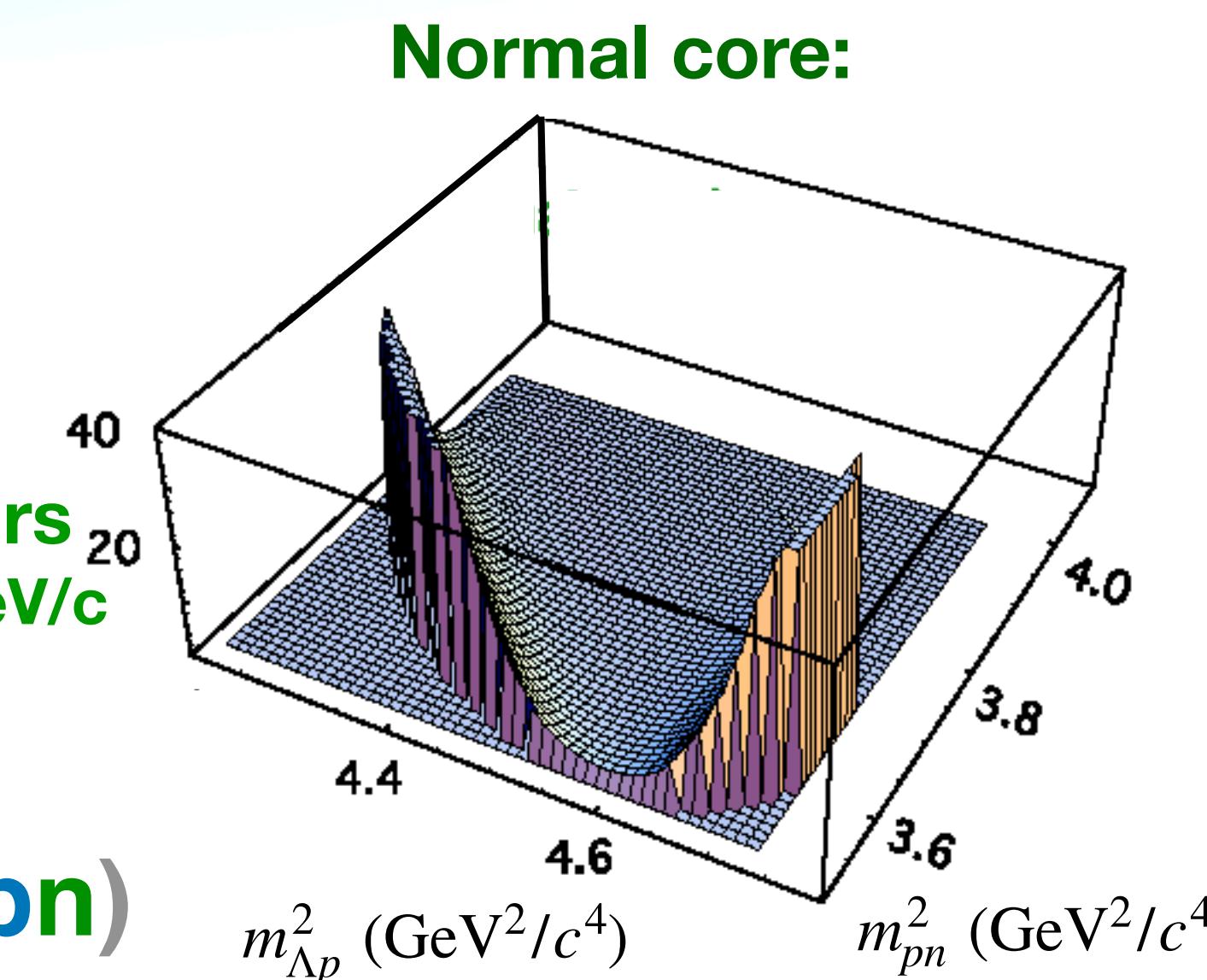
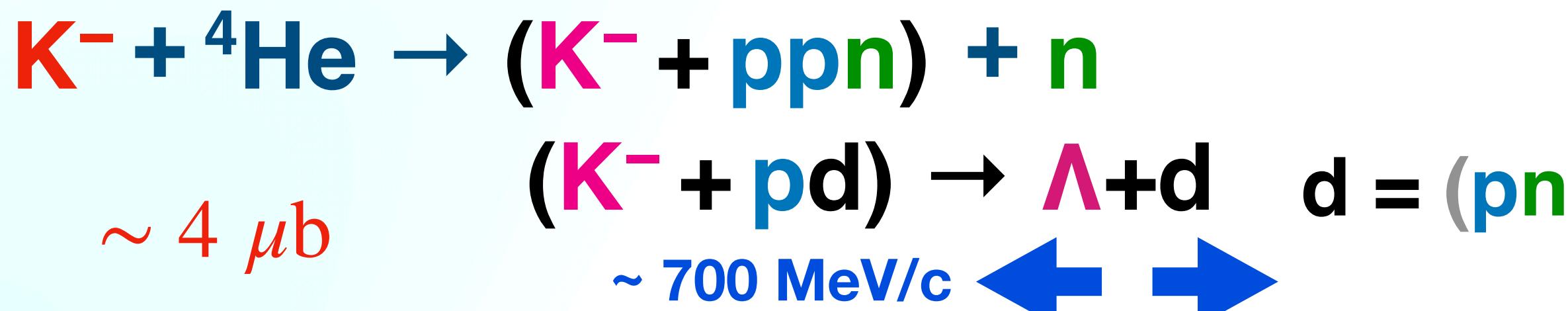
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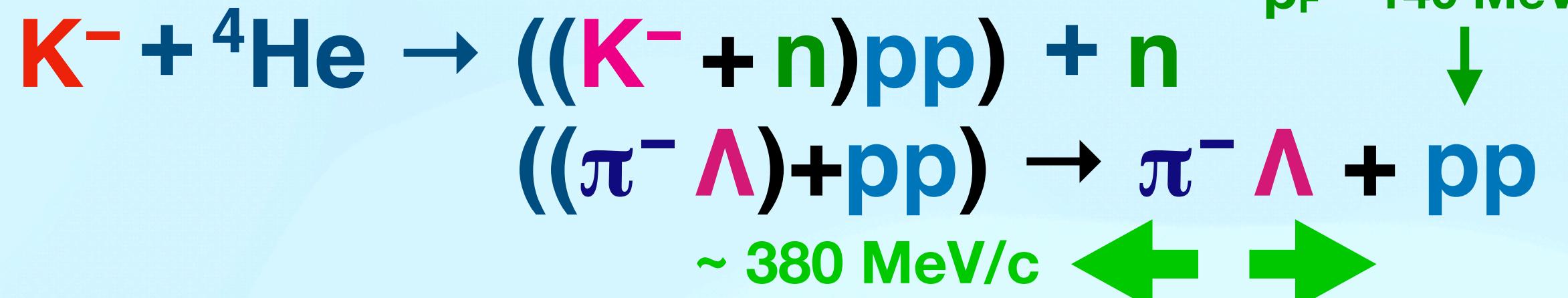
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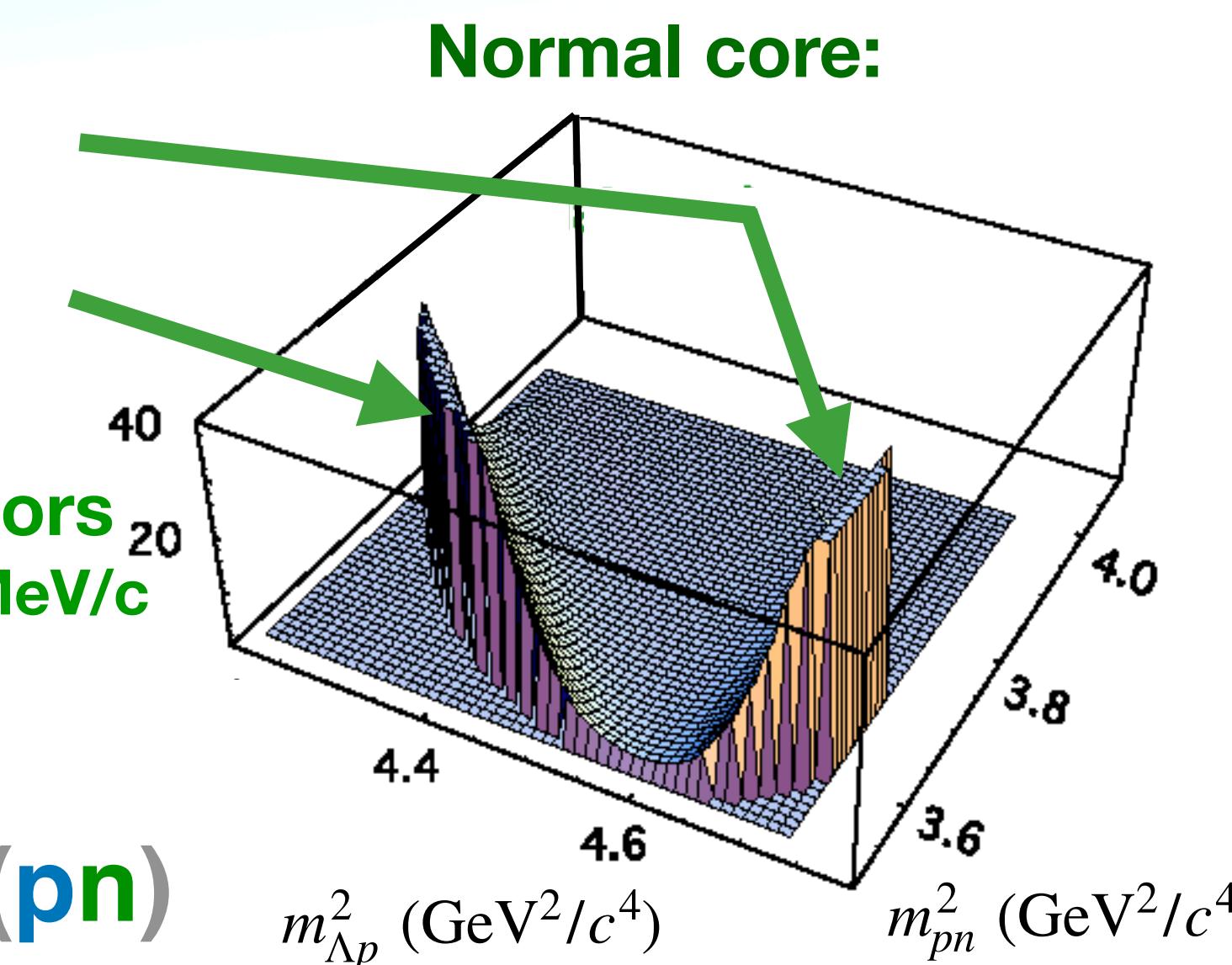
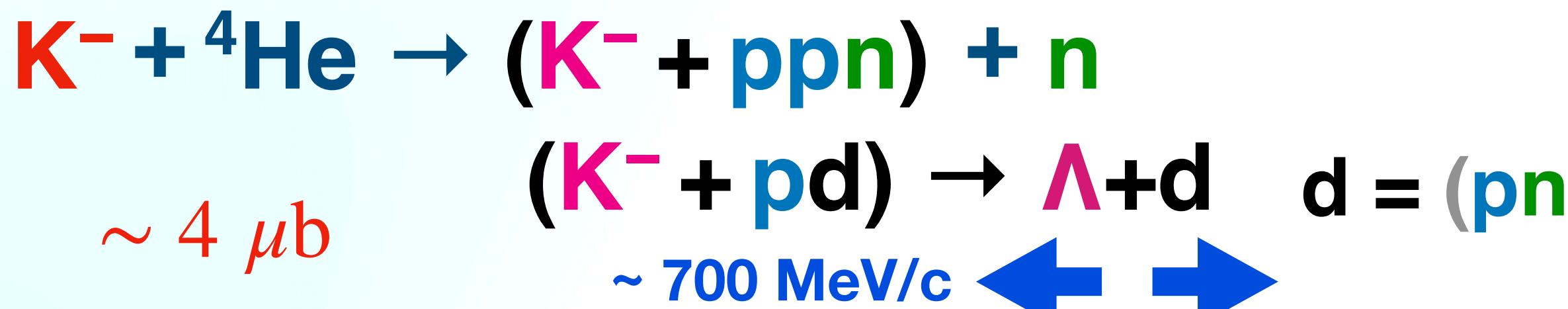
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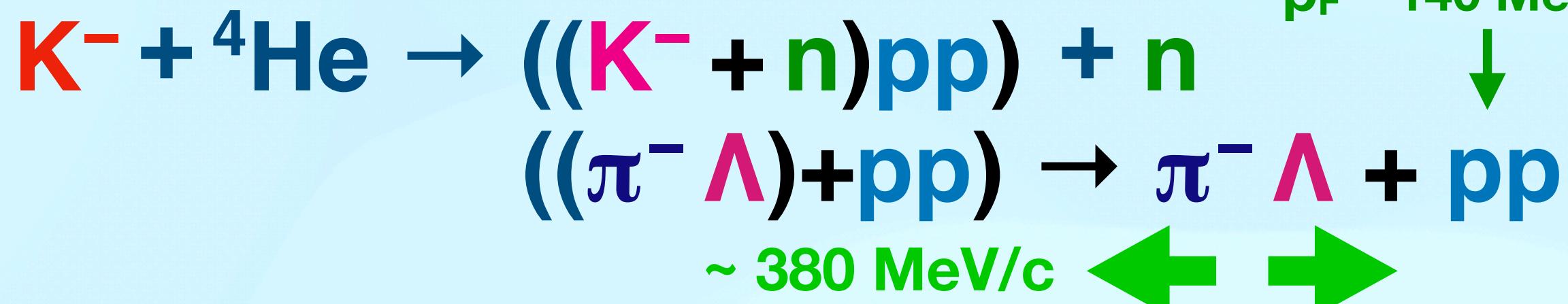
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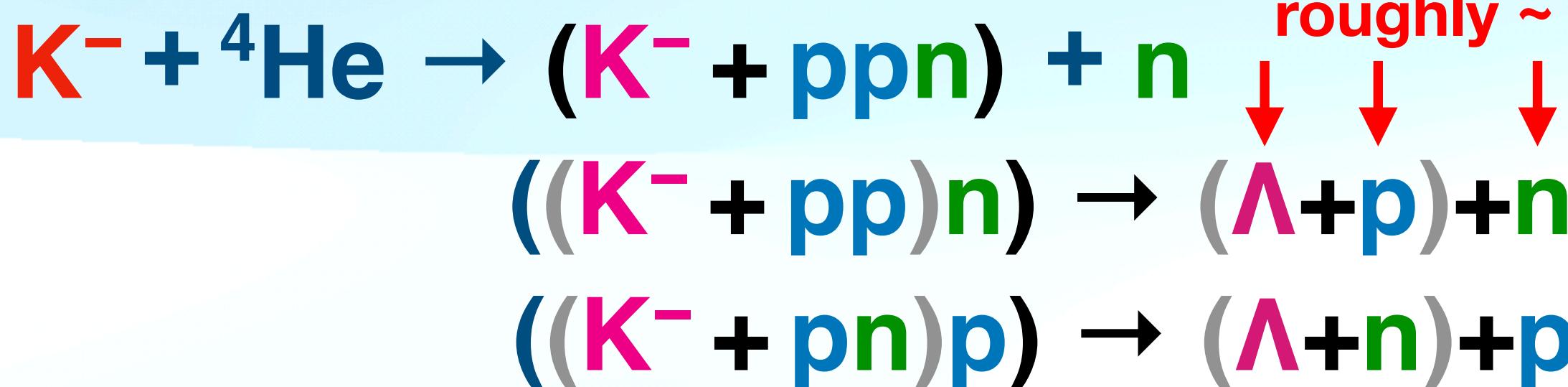


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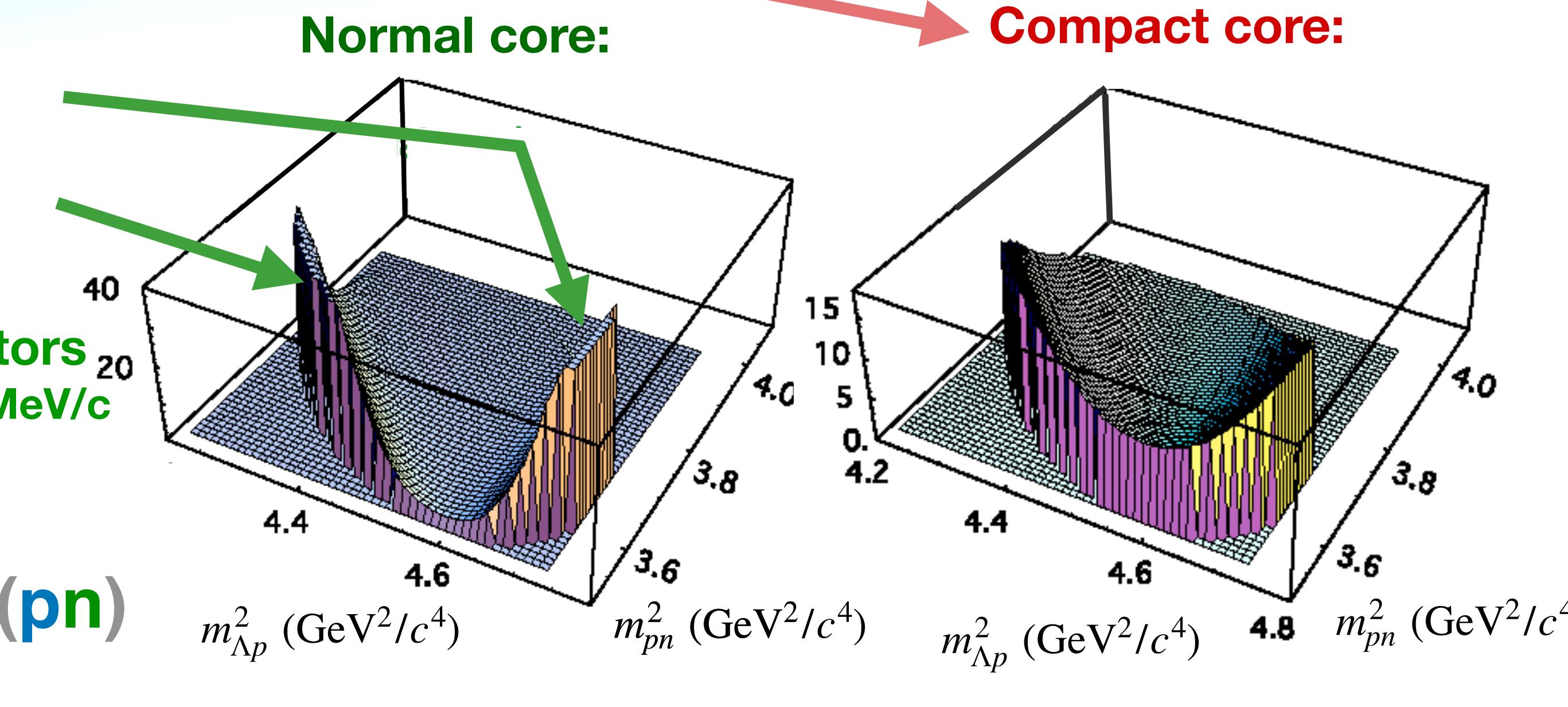
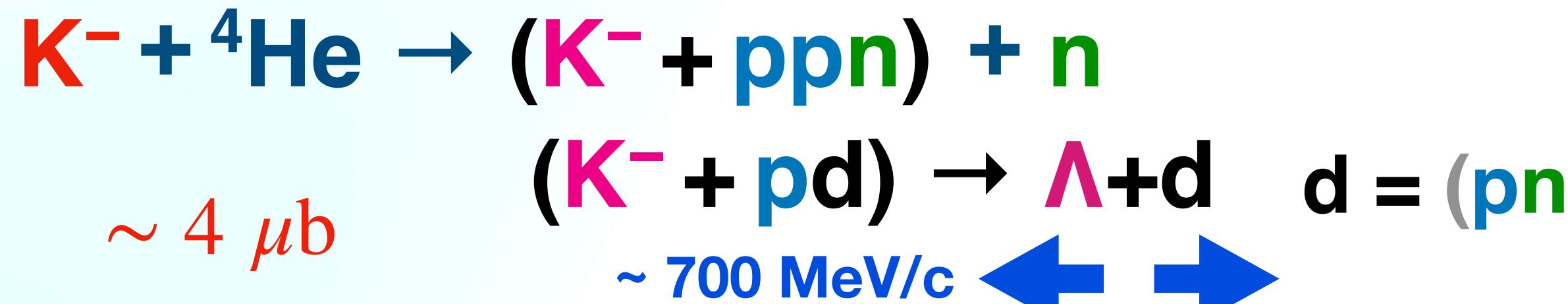
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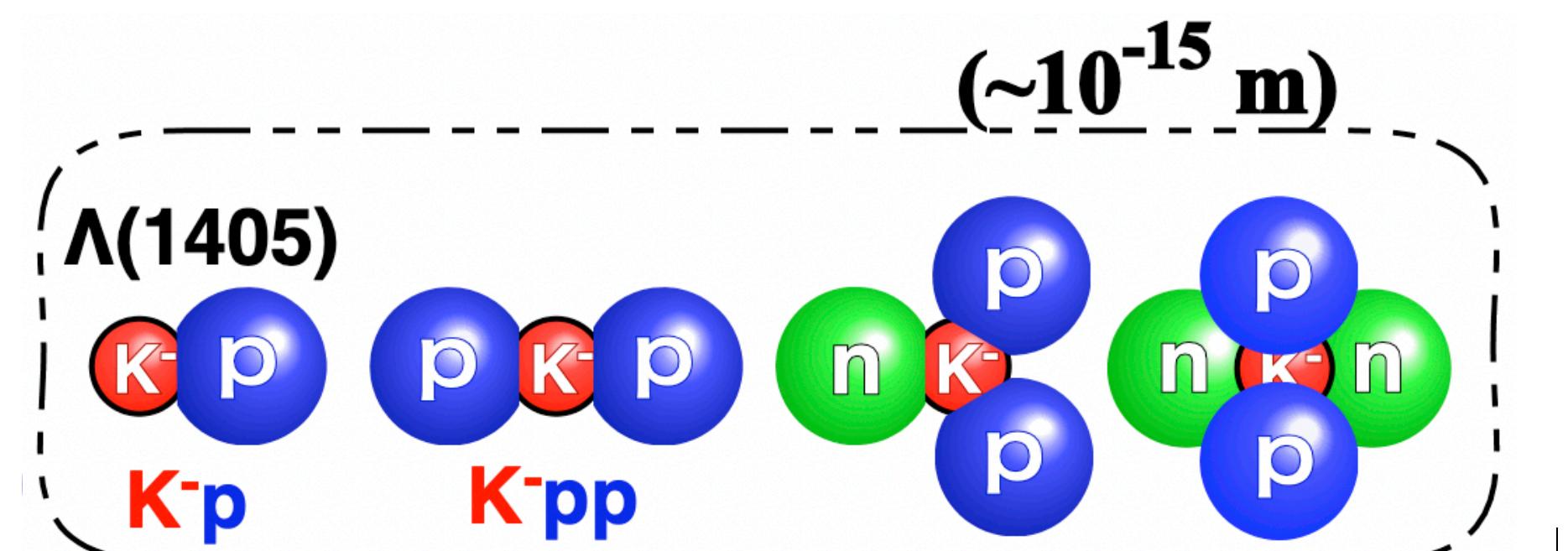
Toward next generation research – New Spectrometer –

Does Isospin-partner ($\bar{K}^0 nn$) _{$I_3 = -1/2$} exist?

Are kaonic nuclei really compact?

Theoretical exploration of Kaonic Nuclei.

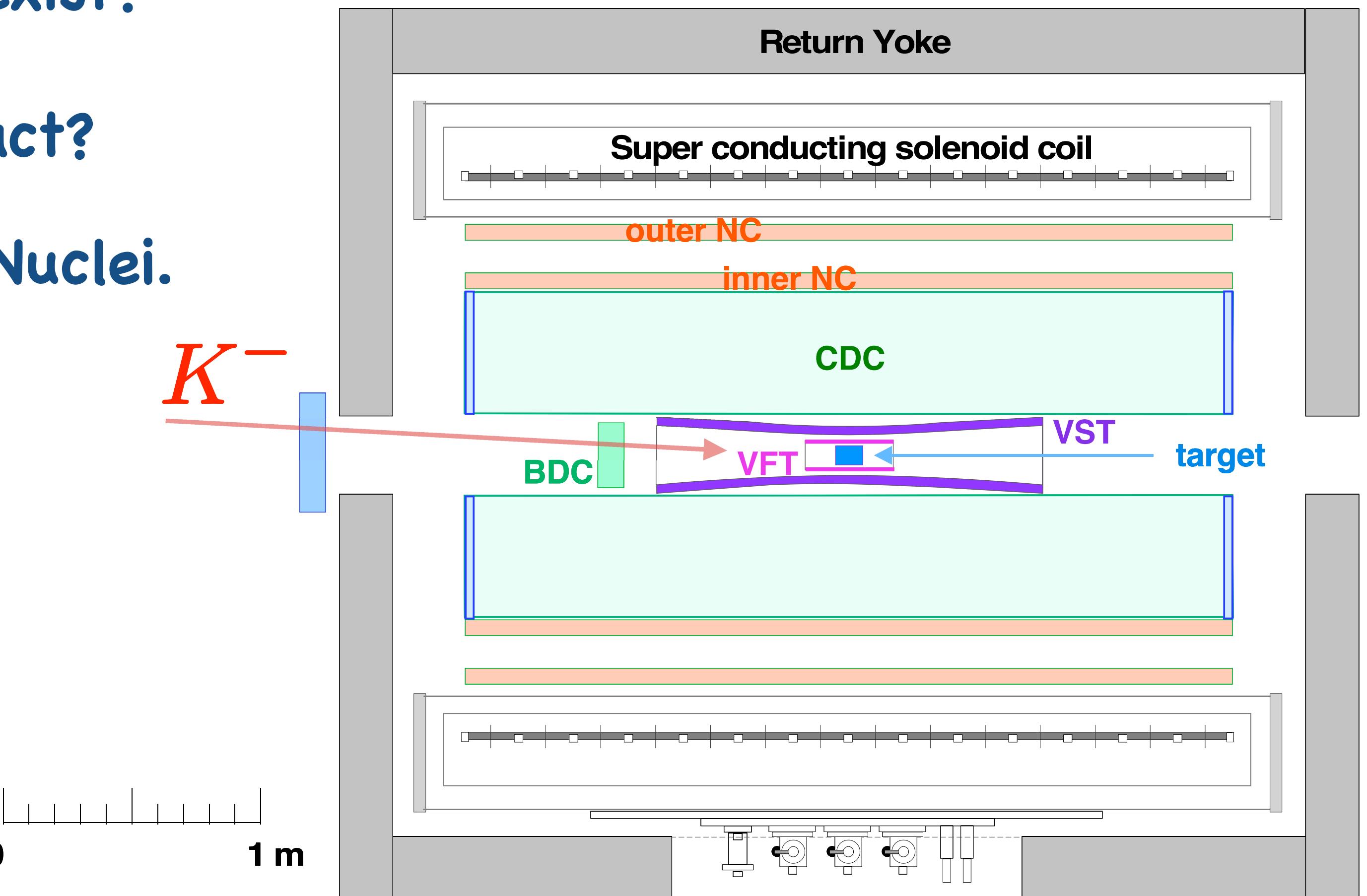
for the systematic study on



K^-

molecule-like hadronic nuclear cluster⁰

"Does it have a unique shape
like a chemical molecule?"



... Construction is in progress, led by *F. Sakuma* 25

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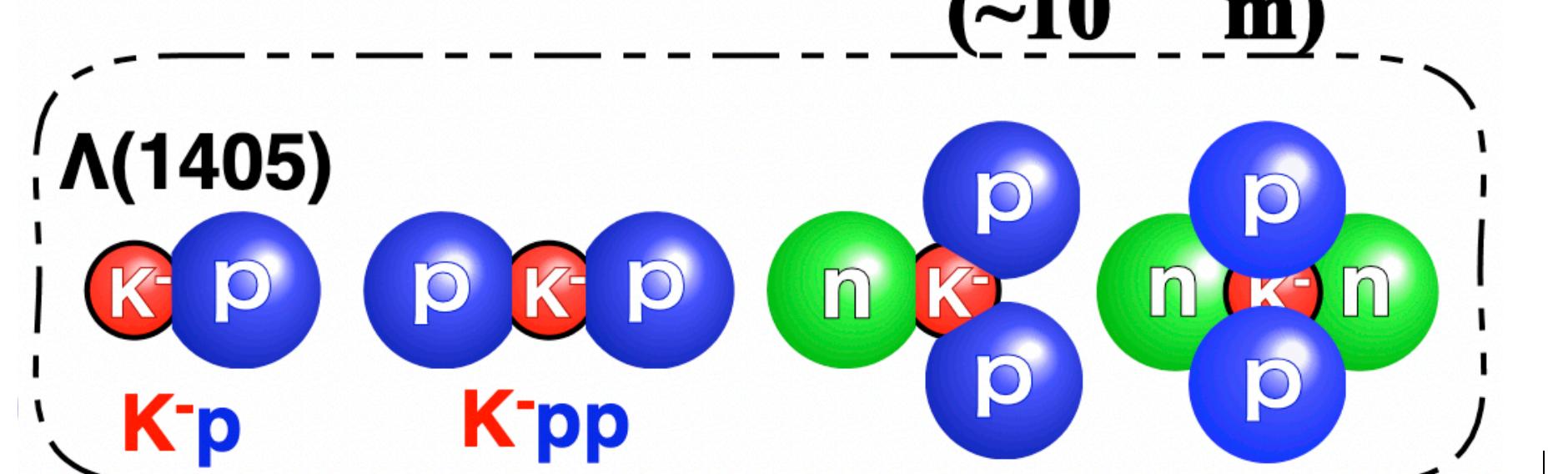
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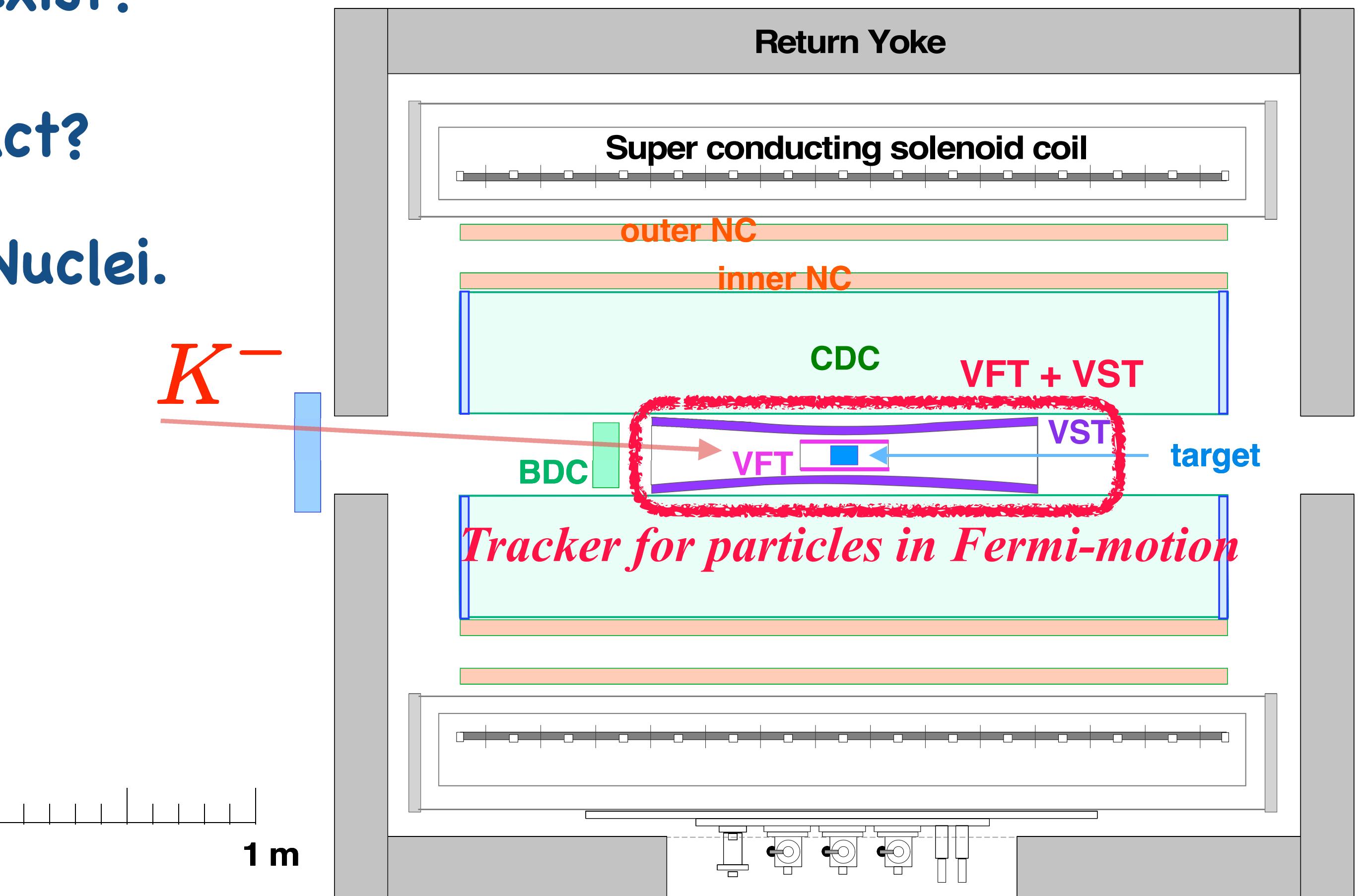
for the systematic study on

(~10⁻¹⁵ m)



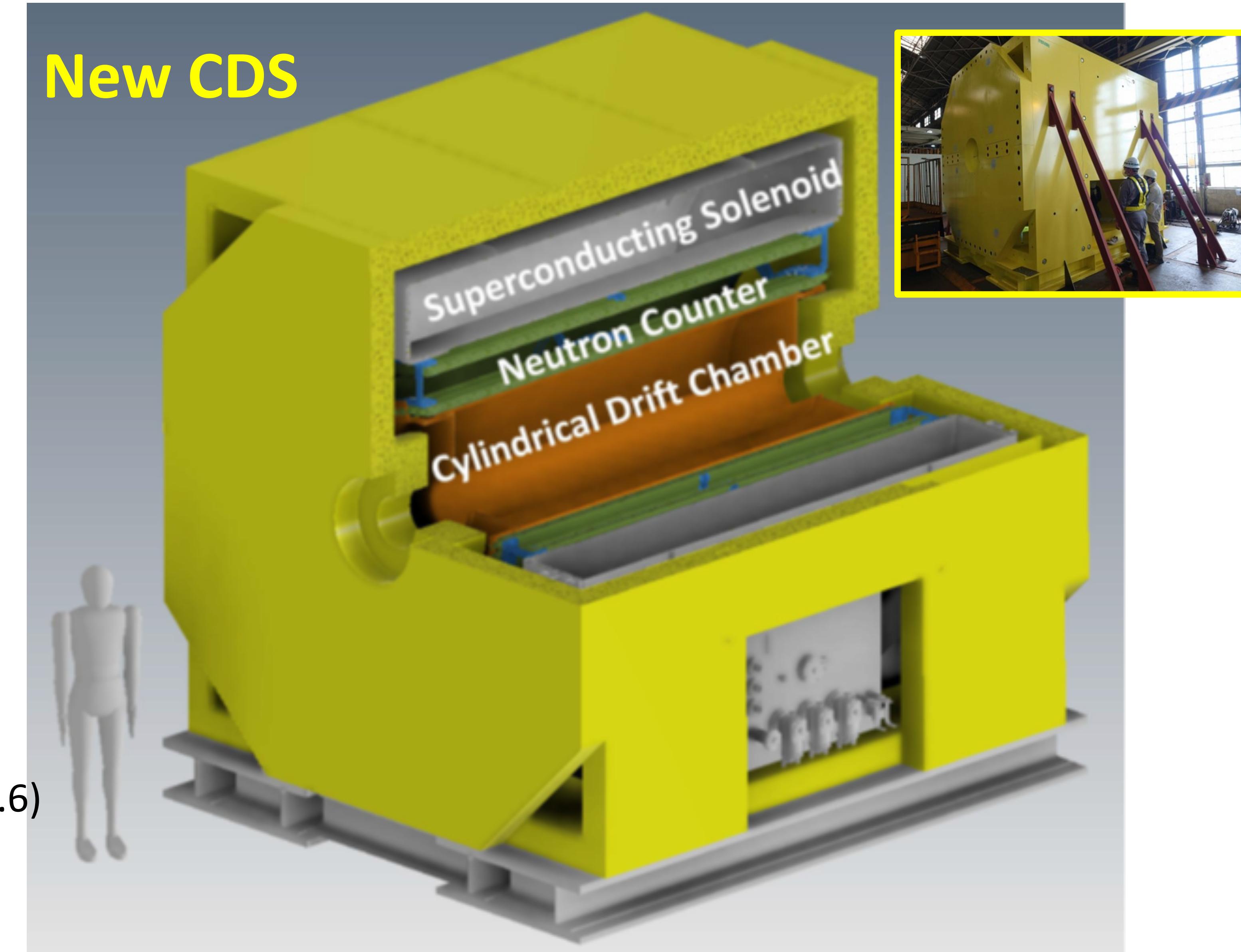
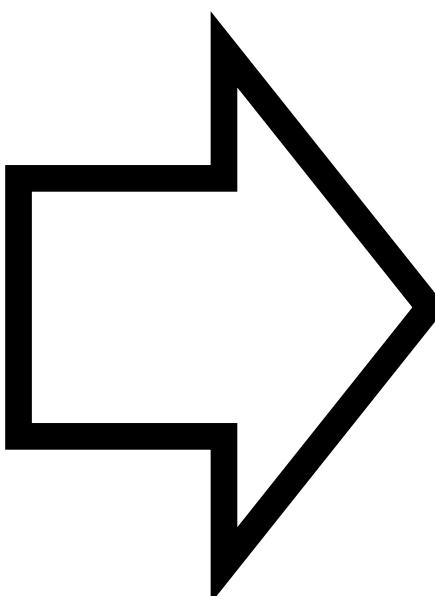
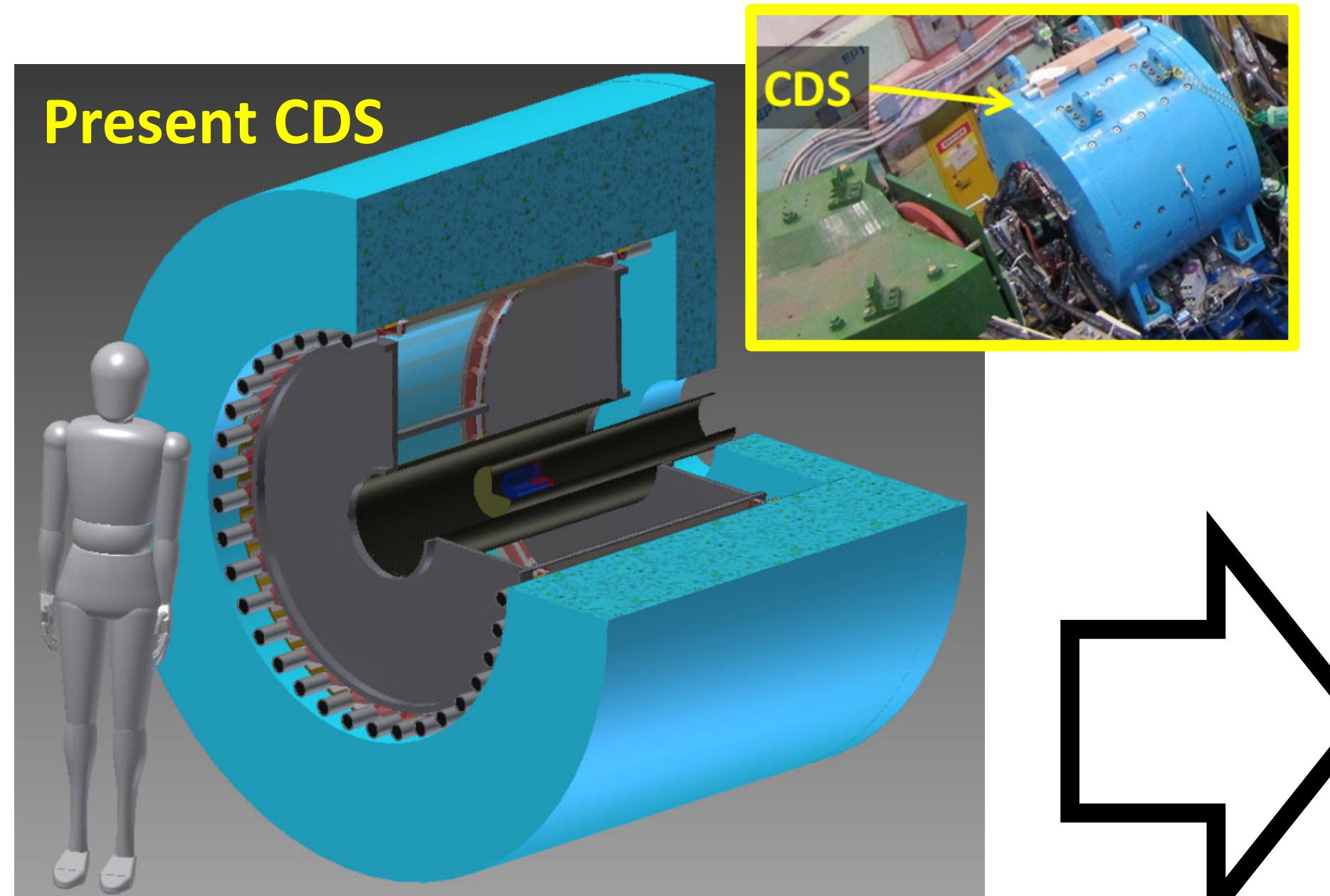
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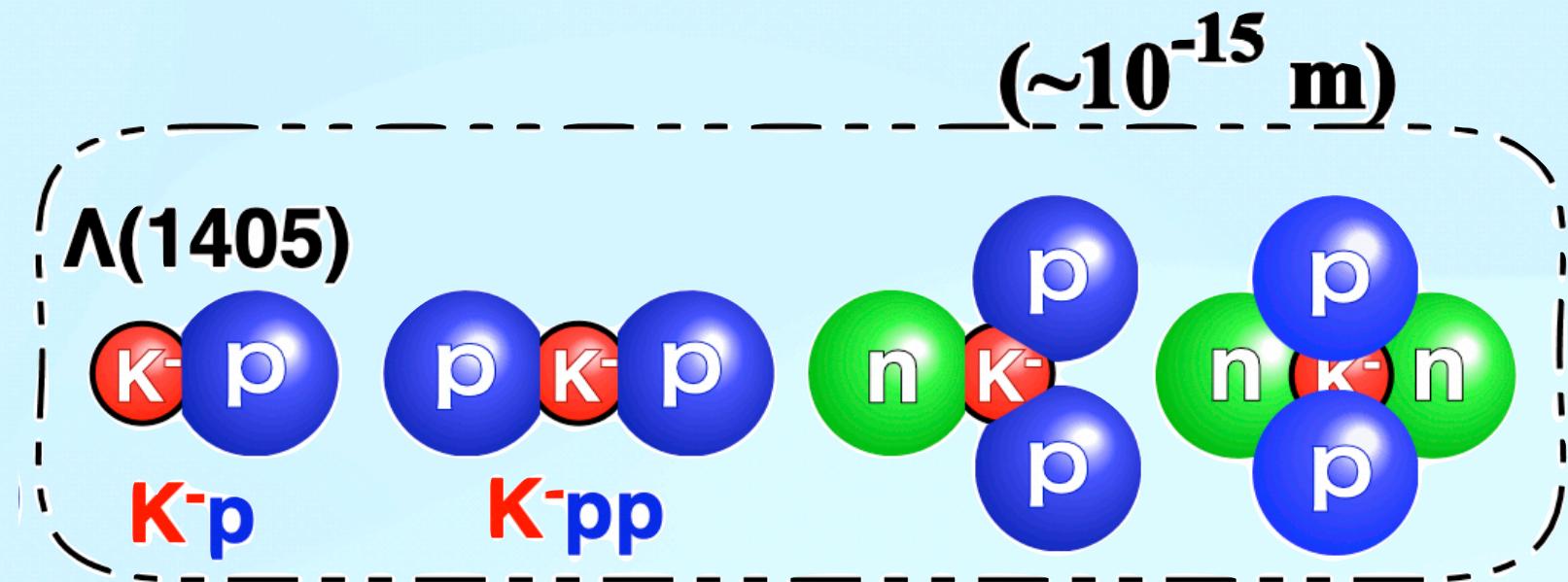
... Construction is in progress, led by *F. Sakuma* 25

New Cylindrical Detector System (CDS)



- ✓ Solid angle: x1.6 (59% → 93%)
- ✓ Neutron eff.: x7 (3% → 12% x 1.6)

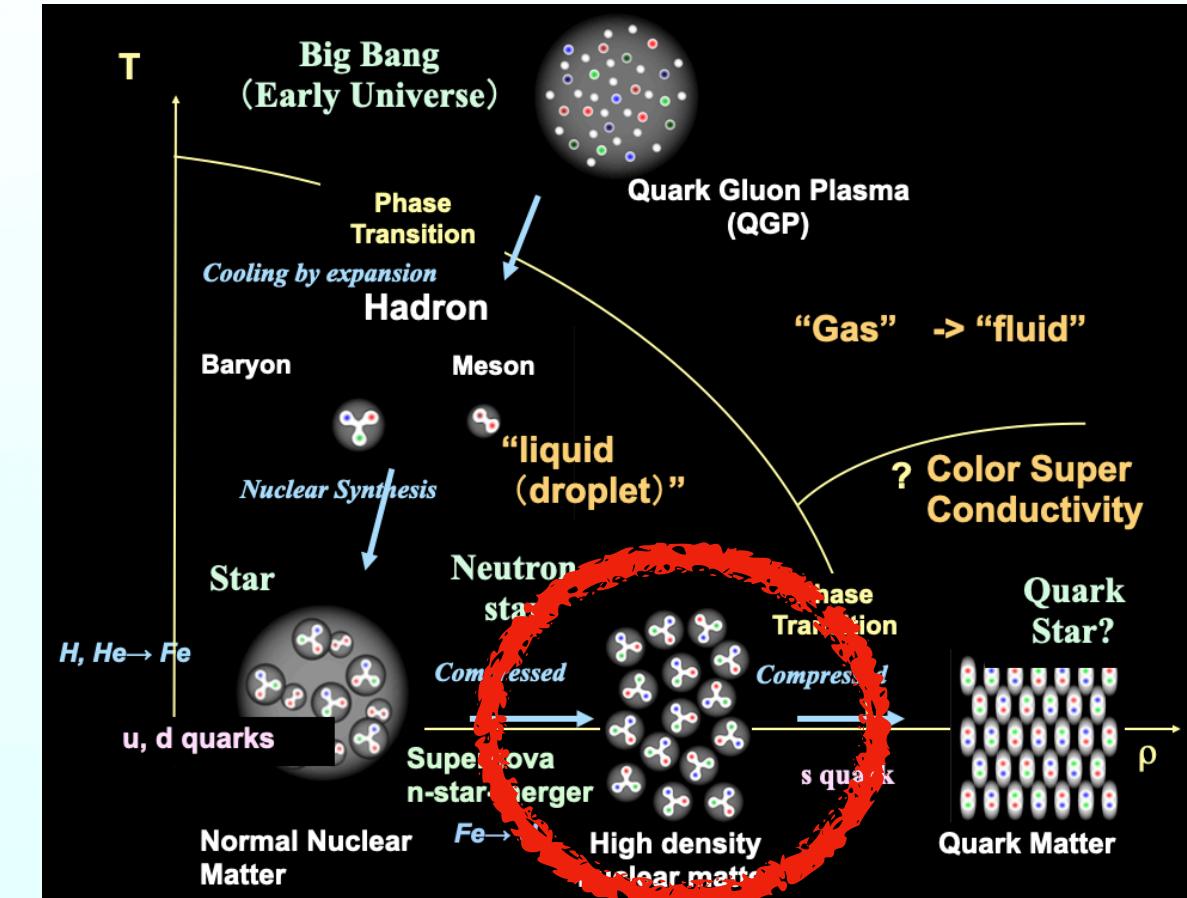
With new spectrometer, we will conduct — a systematic study on light kaonic nuclei —



"Does it have a unique shape like a chemical molecule?"

molecule-like hadronic nuclear cluster

How hadron mass is generated?



Physics at high density?



By establishing the $I(J^P)$, presence of " $K^- nn$ " — a Charge Mirror State of " $K^- pp$ " — and the Detailed Study of Three-Nucleon Kaon Nuclear Bound State as new starting points, we aim to open the door to research on high-density nuclear matter.

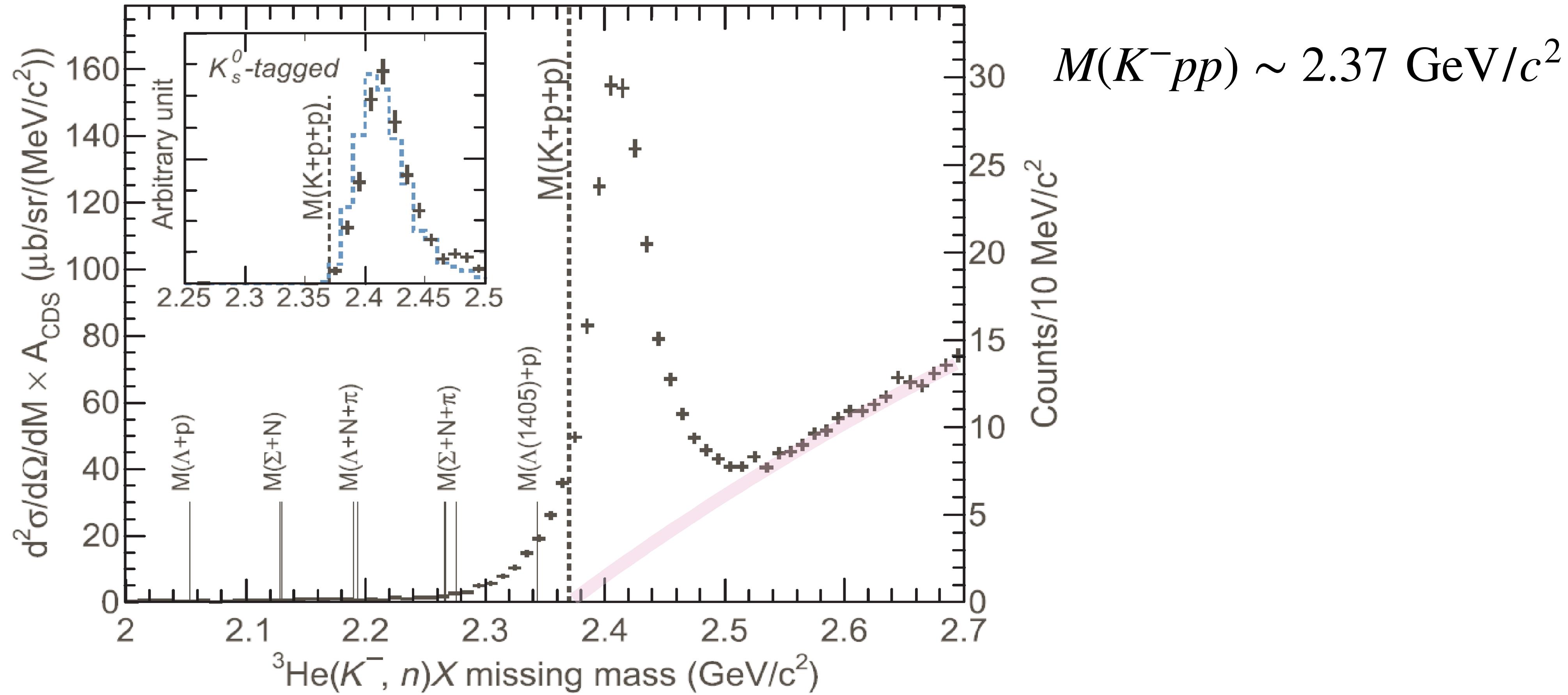
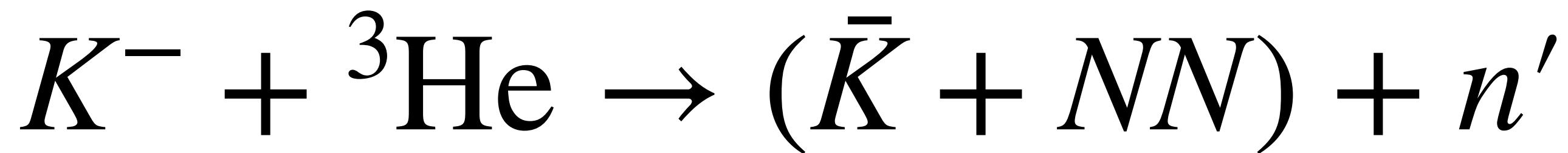


Hannes truly enjoyed his time in Japan, immersing himself in the physics research and embracing new experiences with enthusiasm.

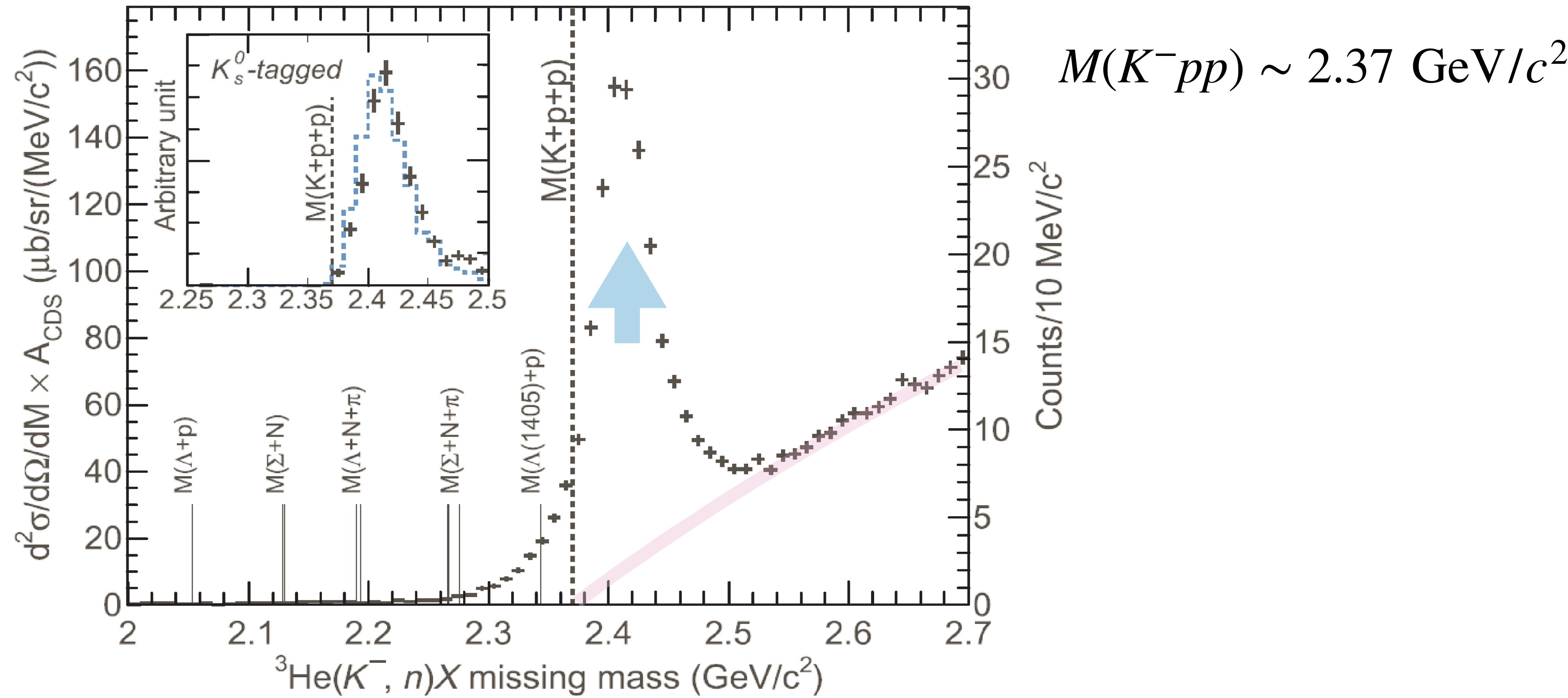
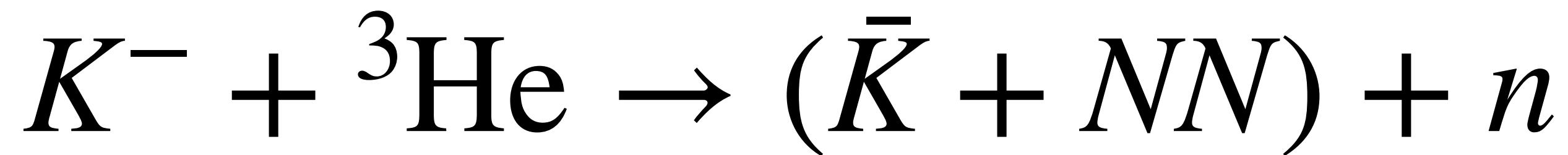
*As our field evolves, the need for long-term projects and collaborative efforts across multiple disciplines has become increasingly vital. Inspired by Hannes's passion for discovery, I am committed to continuing this valuable research and building on what we've accomplished together. I would like to extend an invitation to anyone interested in embarking on new collaborations. Together, we can explore new frontiers in physics and create truly remarkable **Hadronic-Molecule Study – Kaonic Nuclear Bound State**. If you share this vision, I encourage you to join us in this exciting journey!*

Thank you for your attention!

${}^3\text{He}(K^-, n_{\text{NC}})X$ – semi-inclusive

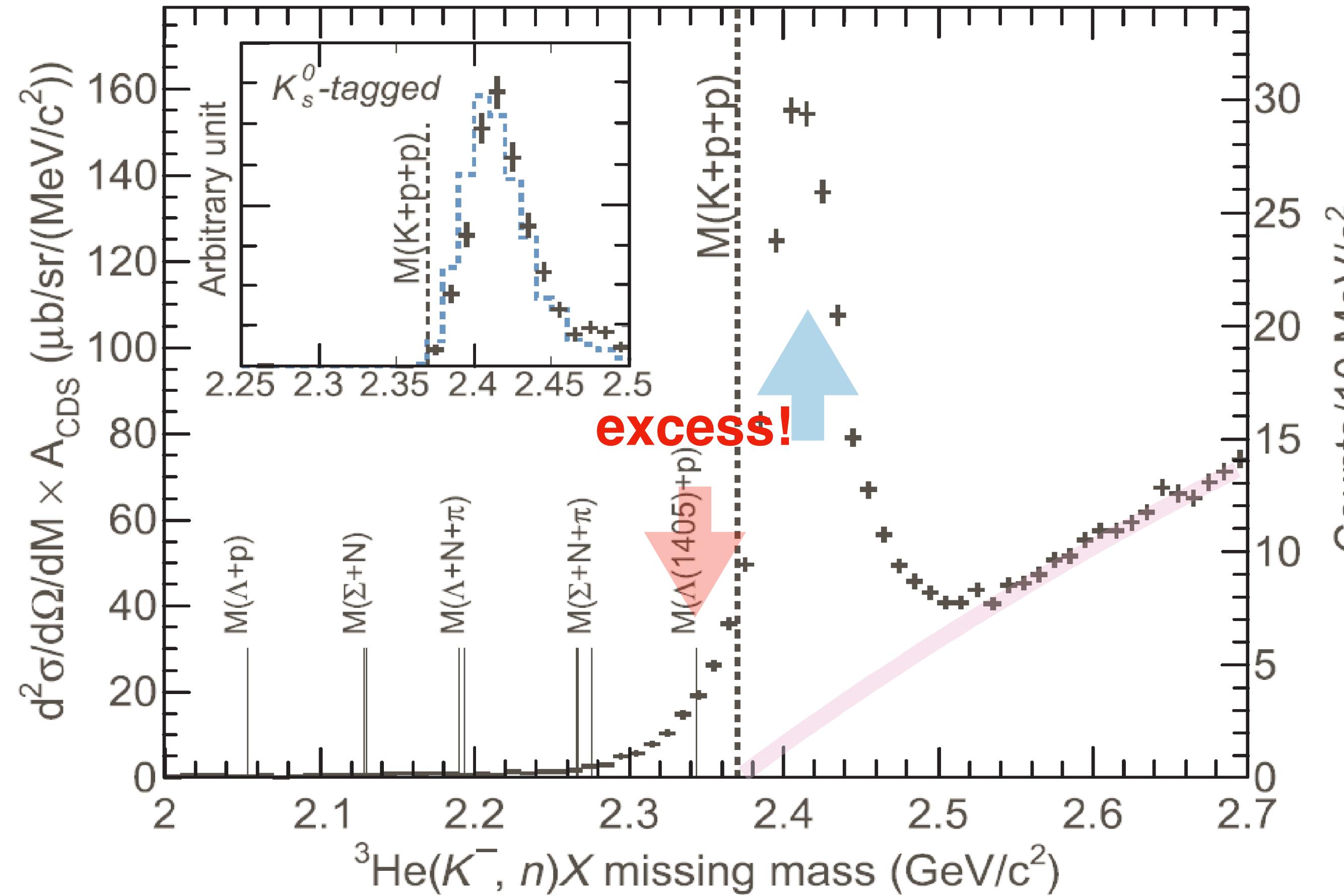
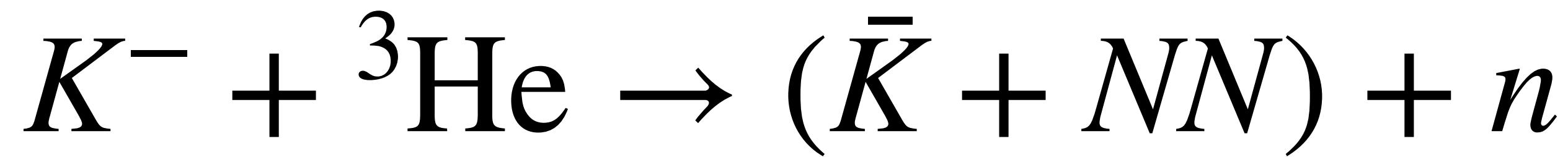


${}^3\text{He}(K^-, n_{\text{NC}})X$ – semi-inclusive



A nucleon knockout reaction $K^-N \rightarrow \bar{K}n'$ is the dominant reaction process

${}^3\text{He}(K^-, n_{\text{NC}})X$ – semi-inclusive

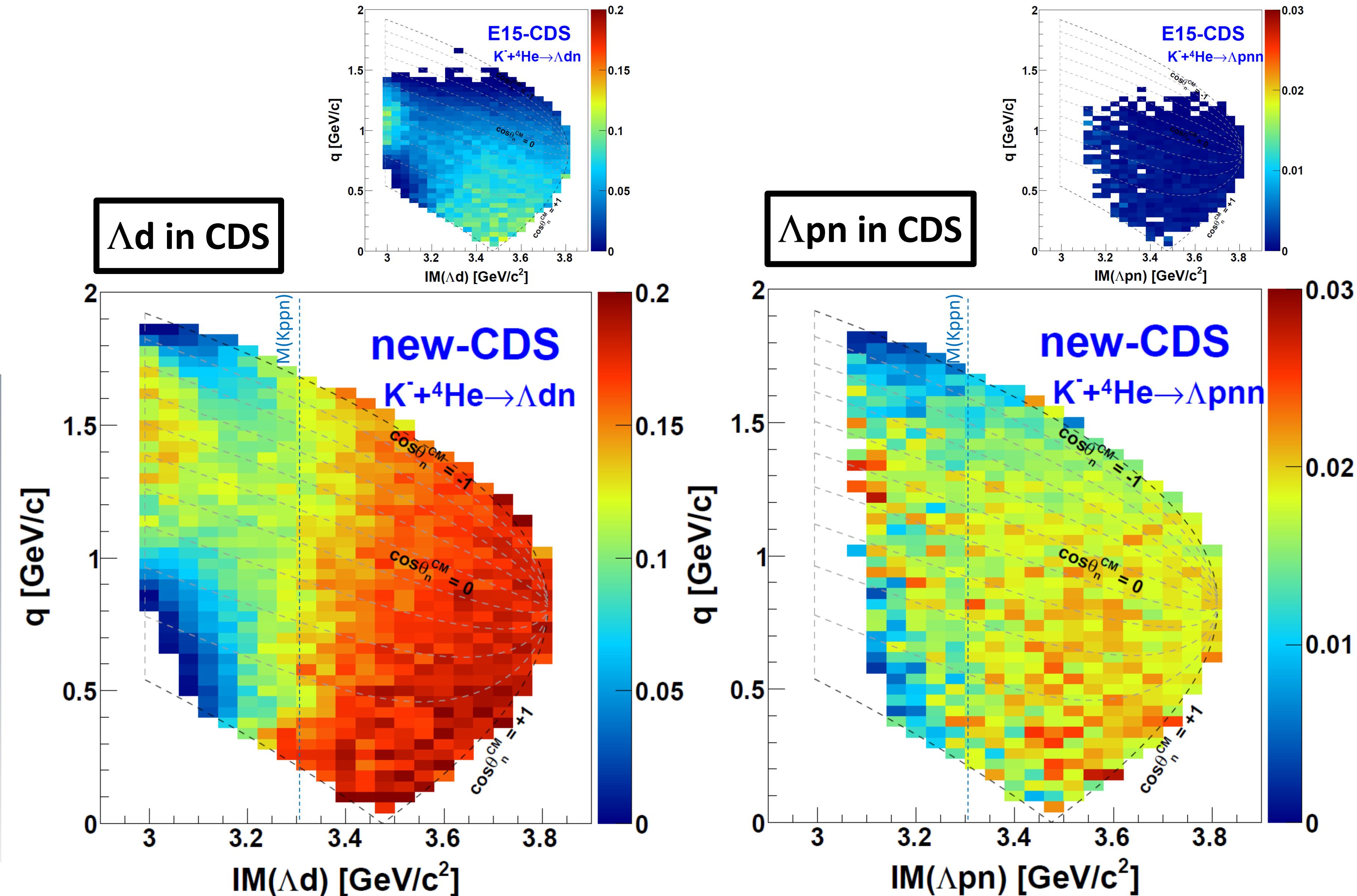
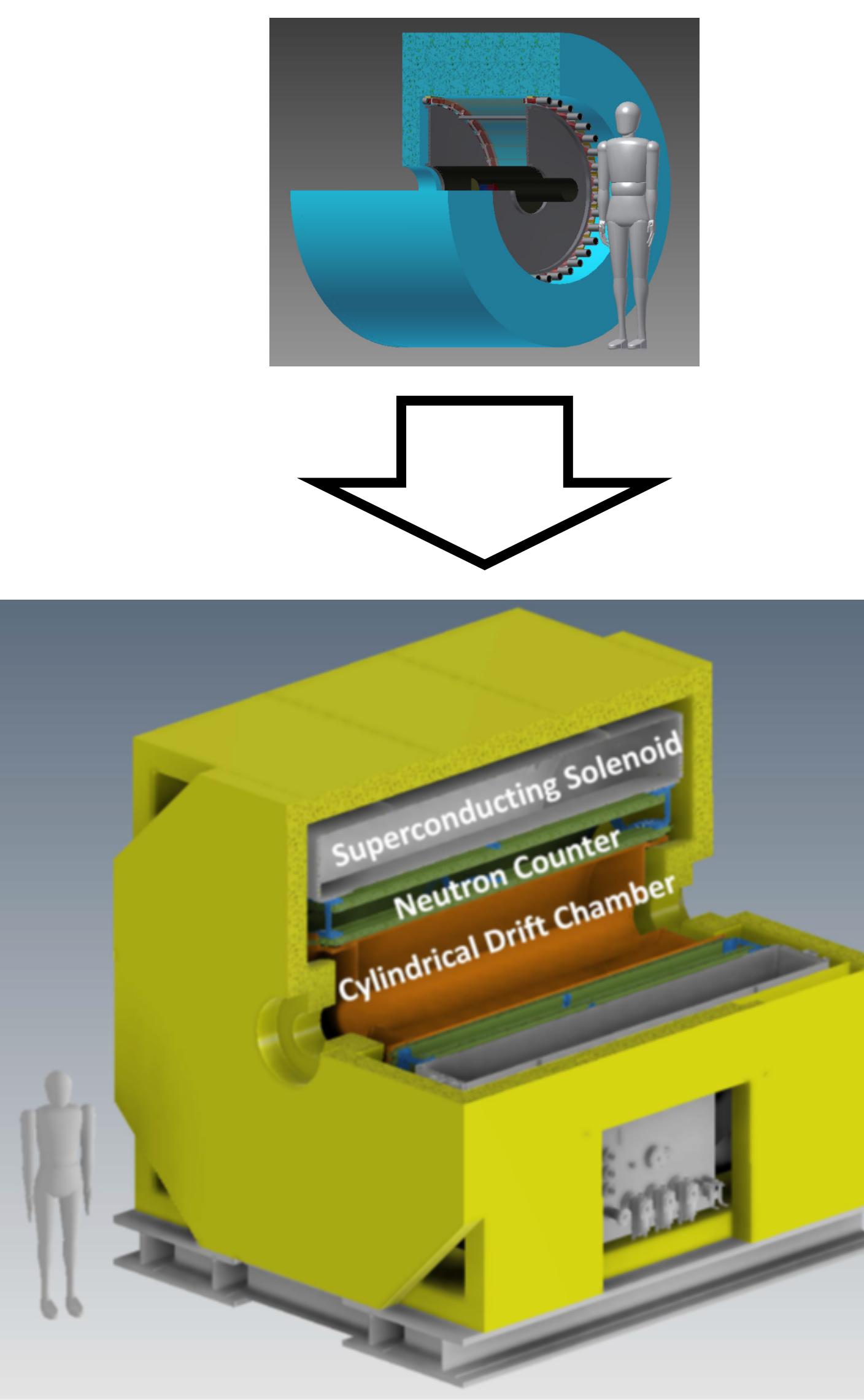


$$M(K^-pp) \sim 2.37 \text{ GeV}/c^2$$

How to study excess:
 $\bar{K} + NN \rightarrow \Lambda p$ happens
only when all the
particles are in the
strong interaction range,
because of energy-
momentum mismatch

A nucleon knockout reaction $K^-N \rightarrow \bar{K}n'$ is the dominant reaction process

Acceptance for $K^- + ^4\text{He} \rightarrow \Lambda d n$ reaction



Possible $I(J^P)$?

$\bar{K}NN : J^P = 0^-, I = 1/2 : I_{NN} = 1, S_{NN} = 0, L_{\bar{K}} = 0$

nucleon isospin symmetric ($I_{NN} = 1$) and spin anti-symmetric ($S_{NN} = 0$)

$\bar{K}NN : J^P = 1^-, I = 1/2 : I_{NN} = 0, S_{NN} = 1, L_{\bar{K}} = 0$

nucleon isospin anti-symmetric ($I_{NN} = 0$) and spin symmetric ($S_{NN} = 1$)

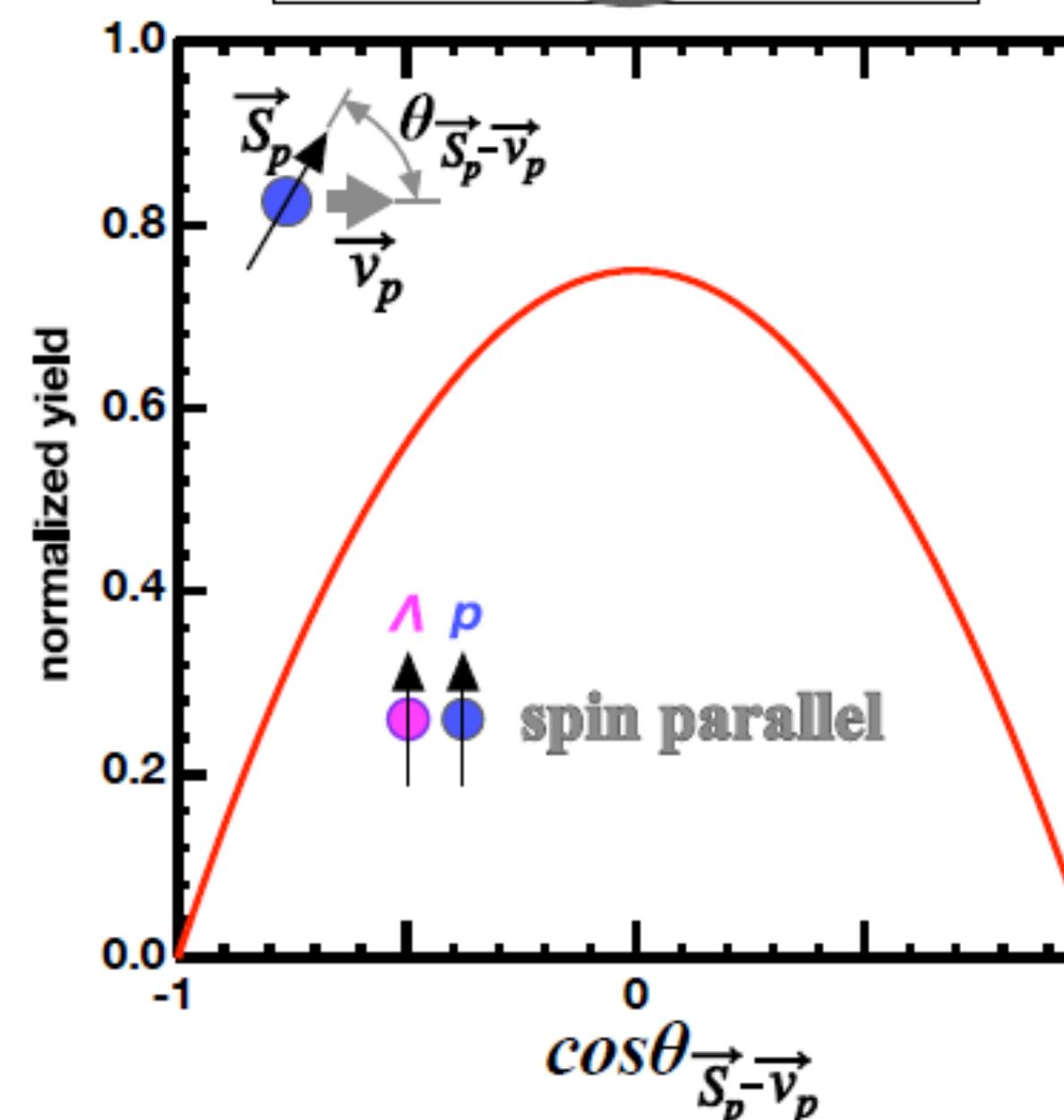
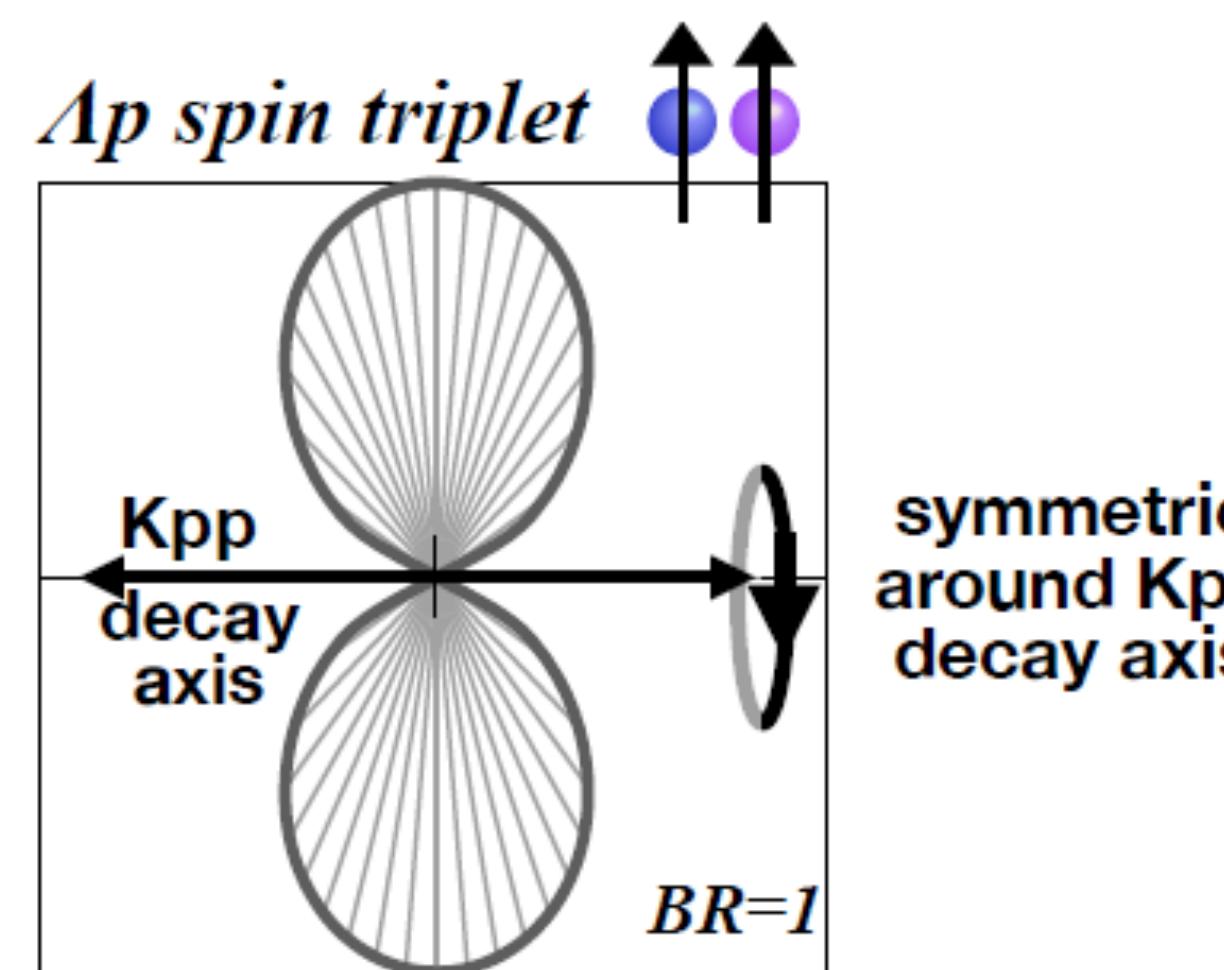
$I(\bar{K}NN) / J^P(\bar{K}NN)$	$(1/2) / (0^-)$	$(1/2) / (1^-)$
NN symmetry	$I(NN) = 1, S(NN) = 0$	$I(NN) = 0, S(NN) = 1$
“ $K^- pp$ ” $I_3(\bar{K}NN) = +\frac{1}{2}$	$-\sqrt{\frac{1}{3}} \left(\sqrt{2} K^- pp + \bar{K}^0 \frac{pn + np}{\sqrt{2}} \right) \otimes \left(\frac{\uparrow\downarrow - \downarrow\uparrow}{\sqrt{2}} \right)$	$\bar{K}^0 \frac{(pn - np)}{\sqrt{2}} \otimes \left(\uparrow\uparrow, \frac{\uparrow\downarrow + \downarrow\uparrow}{\sqrt{2}}, \downarrow\downarrow \right)$
“ $\bar{K}^0 nn$ ” $I_3(\bar{K}NN) = -\frac{1}{2}$	$-\sqrt{\frac{1}{3}} \left(\sqrt{2} \bar{K}^0 nn + K^- \frac{pn + np}{\sqrt{2}} \right) \otimes \left(\frac{\uparrow\downarrow - \downarrow\uparrow}{\sqrt{2}} \right)$	$-K^- \frac{(pn - np)}{\sqrt{2}} \otimes \left(\uparrow\uparrow, \frac{\uparrow\downarrow + \downarrow\uparrow}{\sqrt{2}}, \downarrow\downarrow \right)$
$\bar{K}N$ coupling	$\frac{ I_{\bar{K}N} = 0 ^2}{ I_{\bar{K}N} = 1 ^2} = \frac{3}{1}$	$\frac{ I_{\bar{K}N} = 0 ^2}{ I_{\bar{K}N} = 1 ^2} = \frac{1}{3}$
$\frac{\sigma_{\bar{K}^0 nn}}{\sigma_{K^- pp}}$	$0.13 \sim 0.15$	~ 0.75

“ $K^- pp$ ” $\rightarrow \Lambda p$ requires
 $I = 1/2$, presence of
kaon requires negative
parity, and the Λp
decay must be in P-wave
due to the negative
parity

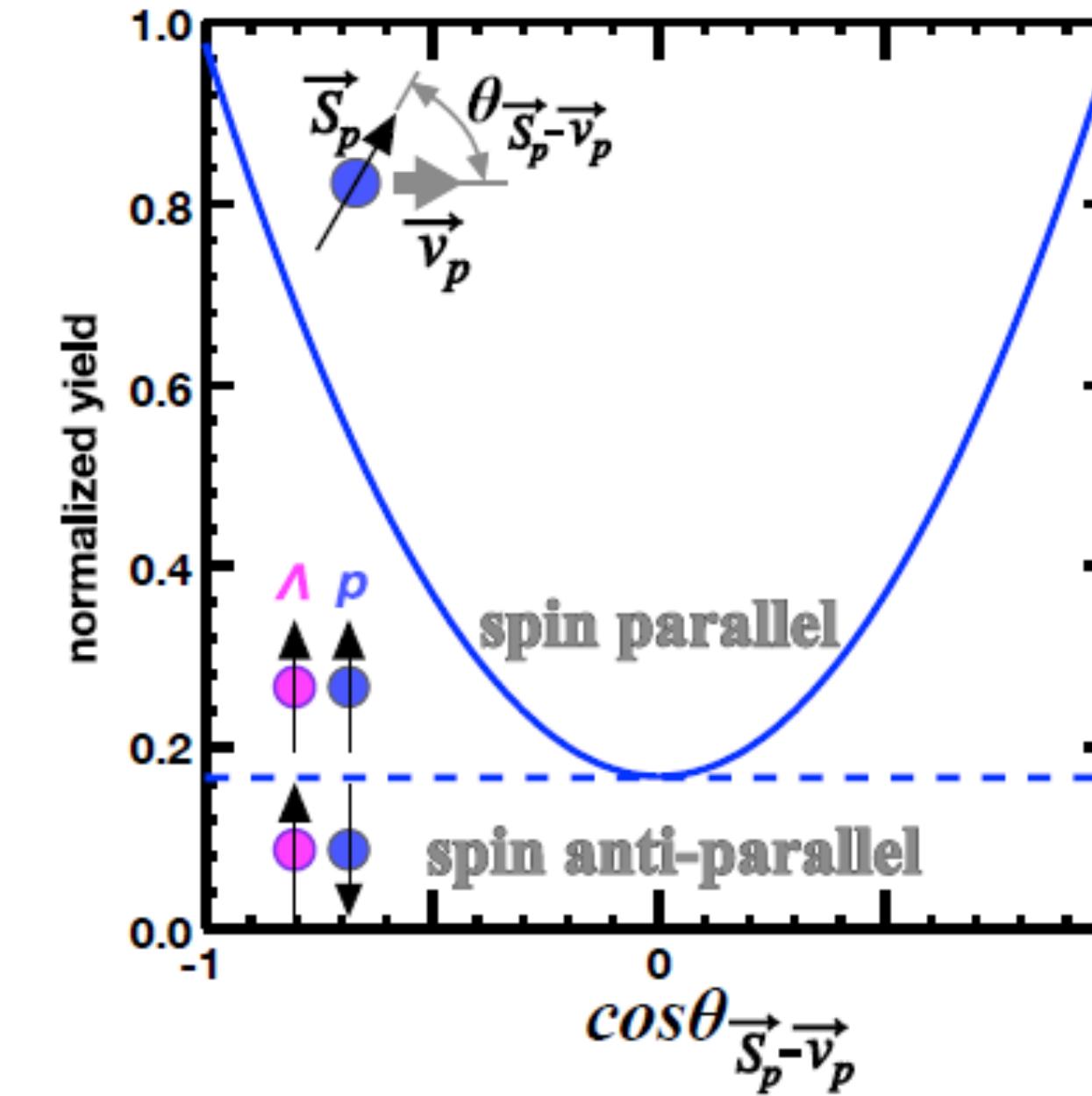
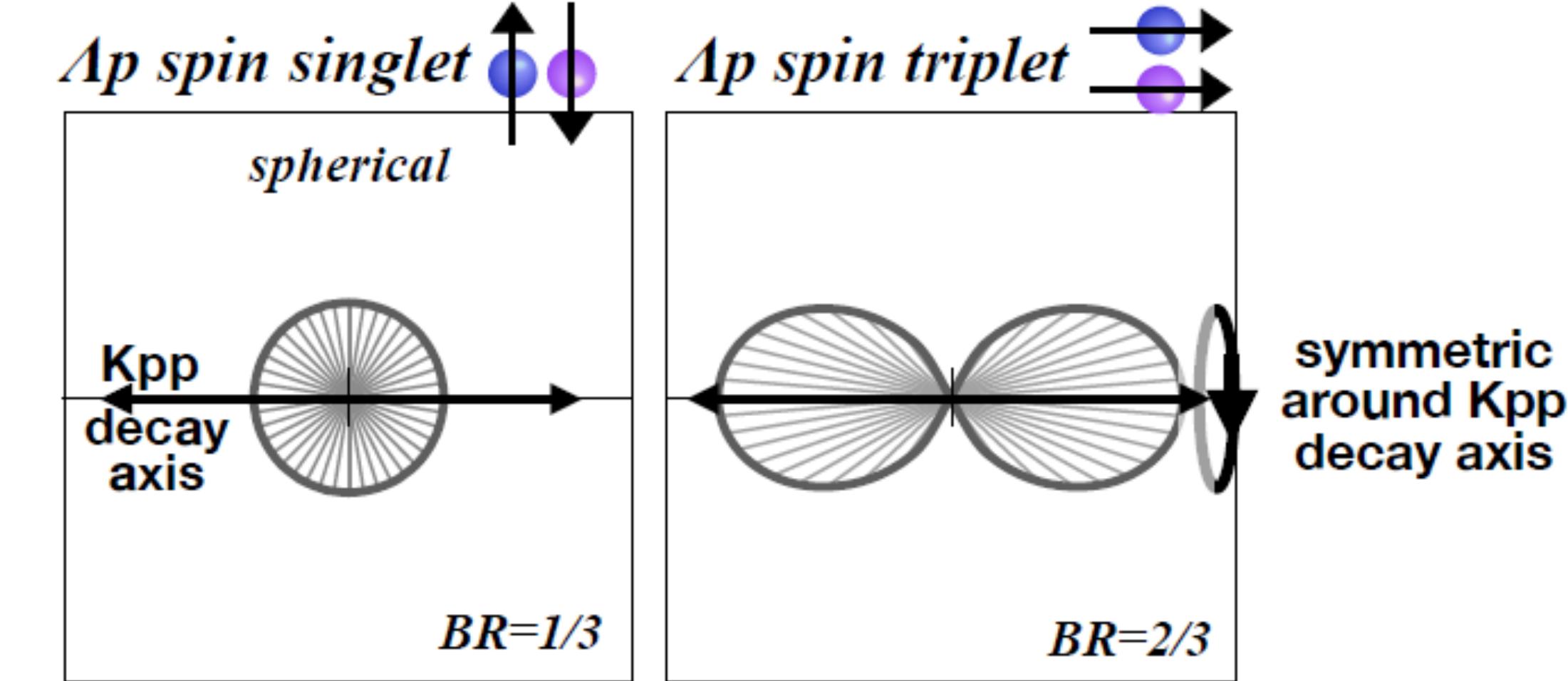
Λp decay axis and spin axis of $\bar{K}NN$ JP

spin axis distribution referring to the decay axis

$\bar{K}NN : J^P = 0^-, I = 1/2: I_{NN} = 1, S_{NN} = 0, L_{\bar{K}} = 0$



$\bar{K}NN : J^P = 1^-, I = 1/2: I_{NN} = 0, S_{NN} = 1, L_{\bar{K}} = 0$

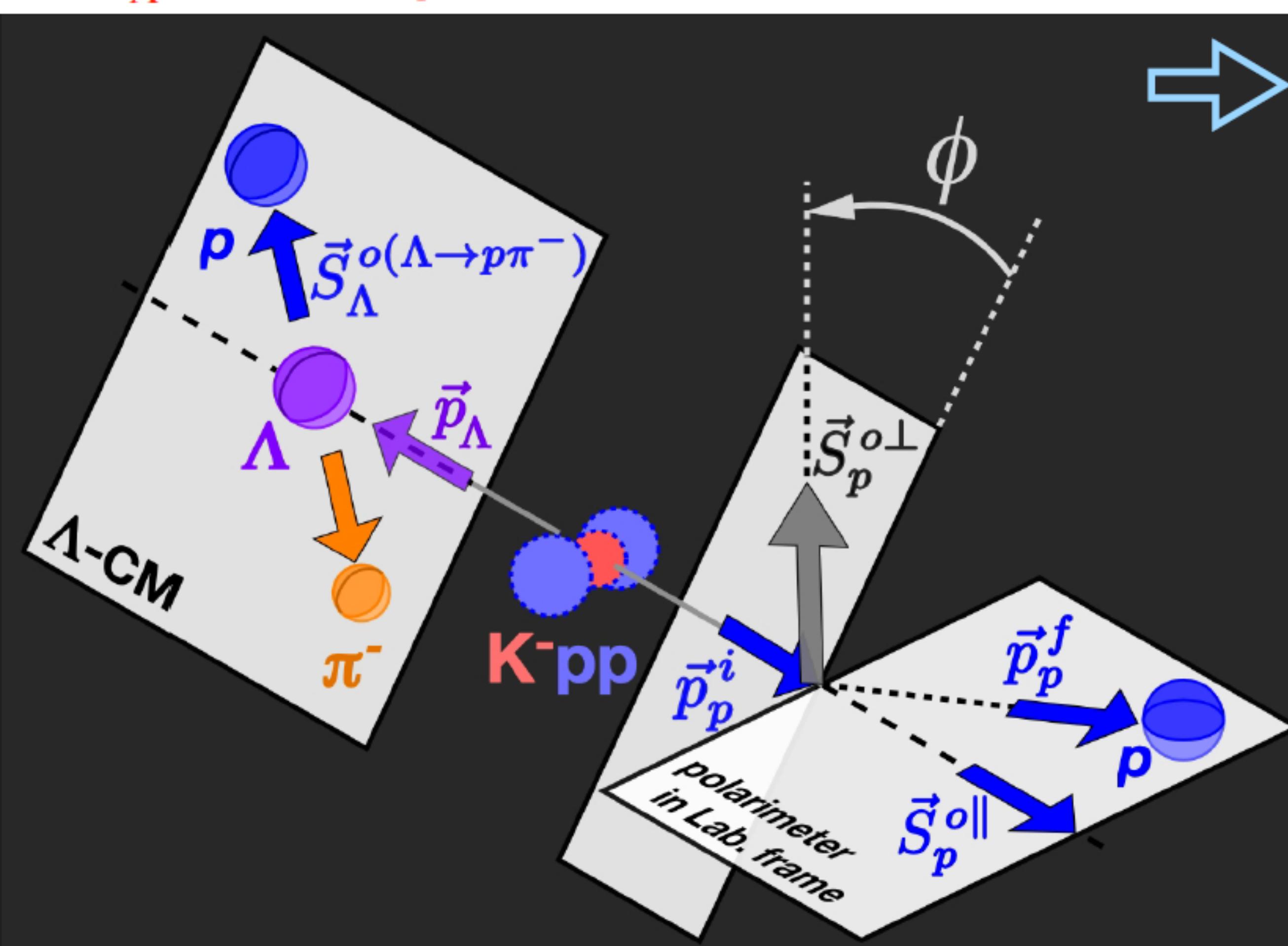


How to measure spin-spin correlation

- spin asymmetry measurement using $\Lambda \rightarrow p\pi^-$ & p-C(H) scattering -

p-C(H) scattering sensitive only on ϕ asymmetry

$$\vec{S}_\Lambda^{o(\Lambda \rightarrow p\pi^-)} \approx \vec{v}_p^{(\Lambda \rightarrow p\pi^-)} (\text{in } \Lambda\text{-CM})$$



$$N(\phi) d\phi \propto (1 + r \cdot \alpha_{\Lambda p} \cos \phi) d\phi$$

r : scaling factor

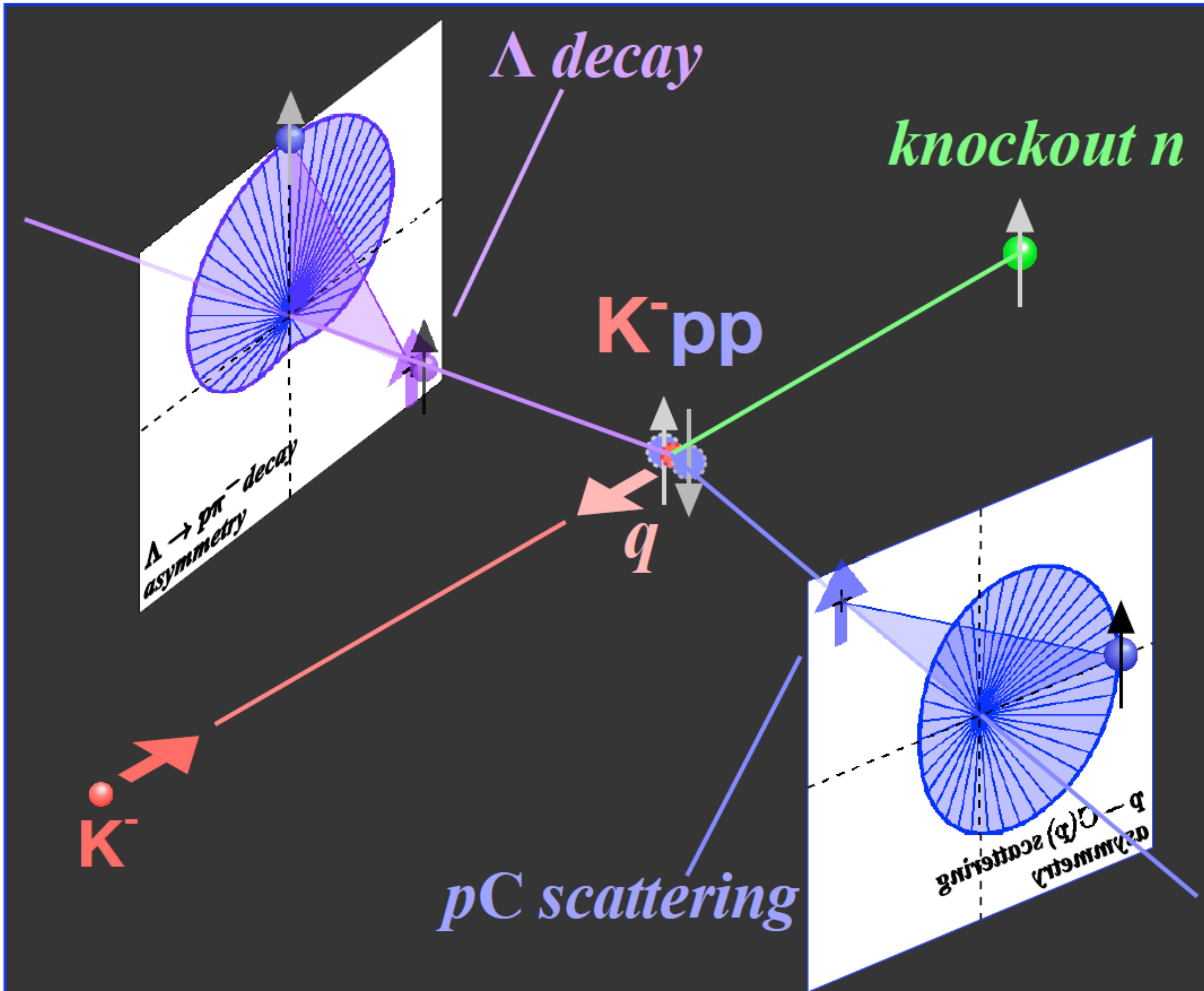
$$r = A_\Lambda \cdot A_{pC} \cdot \vec{S} \cdot \vec{S}^{\parallel} \cdot c_{conv}$$

A_Λ : Λ asymmetry parameter

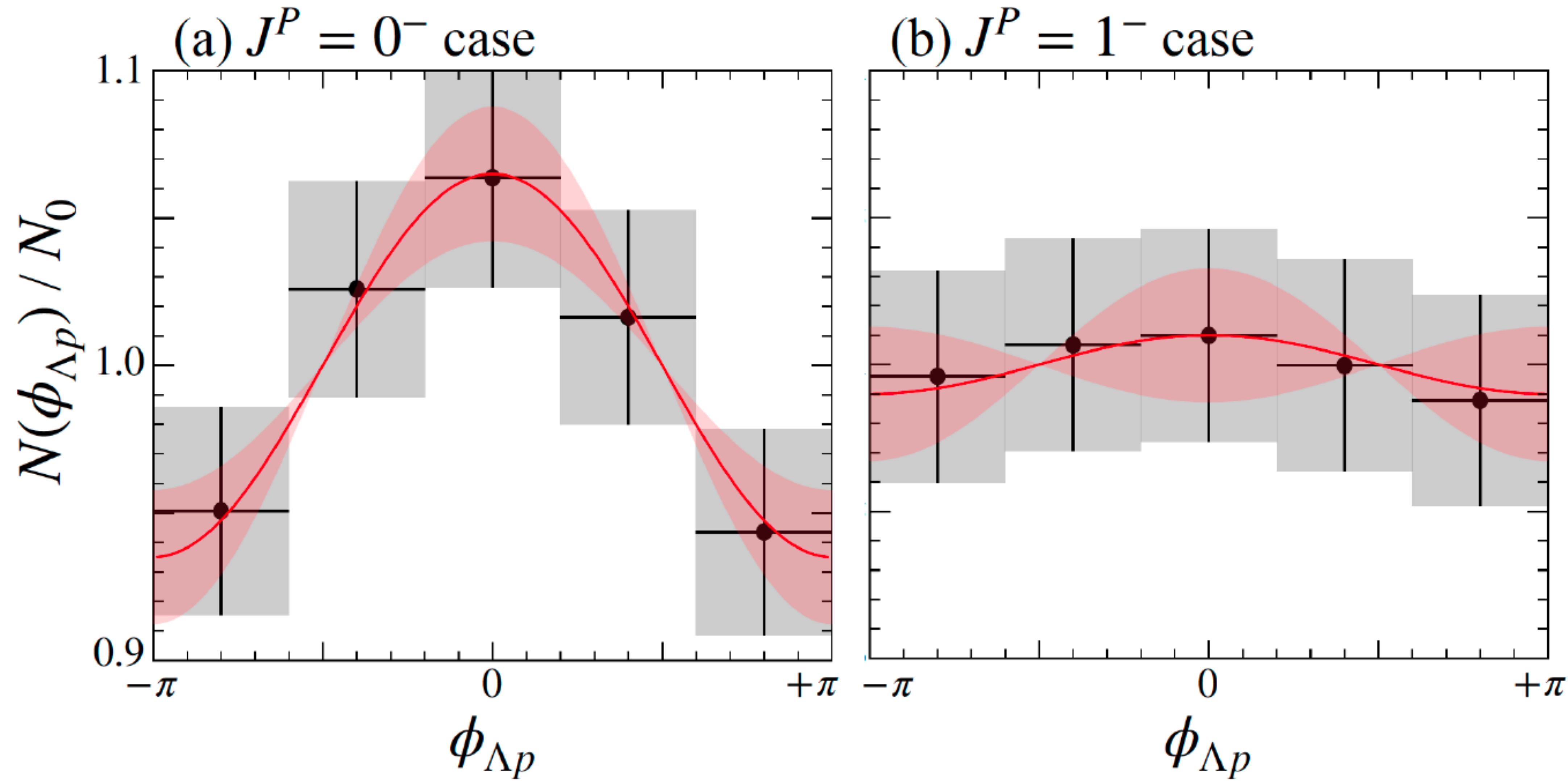
A_{pC} : proton spin-analyzing-power
on carbon (and on p)

$\vec{S} \cdot \vec{S}^{\parallel}$ ($\equiv \vec{S}_p \cdot \vec{S}_p^{\parallel}$) : spin sensitivity
referring to motional axis

c_{conv} : convolution coefficient
between two asymmetries



Λp spin-spin asymmetry



... Construction is in progress, led by F. Sakuma

New Spectrometer under construction

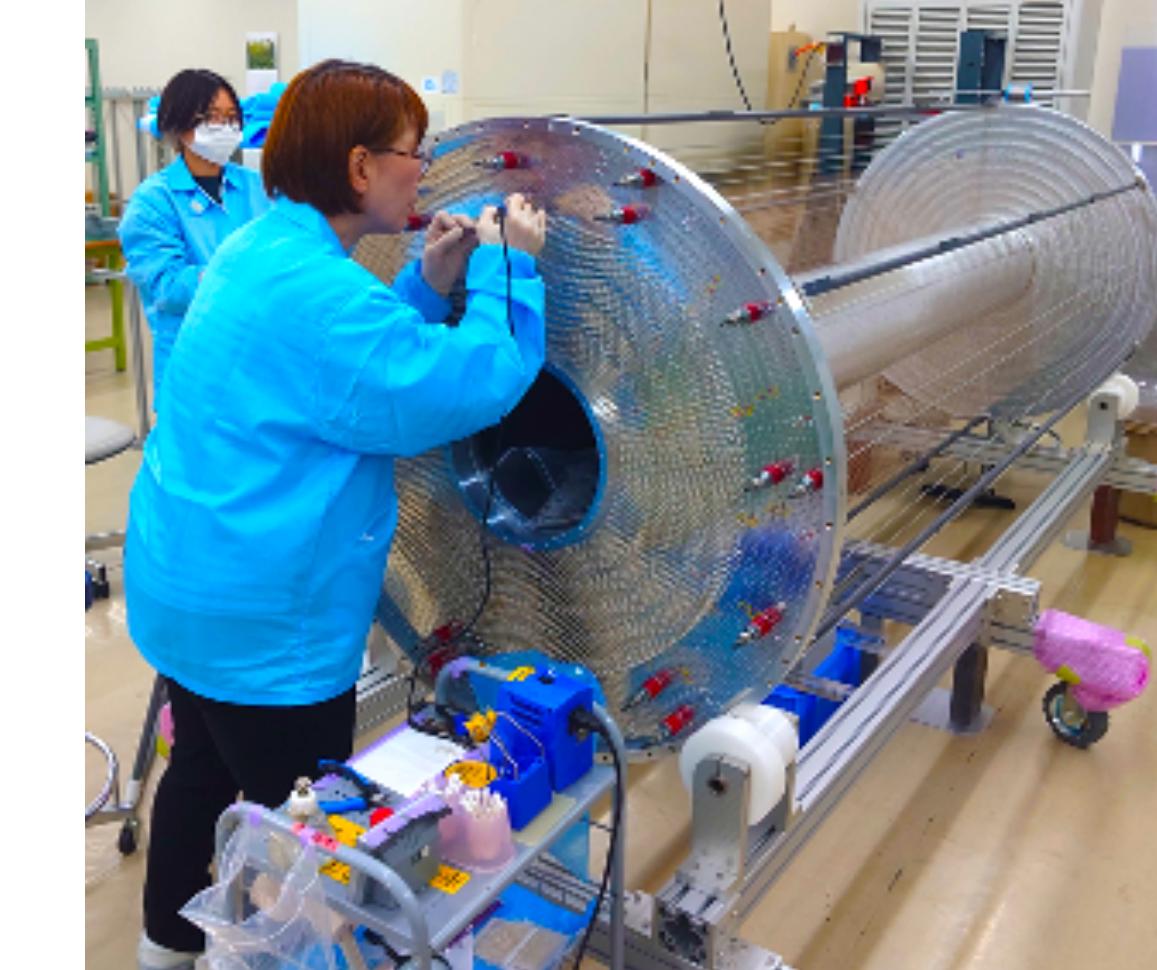
Preparation of planned devices / detectors



Return Yoke



Super-conducting Solenoid Coil



Cylindrical Drift Chamber (CDC)

under construction

New Spectrometer under construction

Preparation of planned devices / detectors



Return Yoke



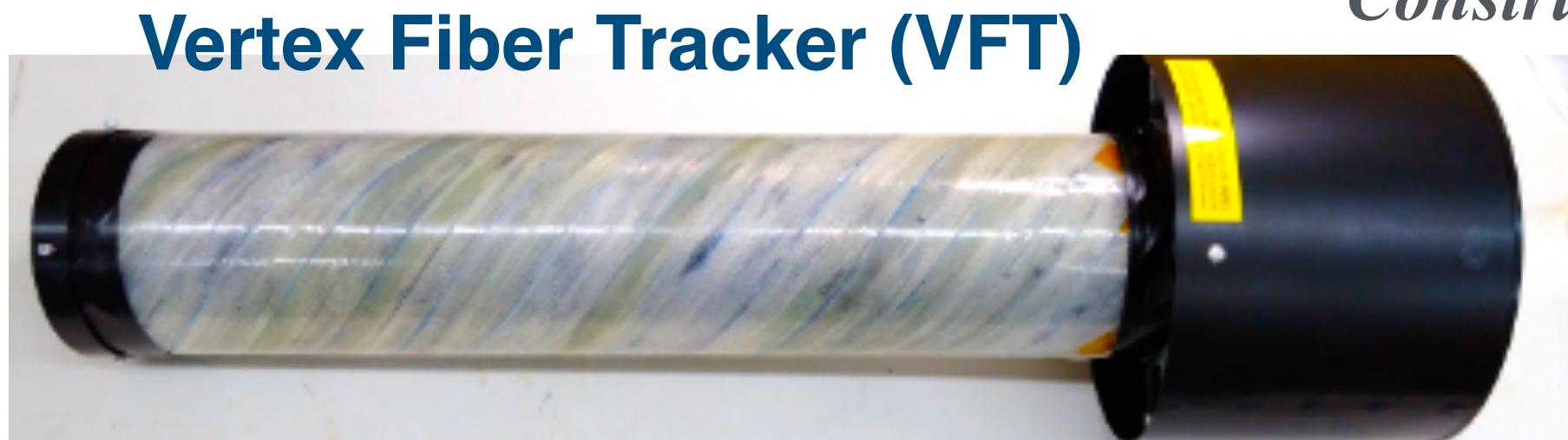
Super-conducting Solenoid Coil



Cylindrical Drift Chamber (CDC)

Additional detectors to improve

To detect the proton in Fermi-motion
& to drastically improve vertex resolution (Λ/Σ^0 separation)

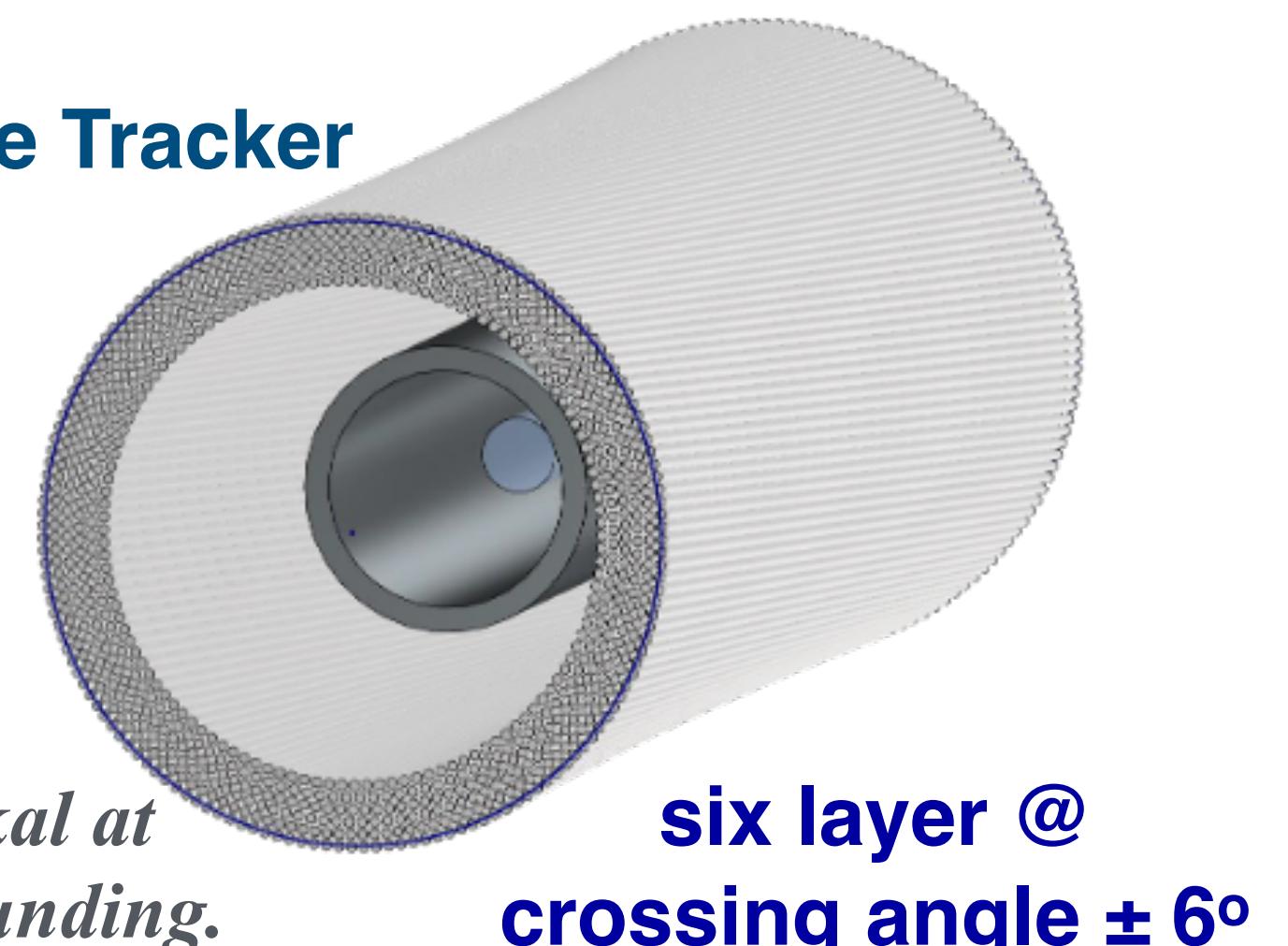


Vertex Fiber Tracker (VFT)

Constructed by T.Hashimoto

double layer
@ crossing
angle $\pm 45^\circ$

Currently, J. Zmeskal at
SMI is applying for funding.



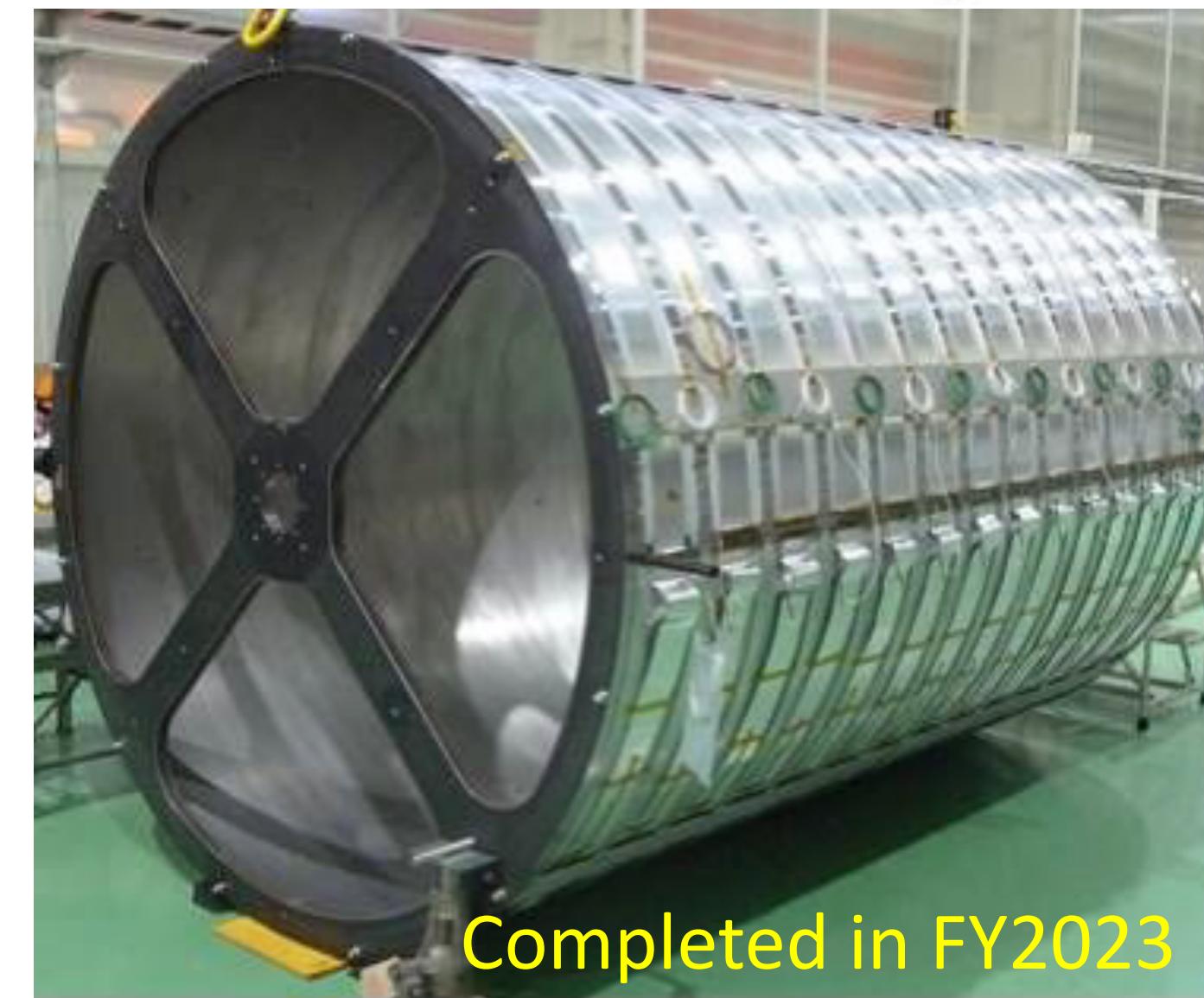
six layer @
crossing angle $\pm 6^\circ$

Superconducting Solenoid Magnet

- Same design as “the detector solenoid magnet” for COMET-I

**being constructed in cooperation with
the J-PARC Cryogenics Section**

- 3.3m x 3.3m x 3.9m, ~108t in total
- Max. field of 1.0T @ center
 - 189A – 10V
- NbTi/Cu SC wire, 98km in total
- **Conduction-cooling with GM*3**
- Semi-active quench-back system
- **Will be completed in FY2024**

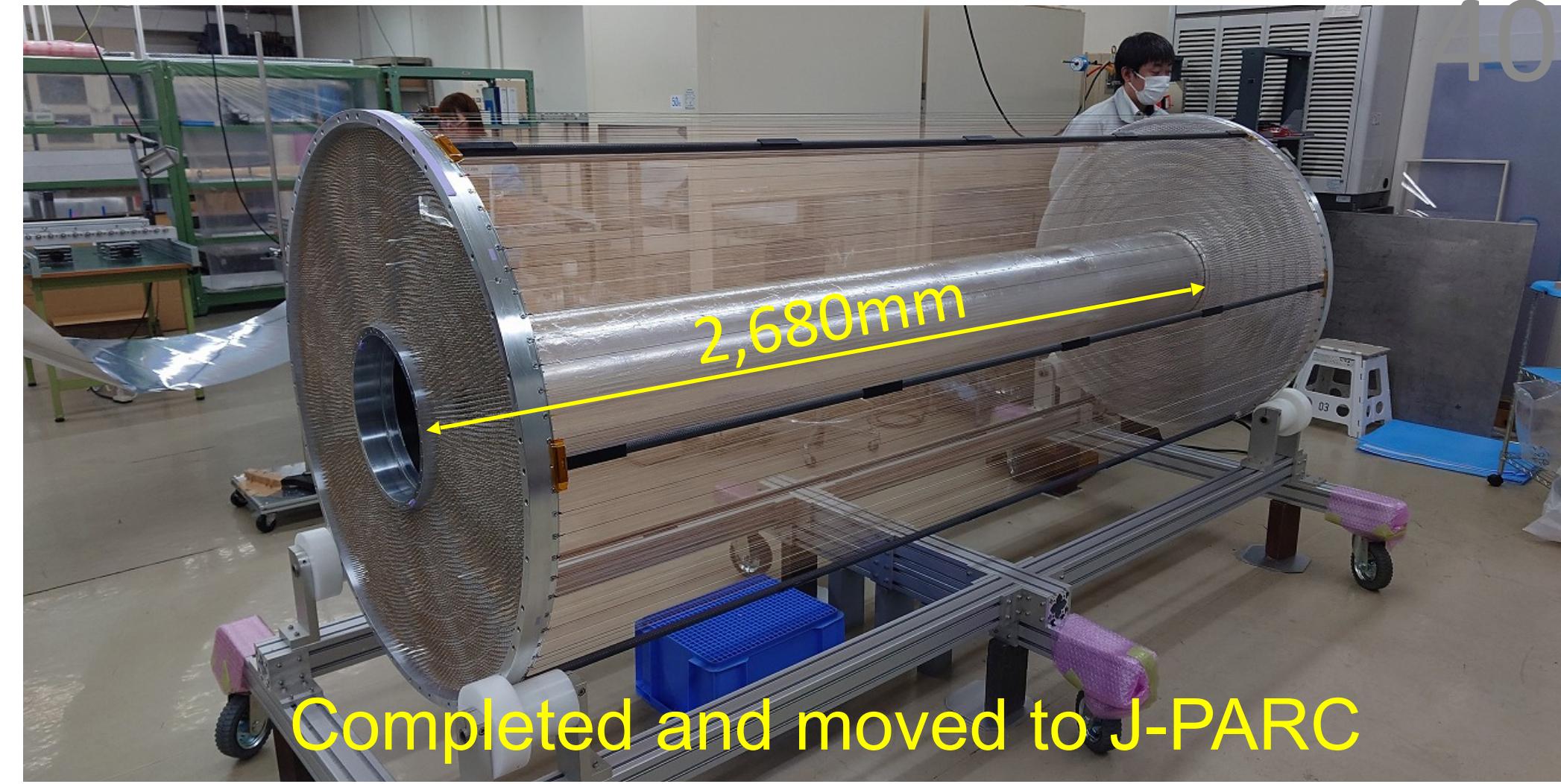


SHI FA-50
(air cooling)

RDE-418D4

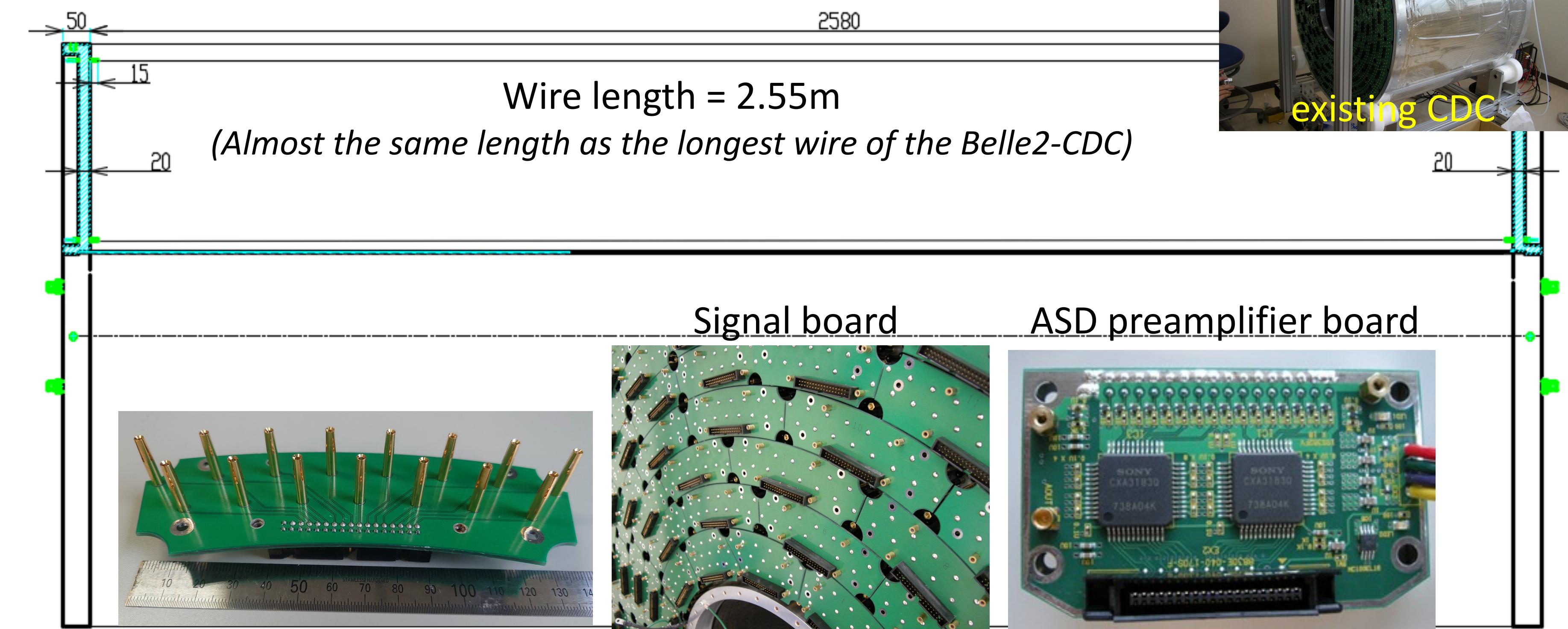
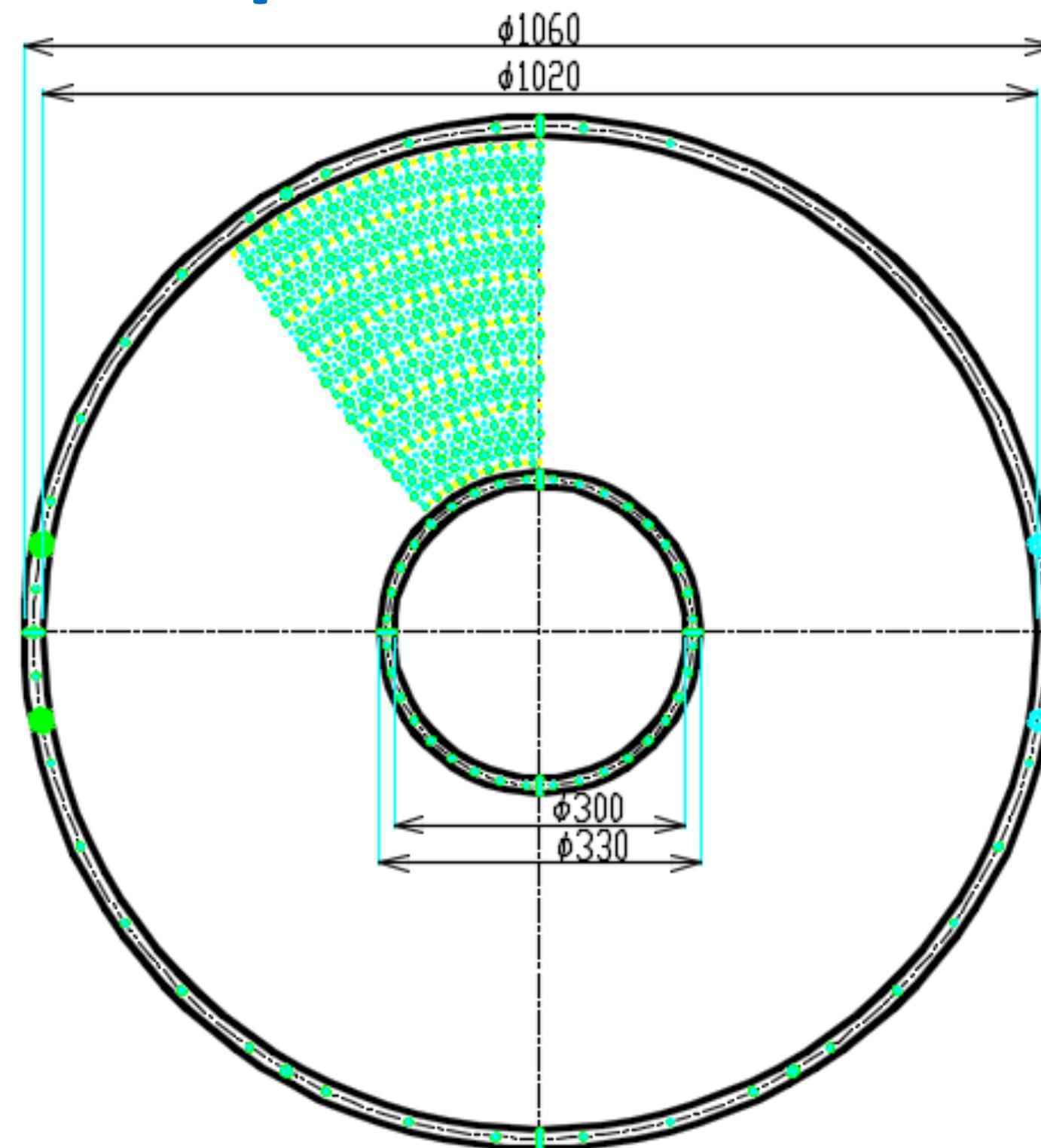
Cylindrical Drift Chamber (CDC)

- 3 times the length of the existing CDC
 - Gas: Ar/CO₂=90/10
- The same design of the present end-cap
- Readout systems are reused



Completed and moved to J-PARC

Completed this month, and commissioning will soon start @ J-PARC



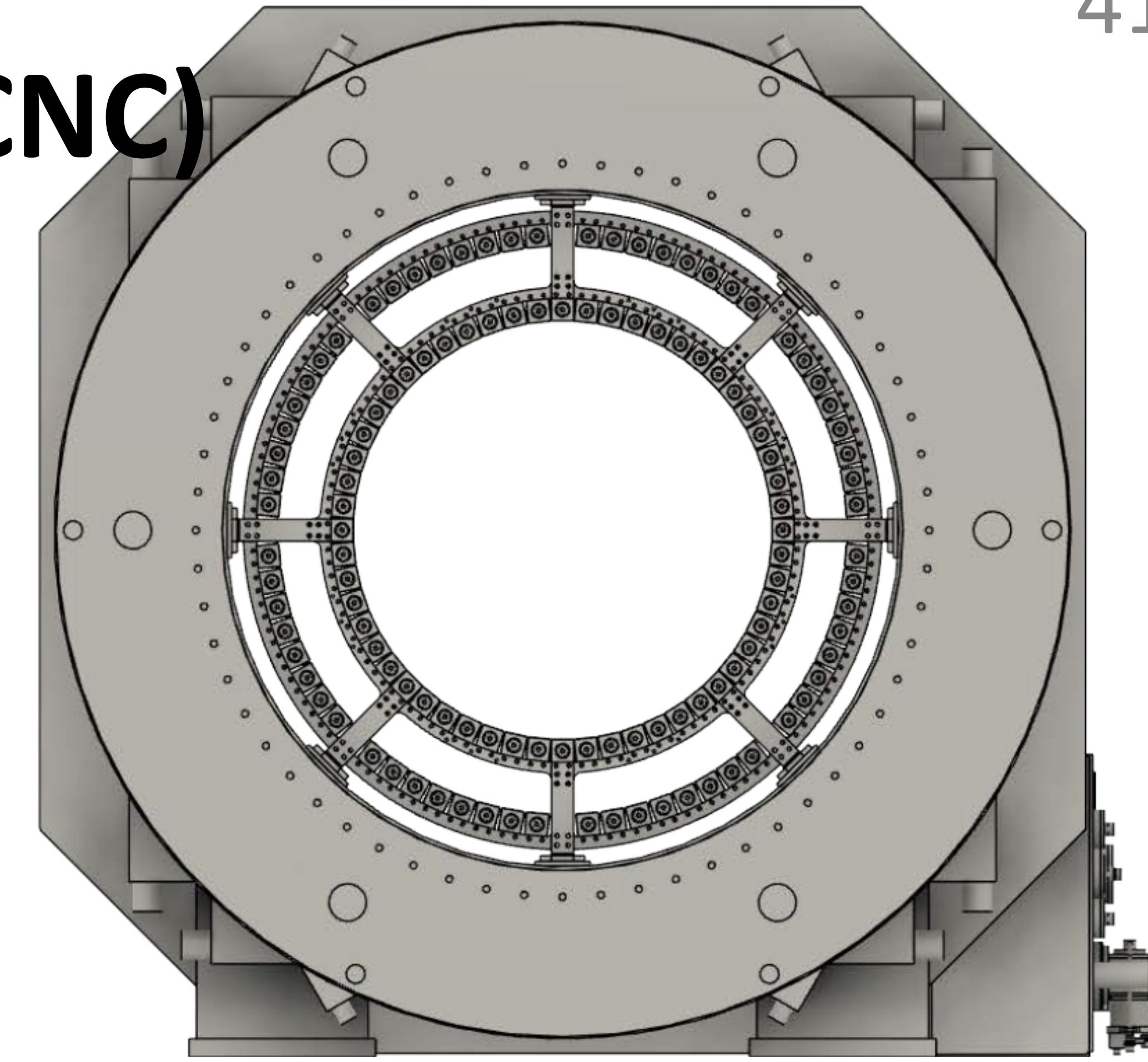
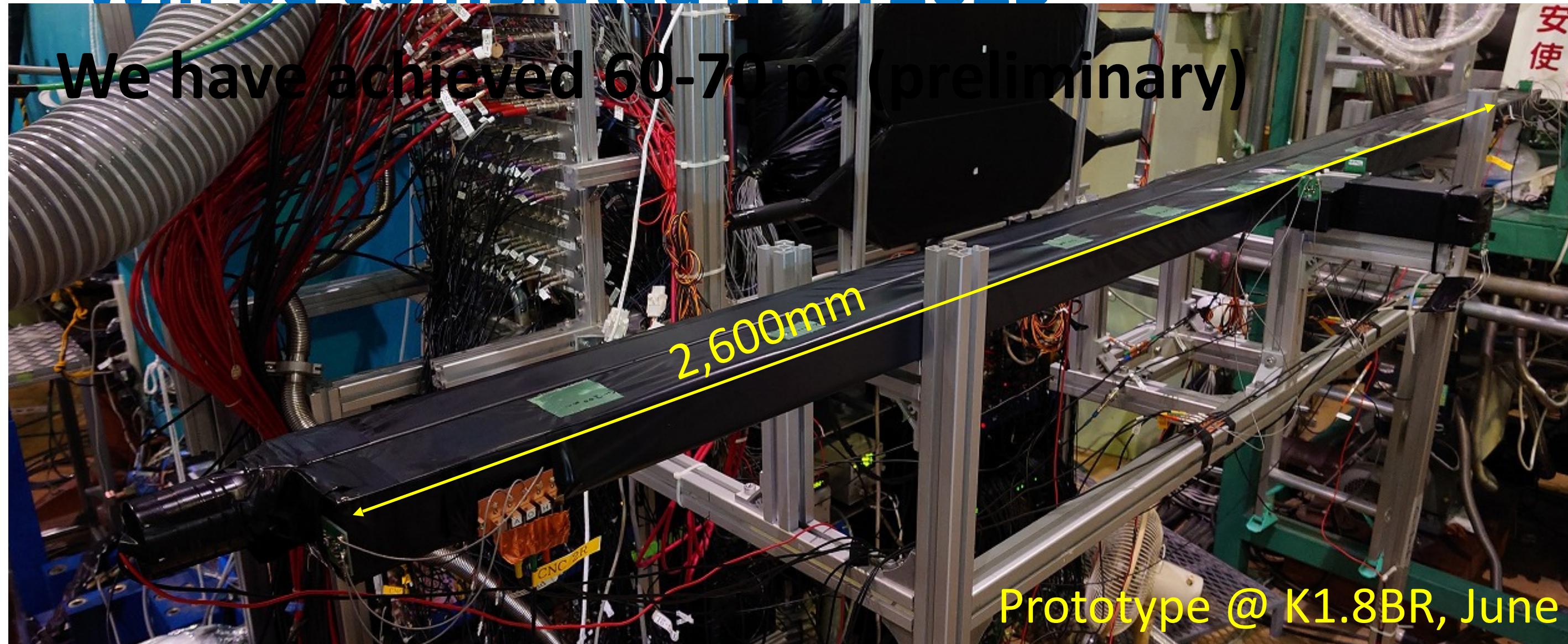
Signal board

ASD preamplifier board



Cylindrical Neutron Counter (CNC)

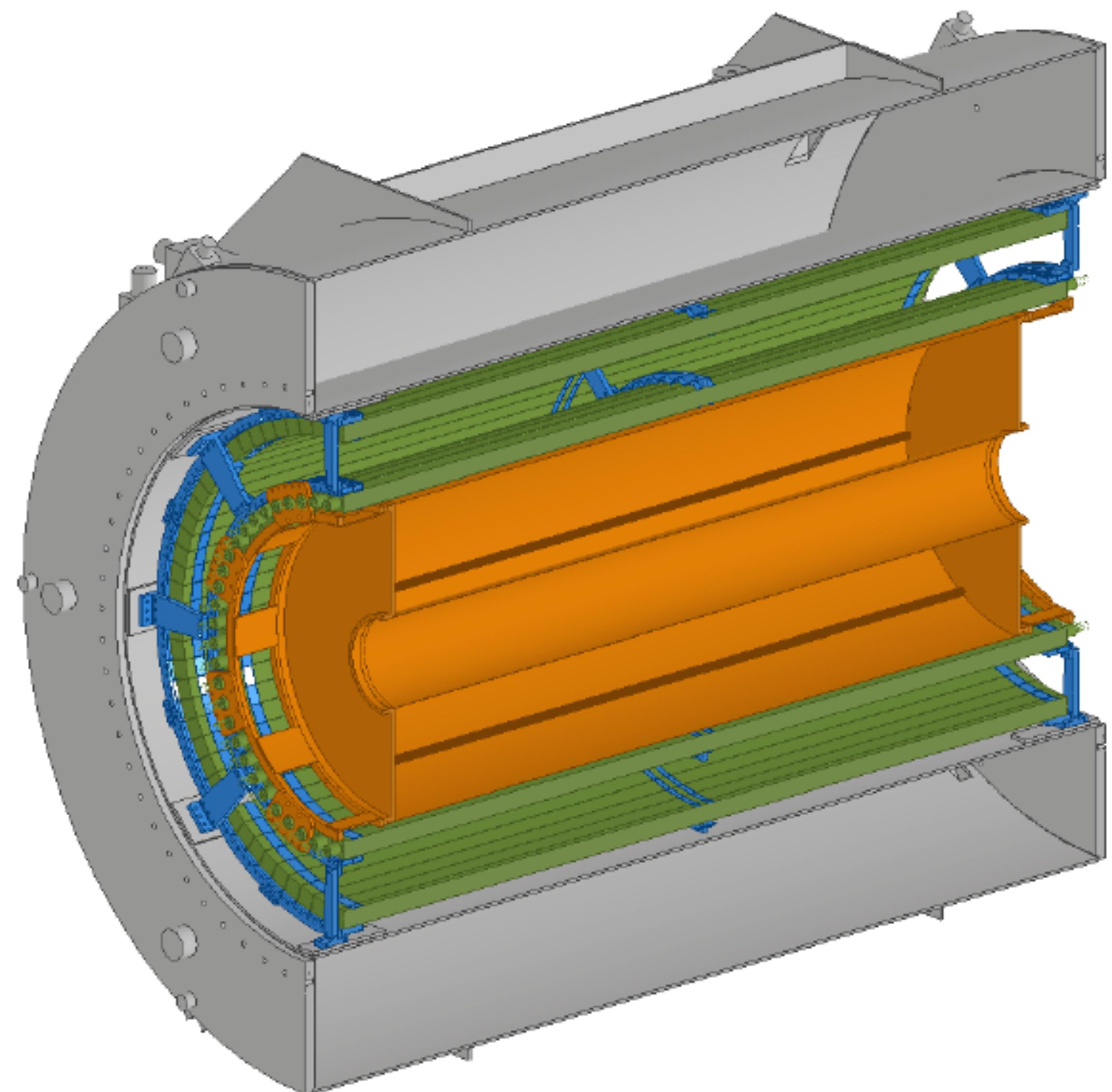
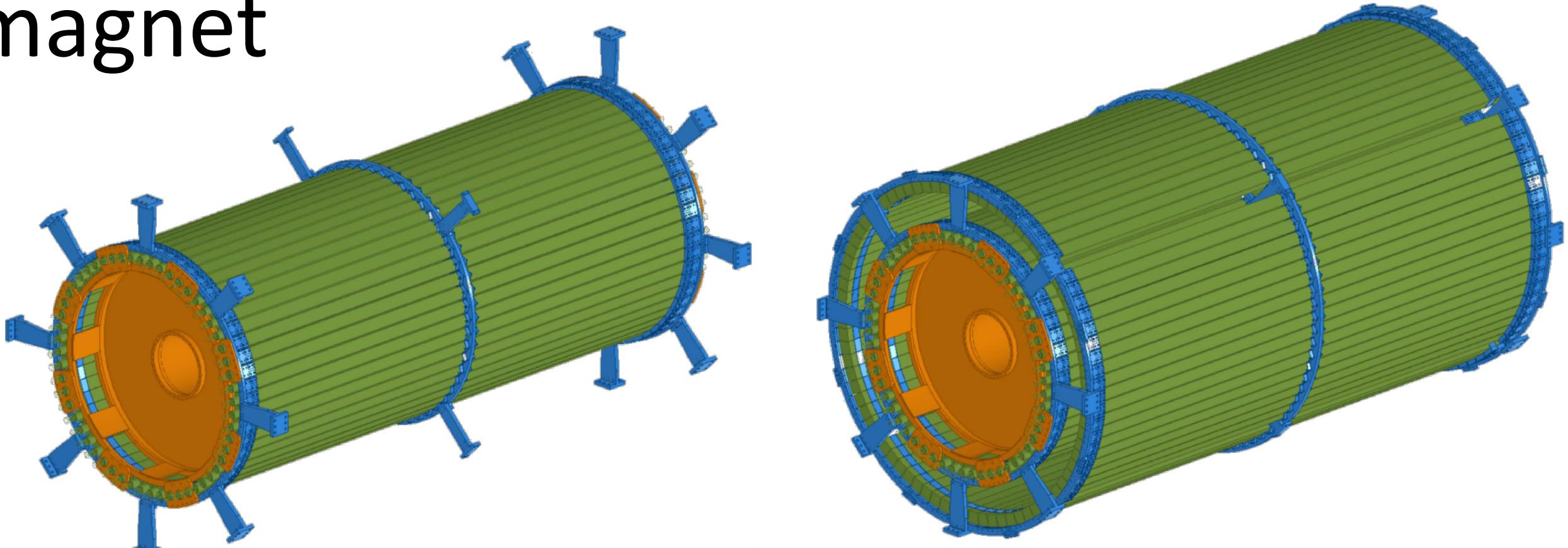
- scintillator array: 2 layers, 12cm thickness
- Neutron detection efficiency of 12~36%
- $56+80=136$ modules
 - ELJEN EJ-200: (T)60mm, (W)60mm, (L)3,000mm
- 1.5-inch FM-PMT [H8409(R7761)] & MPPC array [S13361-6050AE-04]
- **Will be completed in FY2025**



- 136 scintillators in total
- 56 segments @ $r548\sim608\text{mm}$
 - 112 FM-PMTs
 - 80 segments @ $r780\sim840\text{mm}$
 - 160 MPPC-arrays

Support Structure

- CNC is supported at upstream, downstream and middle position
 1. pillars are mounted on the inner cylinder of the magnet
 2. ring structures are installed on the pillars
 3. each module is mounted on the ring structures
- CDC is installed by inserting a long frame bar into the center of the CDC and magnet



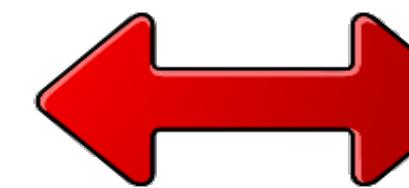
Will be prepared in FY2025-26

Summary of Experimental Status

Negative results

AMADEUS@DAΦNE

$^{12}\text{C}(\text{K}^-, \Lambda\text{p})$ EPJC79(2019)190



HADES@GSI

$\text{p} + \text{p} \rightarrow (\Lambda + \text{p}) + \text{K}^+$ @ 3.5GeV

PLB742(2015)242



LEPS@SPring-8

$d(\gamma, \pi^-\text{K}^+)\text{X}$ @ 1.5-2.4 GeV

PLB728(2014)616



Positive results

FINUDA@DAΦNE

$^{6}\text{Li}/^{7}\text{Li}/^{12}\text{C}(\text{K}^-, \Lambda\text{p})$ PRL94(2005)212303

Multi-NA processes?

DISTO@SATURNE

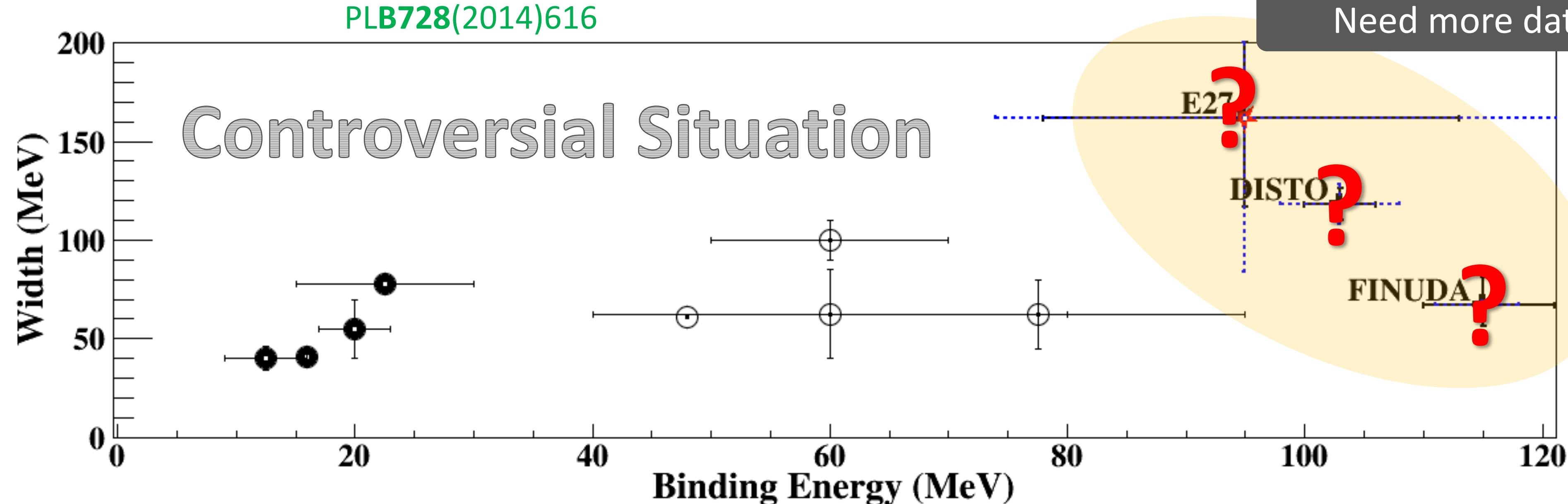
$\text{p} + \text{p} \rightarrow (\Lambda + \text{p}) + \text{K}^+$ @ 2.85GeV

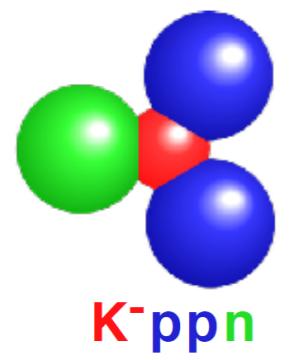
Intermediate $\text{N}^* \rightarrow \Lambda\text{K}^+?$

E27@J-PARC

$d(\pi^+, \text{K}^+)\Sigma^0\text{p}$ @ 1.69 GeV/c

Need more data

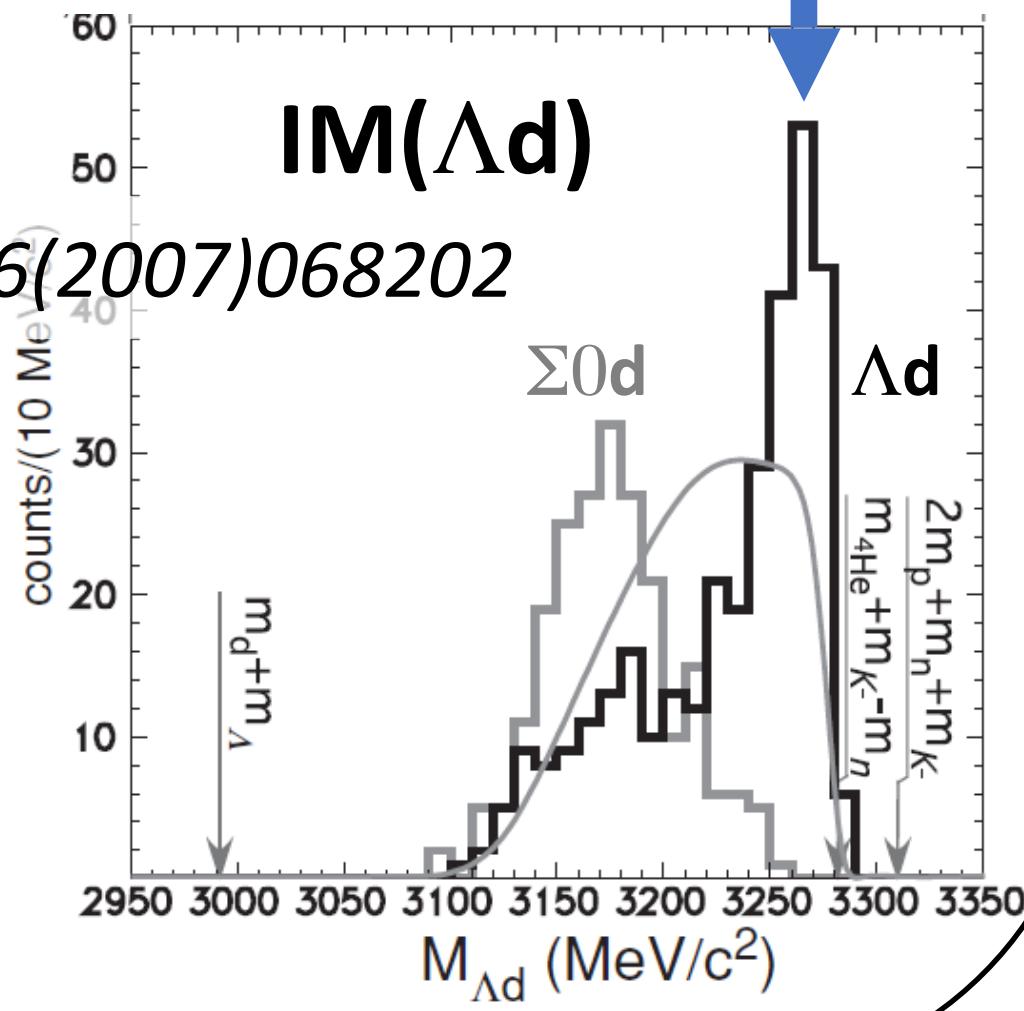
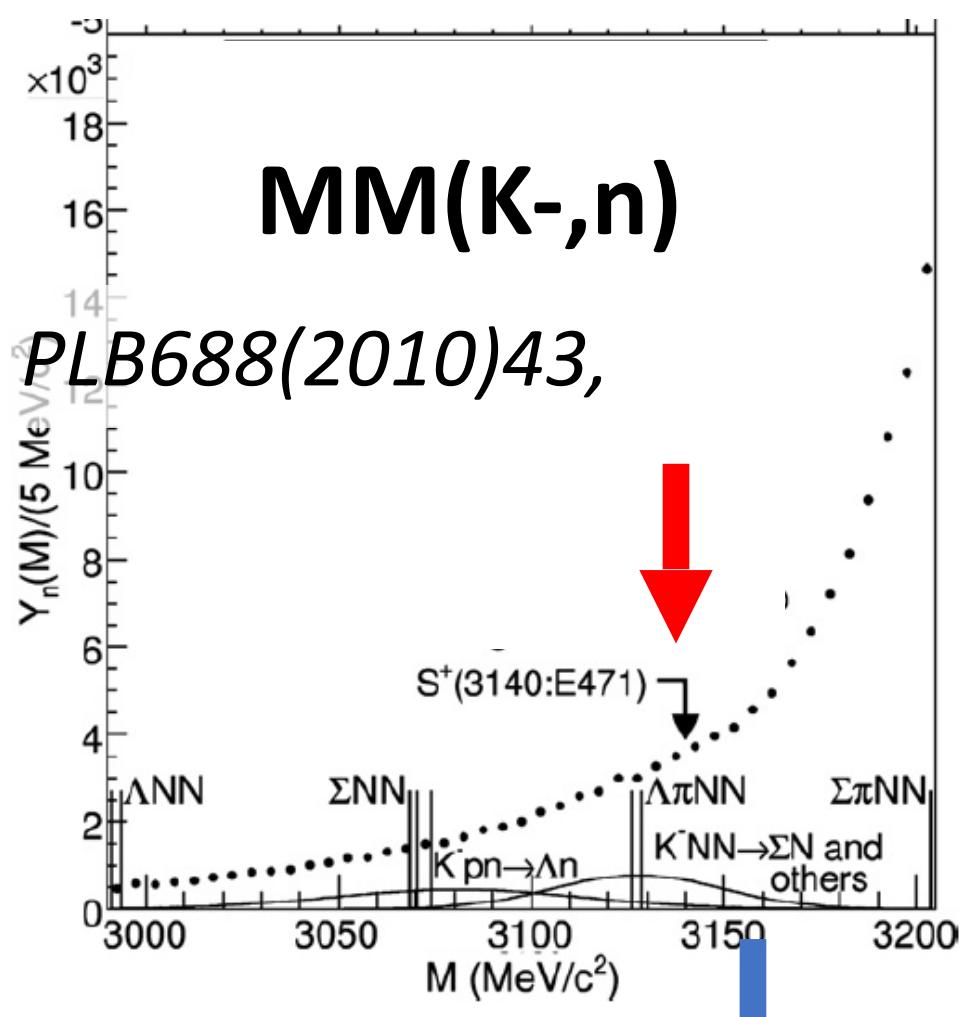
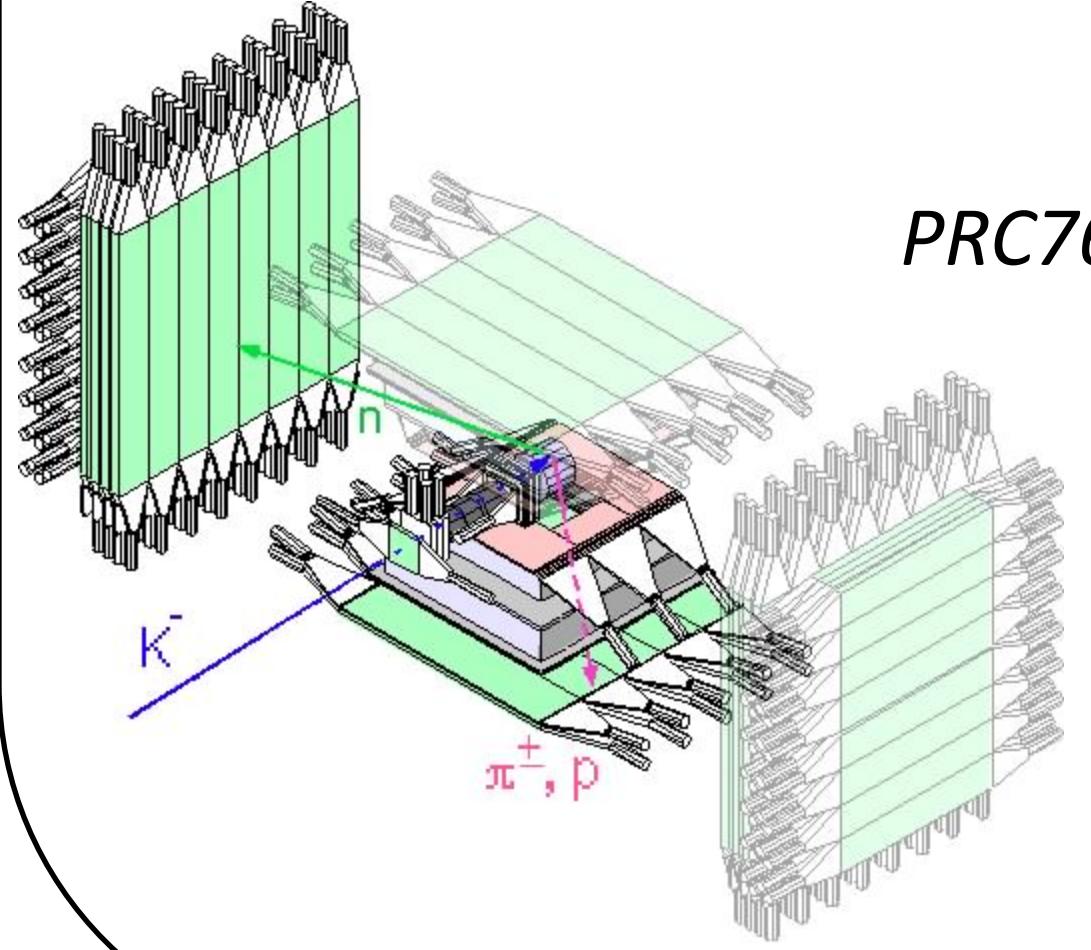
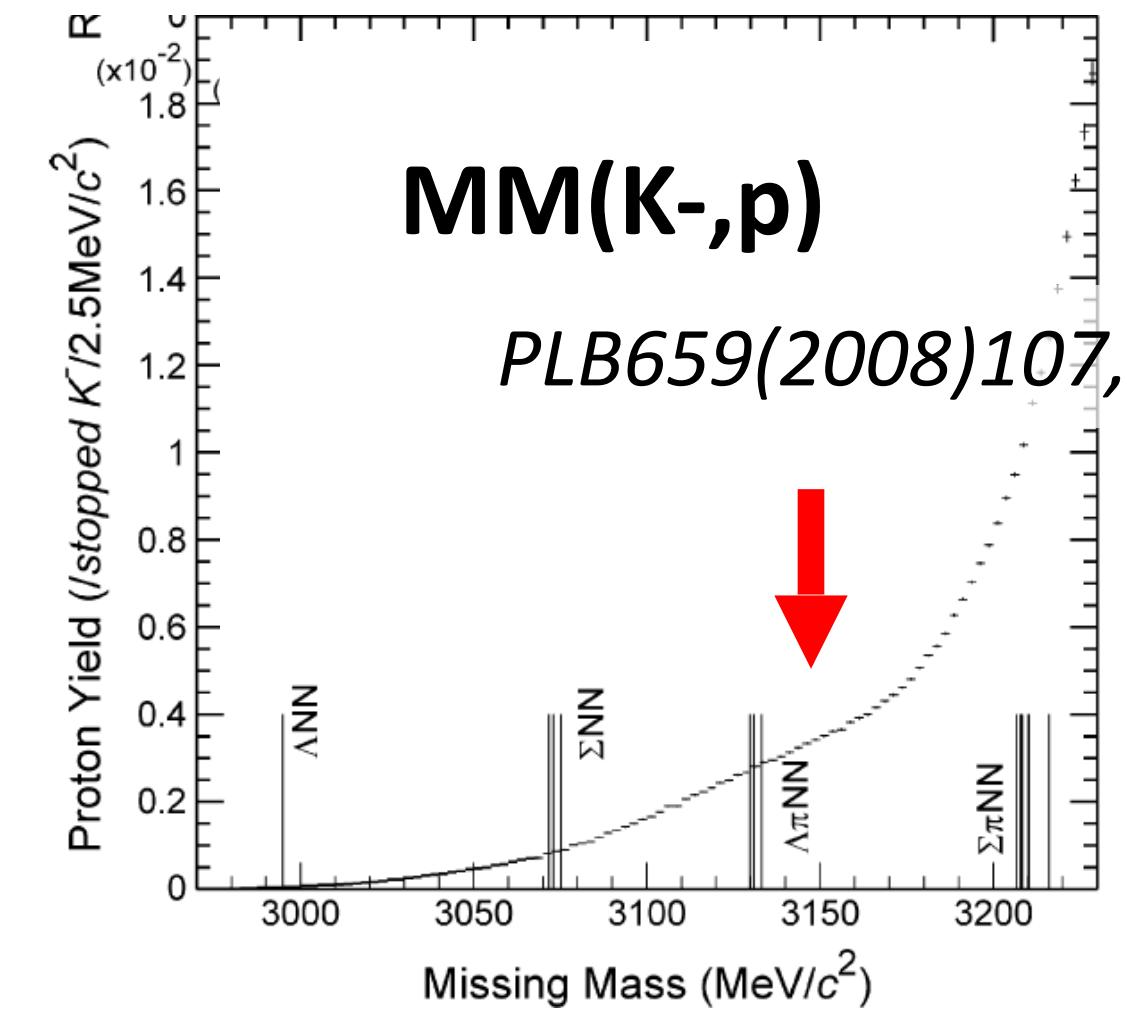




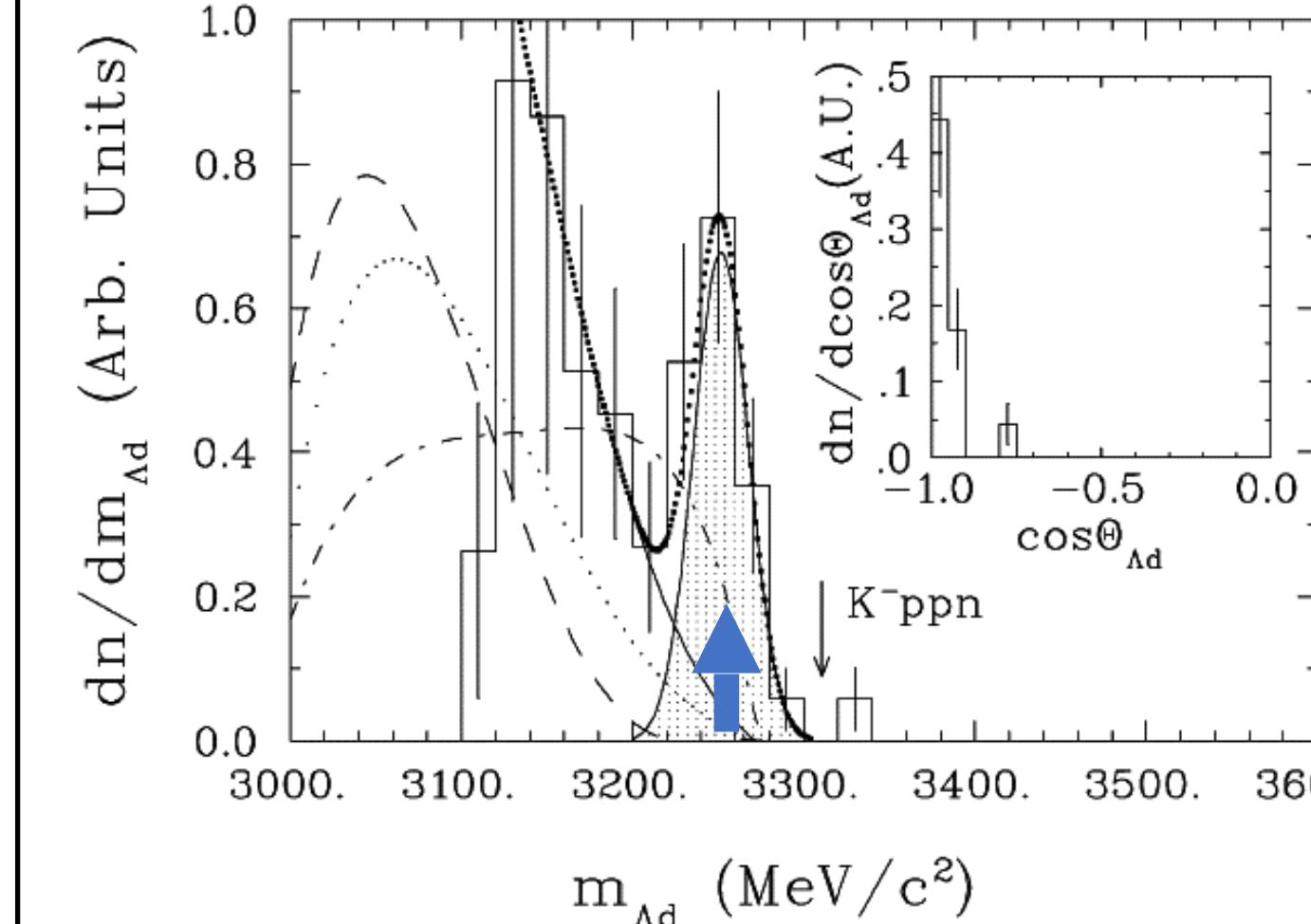
$\bar{K}NNN$ Searches so far

E471/E549@KEK

$^4\text{He}(\text{stopped-K}, \text{p}/\text{n}/\Lambda\text{d})$

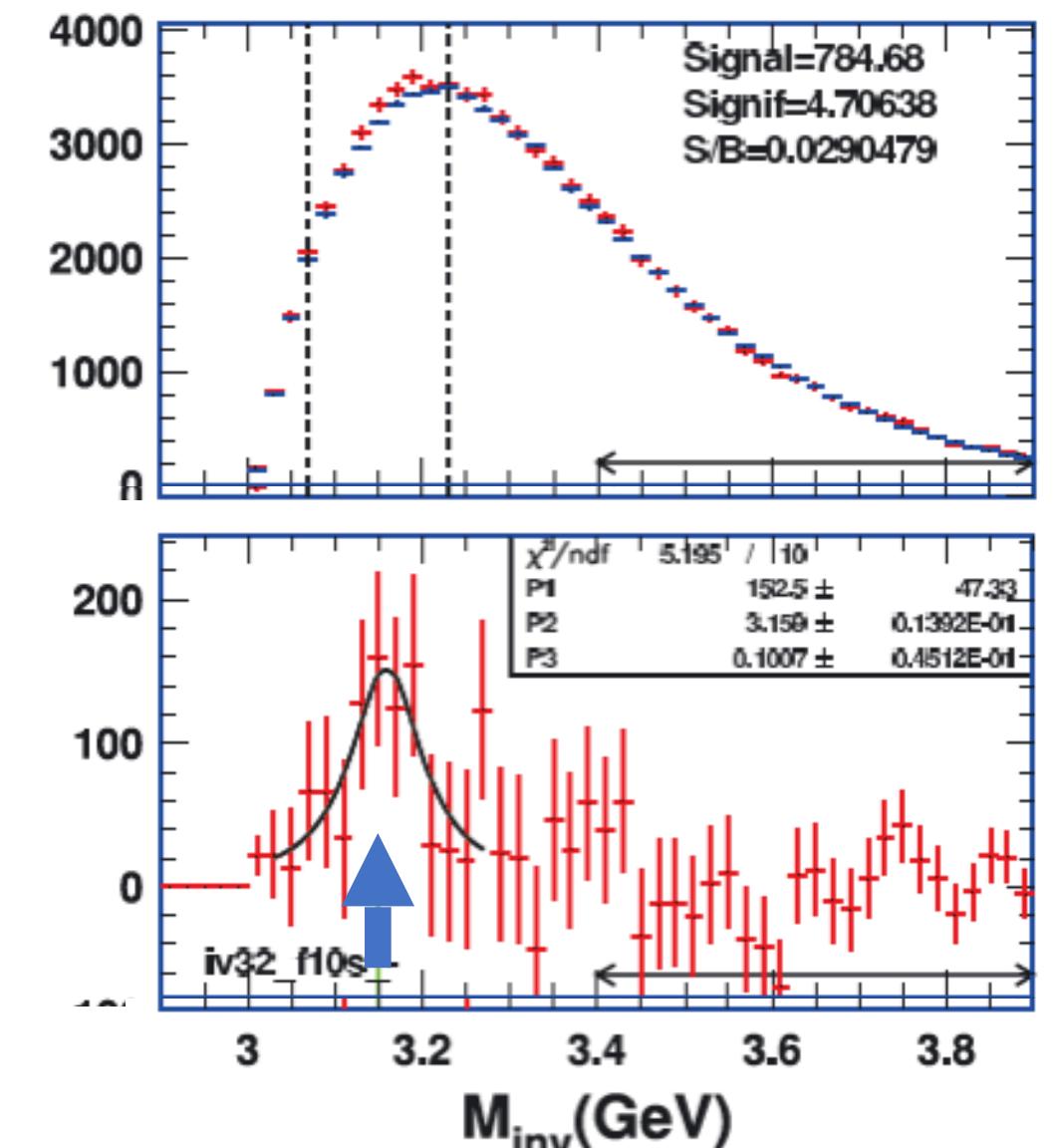


FINUDA@DAΦNE
Li/C(stopped-K, Λd)



PLB654(2007)80

FOPI@GSI
 Λd in $\text{Ni}+\text{Ni}$



EXA05 Conference (2005)

No conclusive results.

multi-N absorption in stopped-K reaction makes interpretation difficult

The detail can be found:

— in a review —

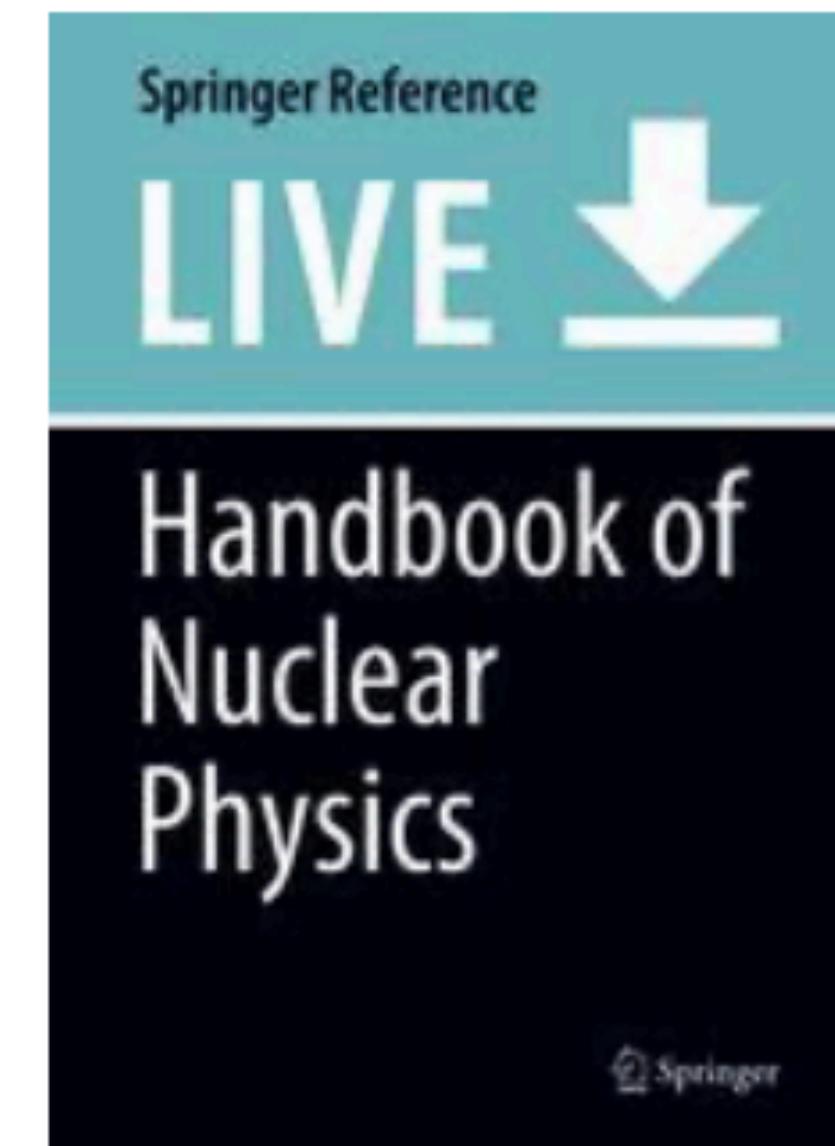
$\bar{K}N$ interaction study via kaonic atom

Search for $\bar{K}NN$ nuclear bound state as a natural extension of $\Lambda(1405) \equiv \bar{K}N$

Recent results on \bar{K} bound state

This is a fee-based literature.

Kaonic Nuclei from the Experimental Viewpoint



Iwasaki, M. (2022).

Kaonic Nuclei from the Experimental Viewpoint.

In: Tanihata, I., Toki, H., Kajino, T. (eds)

Handbook of Nuclear Physics . Springer, Singapore.

https://doi.org/10.1007/978-981-15-8818-1_37-1

https://link.springer.com/referenceworkentry/10.1007/978-981-15-8818-1_37-1

Proposed K1.8BR Upgrade

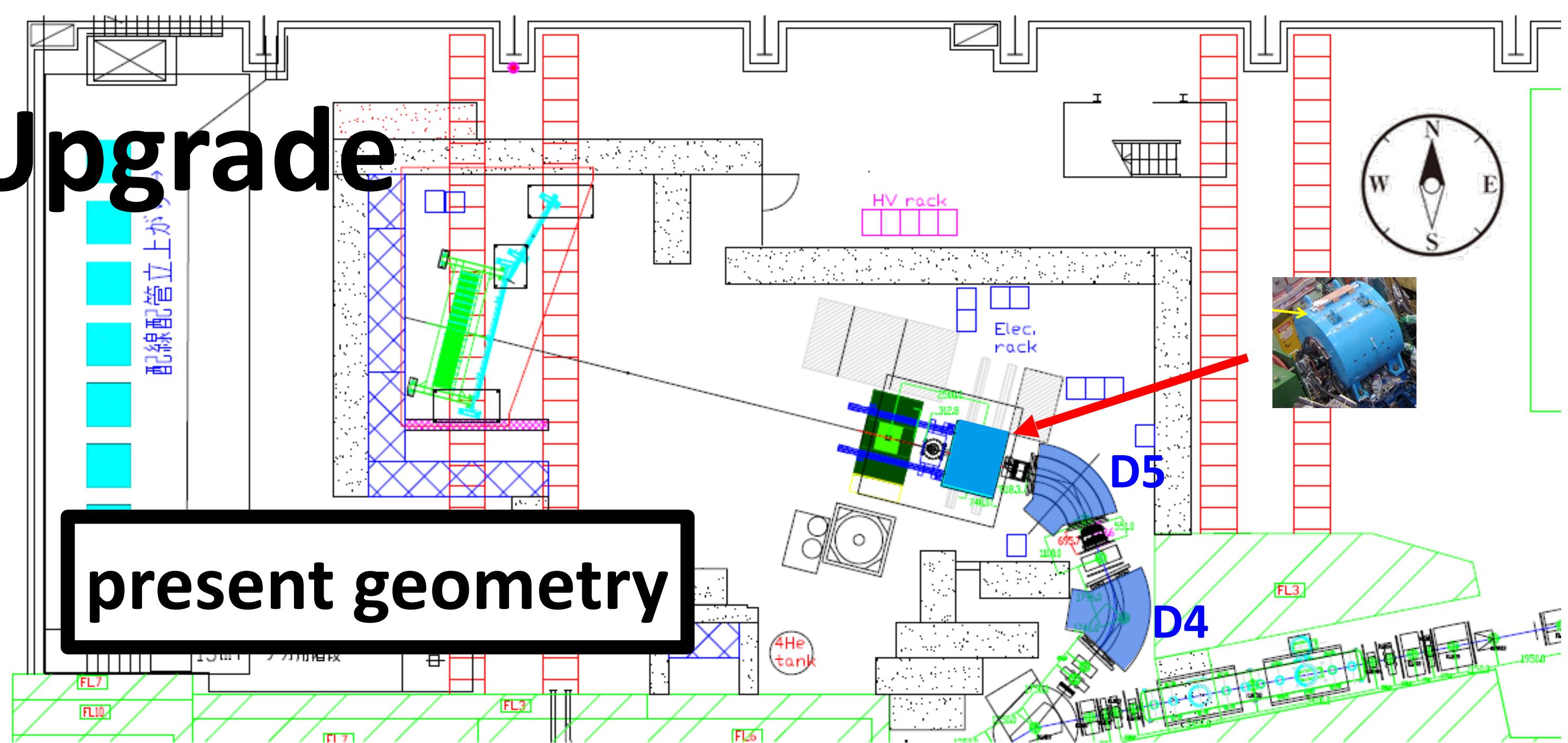
- Shortened beam line to enhance Kaon yield

Shorten the beam line ($\sim 2.5\text{m}$) by removing the final D5 magnet

with π/K ratio ~ 2

Relative beam-line length (m)	@ D5	@ D4
Present CDS	0	-3.7
New CDS	+1.2	-2.5

➤ K- yield increases by
~ 1.4 times @ 1.0 GeV/c



Proposed K1.8BR Upgrade

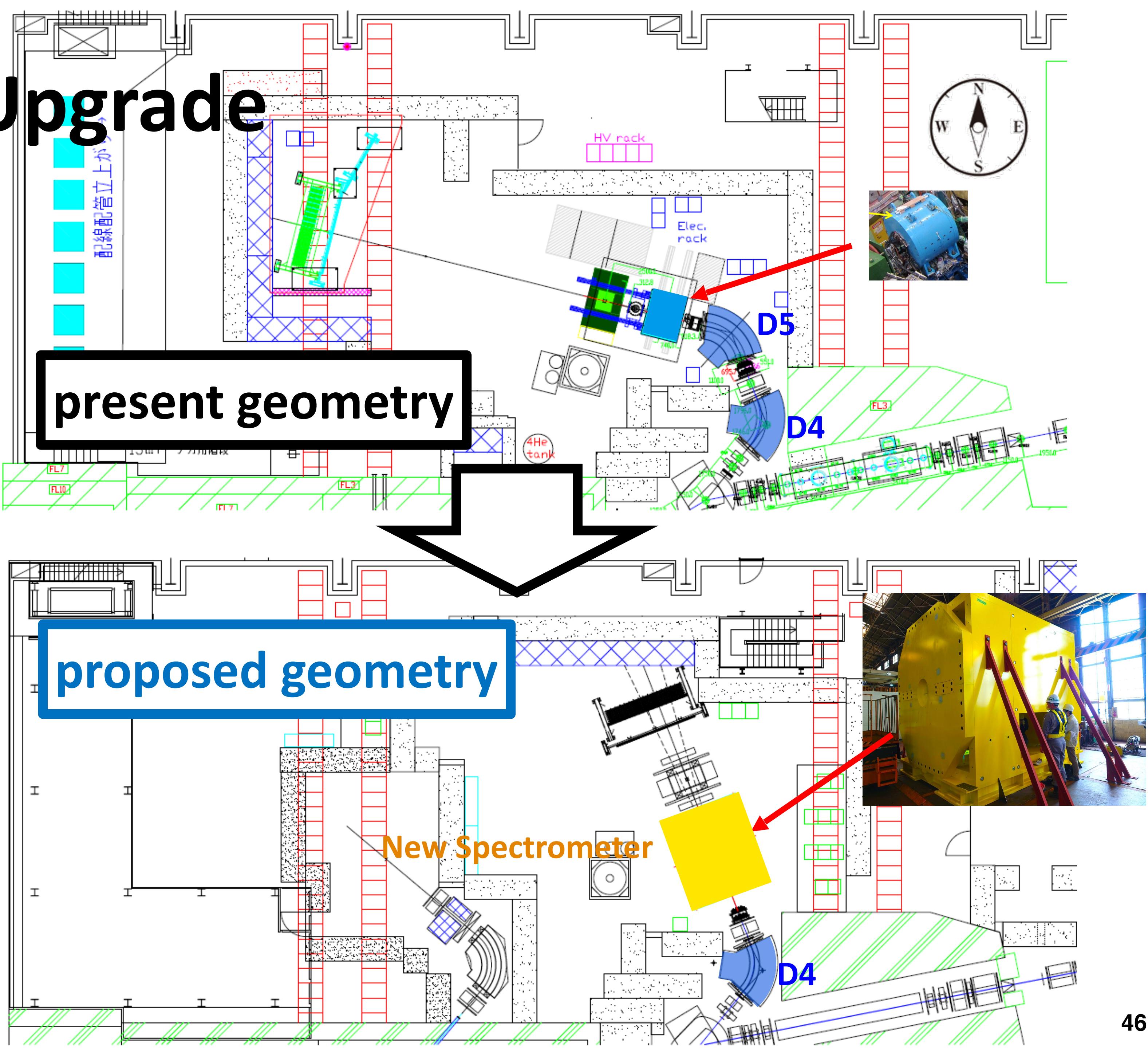
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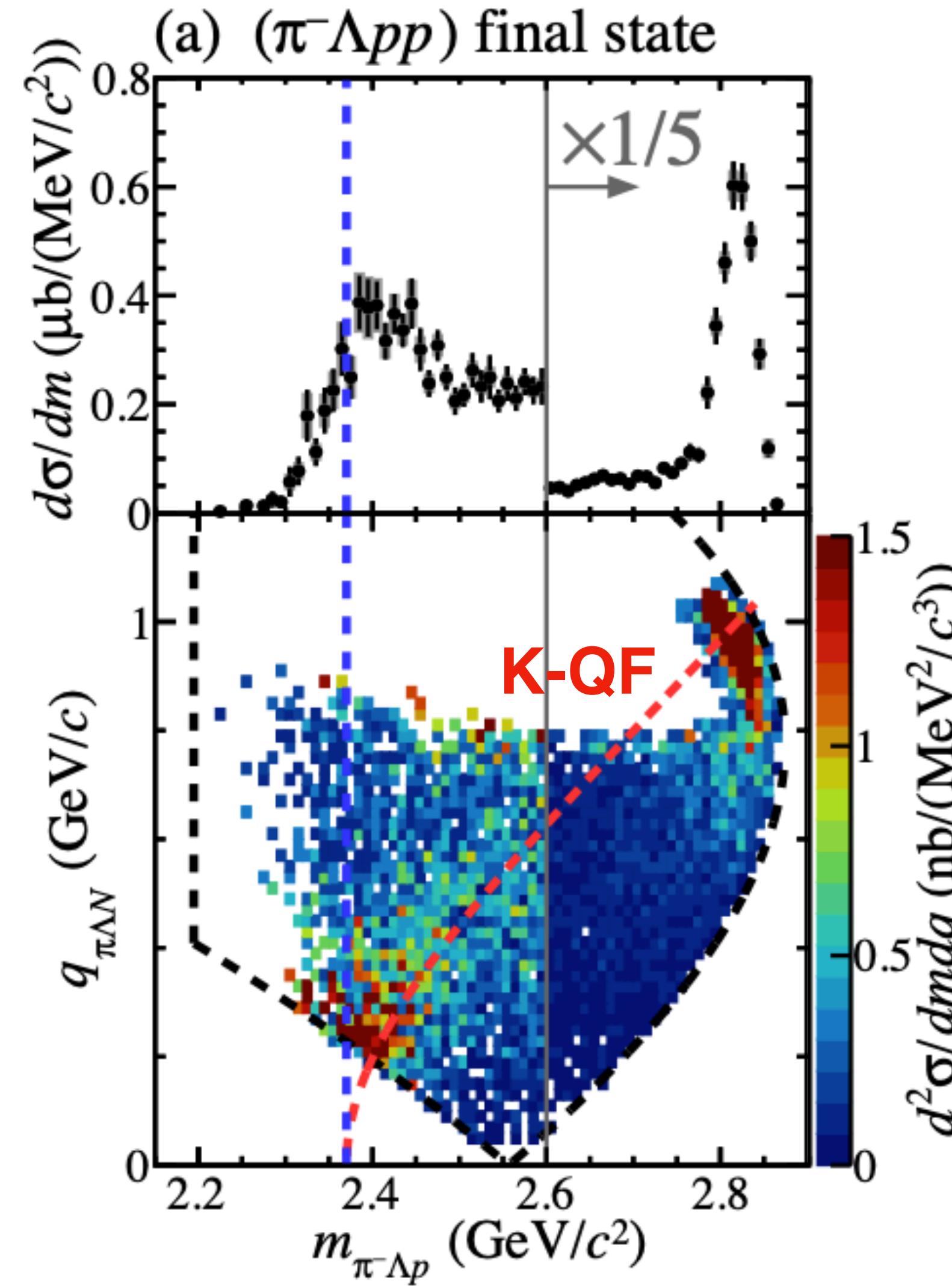
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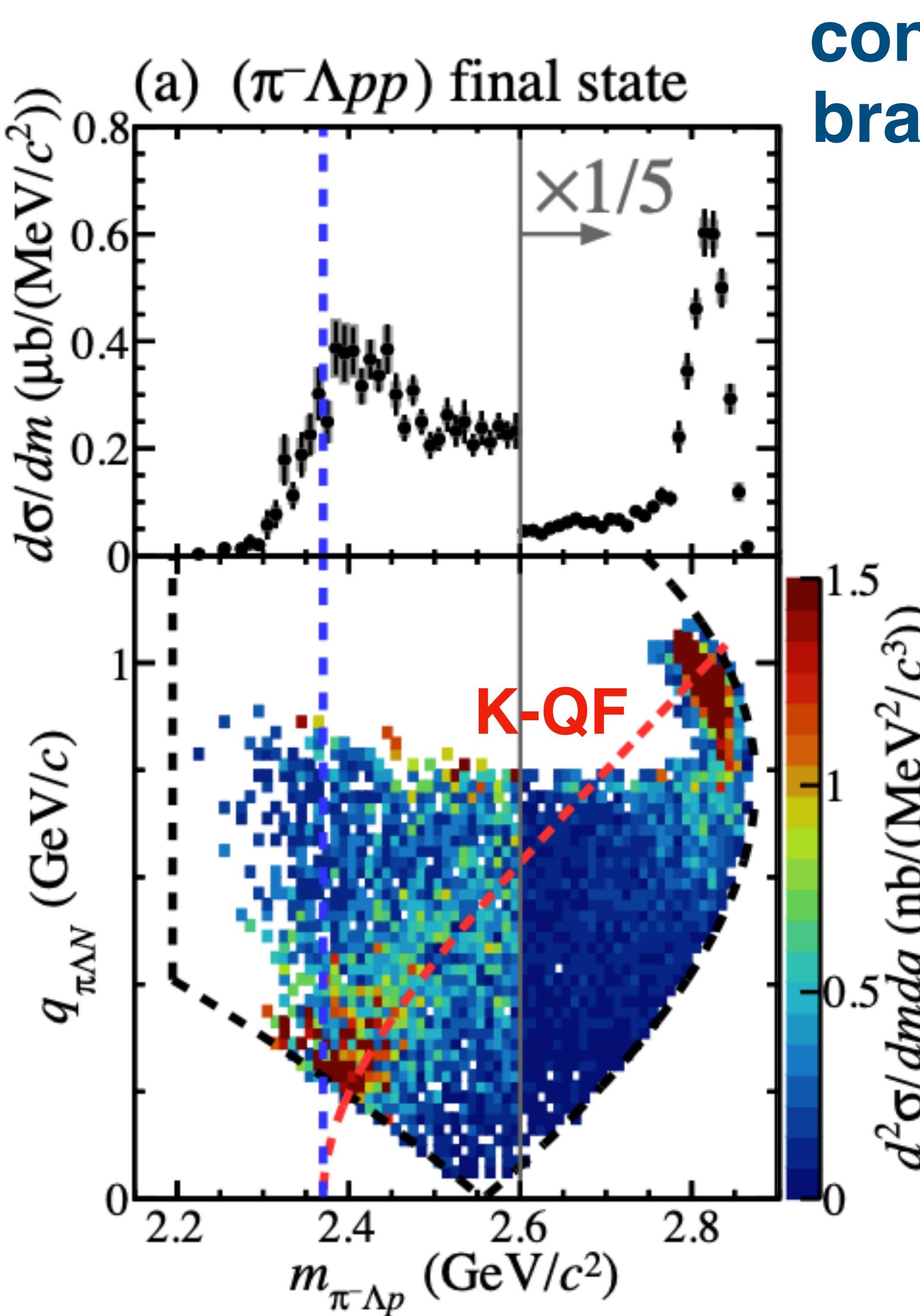


$K^- + {}^3He \rightarrow (\pi^- \Lambda p) + p$ reaction



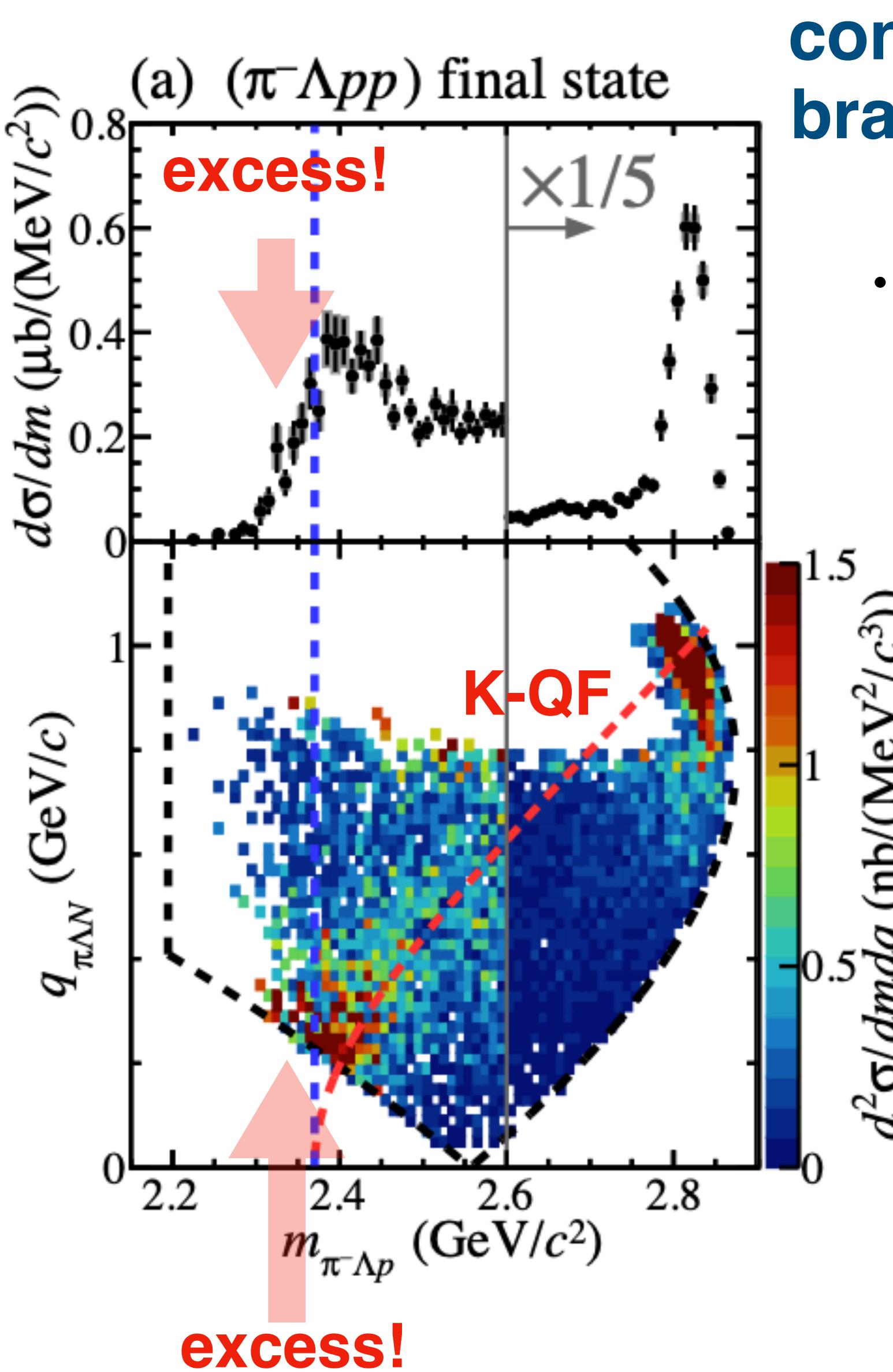
... analyzed by T. Yamaga

$K^- + {}^3He \rightarrow (\pi^- \Lambda p) + p$ reaction



consistent with $K^- + {}^3He \rightarrow \Lambda pn$ reaction
branch seems to be oder bigger

$K^- + {}^3He \rightarrow (\pi^- \Lambda p) + p$ reaction



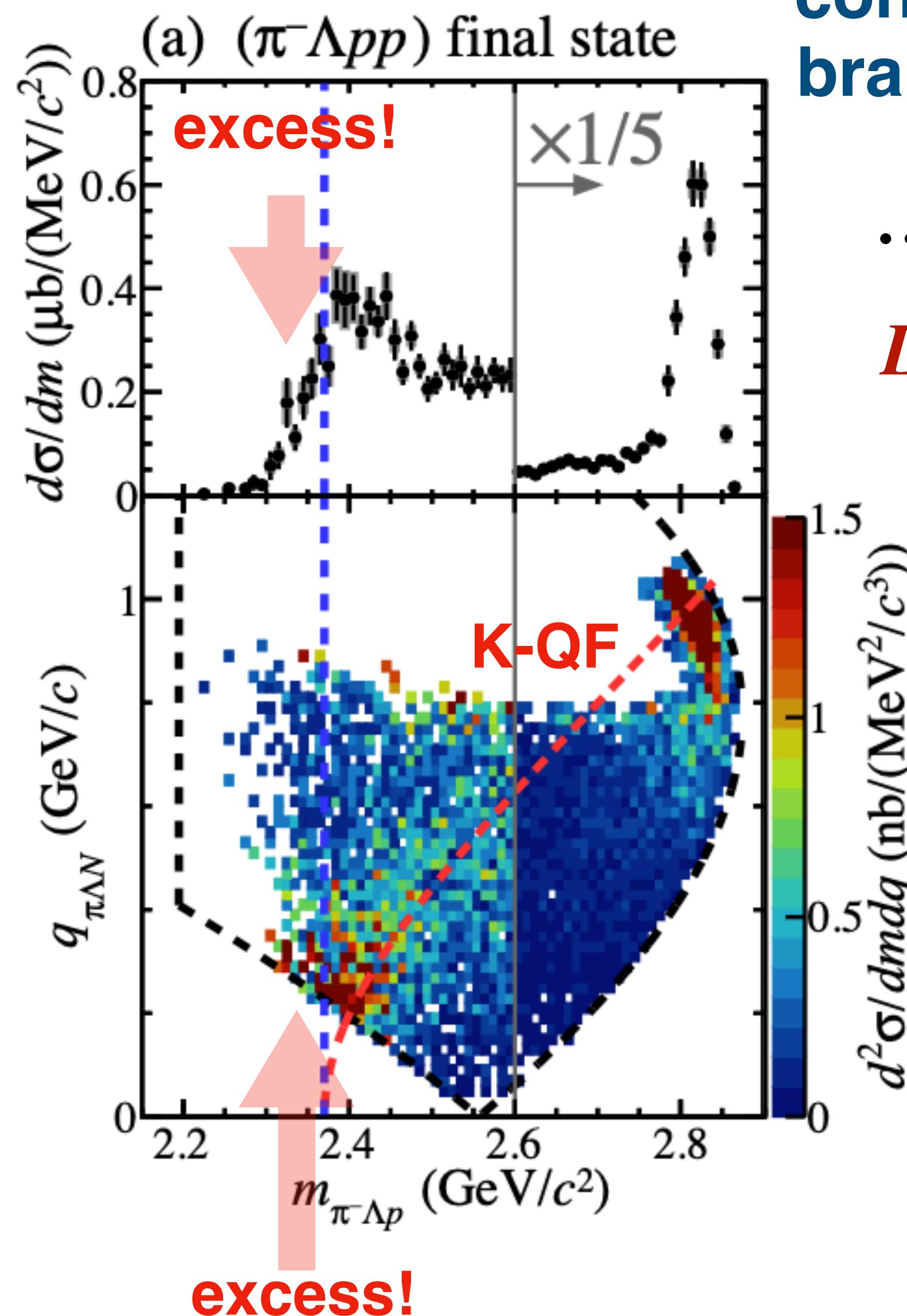
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... *excess is not easy to see ...*

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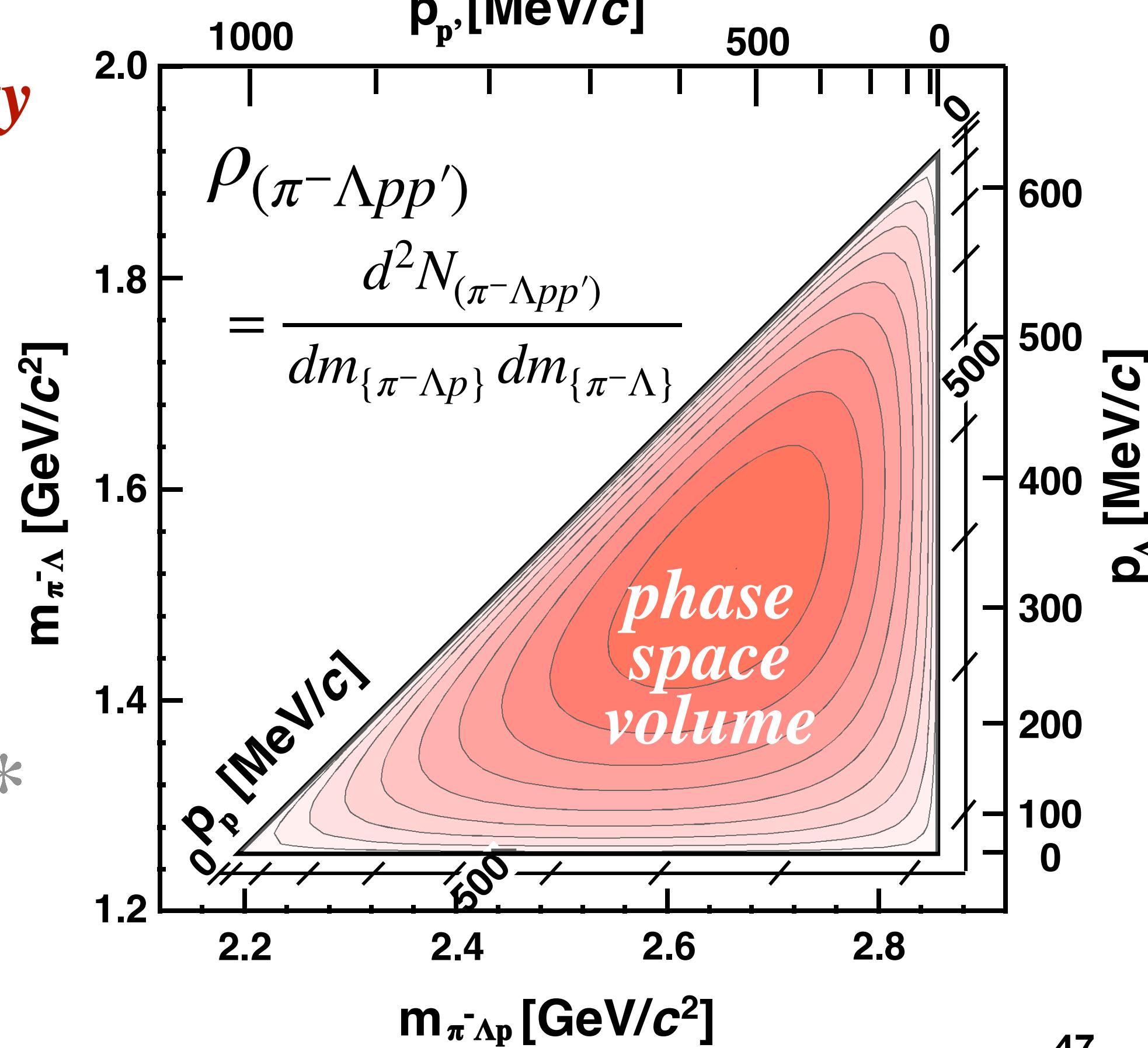
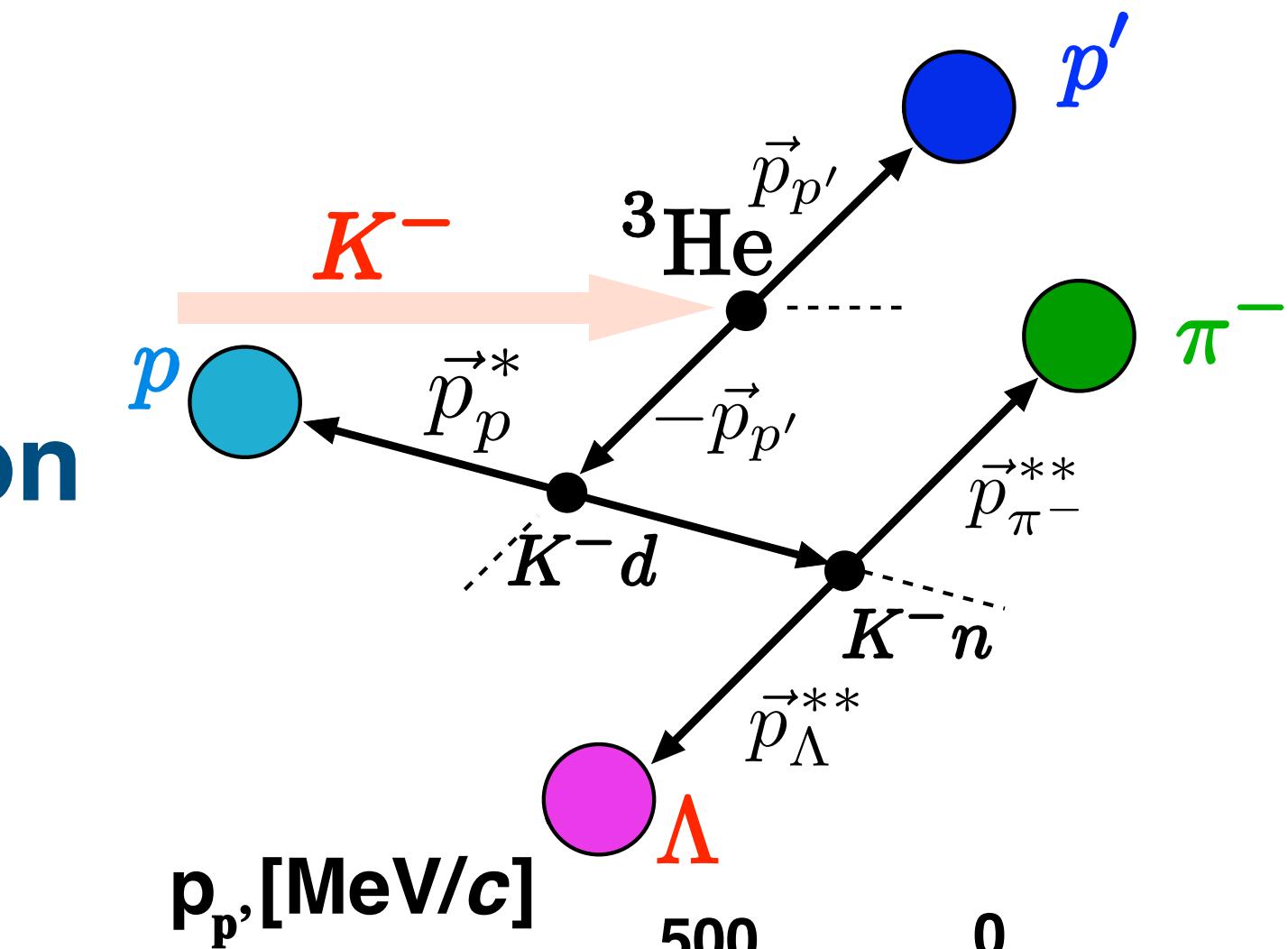
... excess is not easy to see ...

Let's normalize event density by 4-body phase space

The normalization by 4-body phase space, i.e., final-state-density

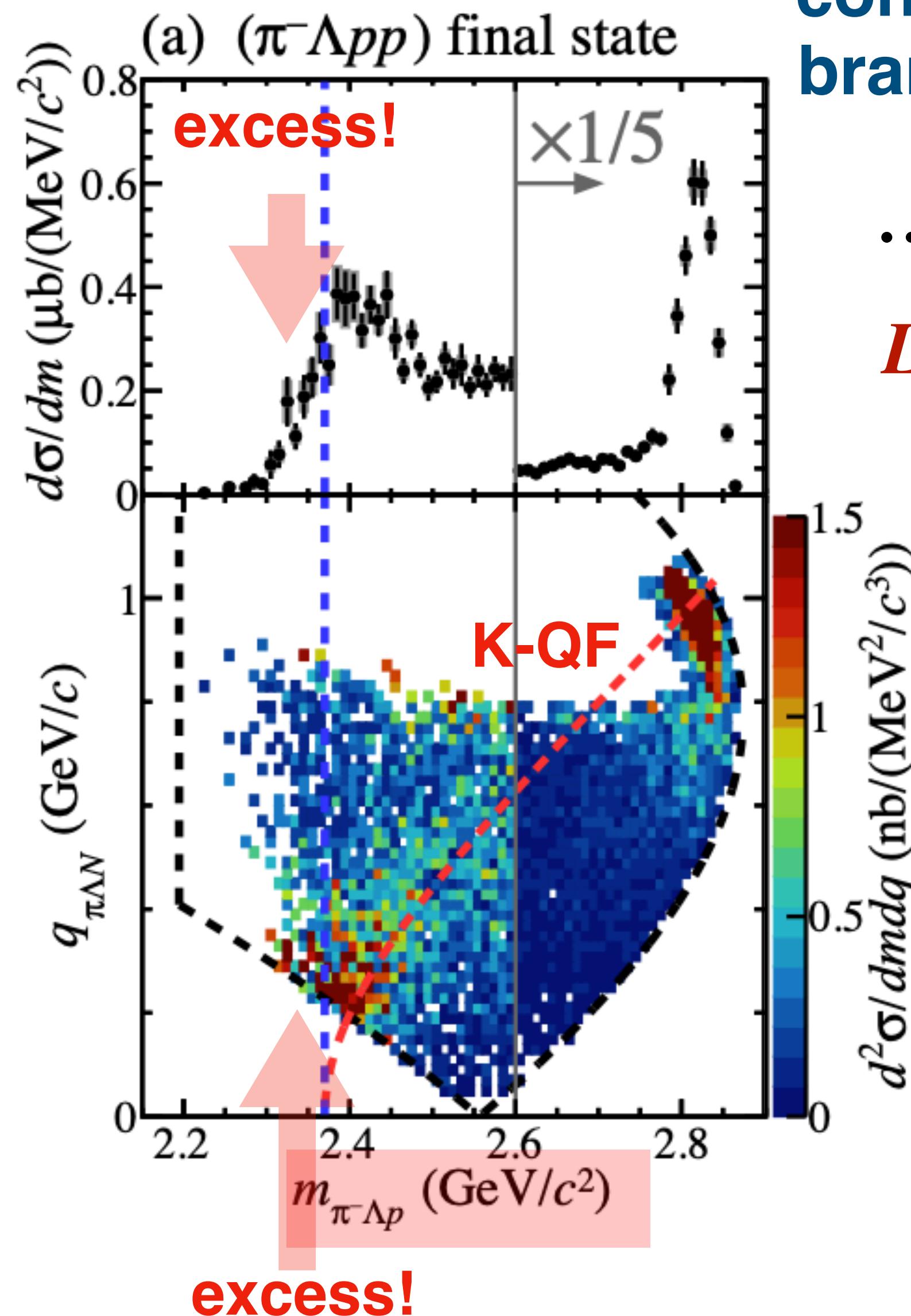
$$\rho_{(\pi^-\Lambda pp')} = \frac{d^2N_{(\pi^-\Lambda pp')}}{dm_{\{\pi^-\Lambda p\}} dm_{\{\pi^-\Lambda\}}} \propto p_{p'} \times p_p^* \times p_\Lambda^{**}$$

... analyzed by T. Yamaga



$K^- + {}^3\text{He} \rightarrow (\pi^-\Lambda p) + p$ reaction

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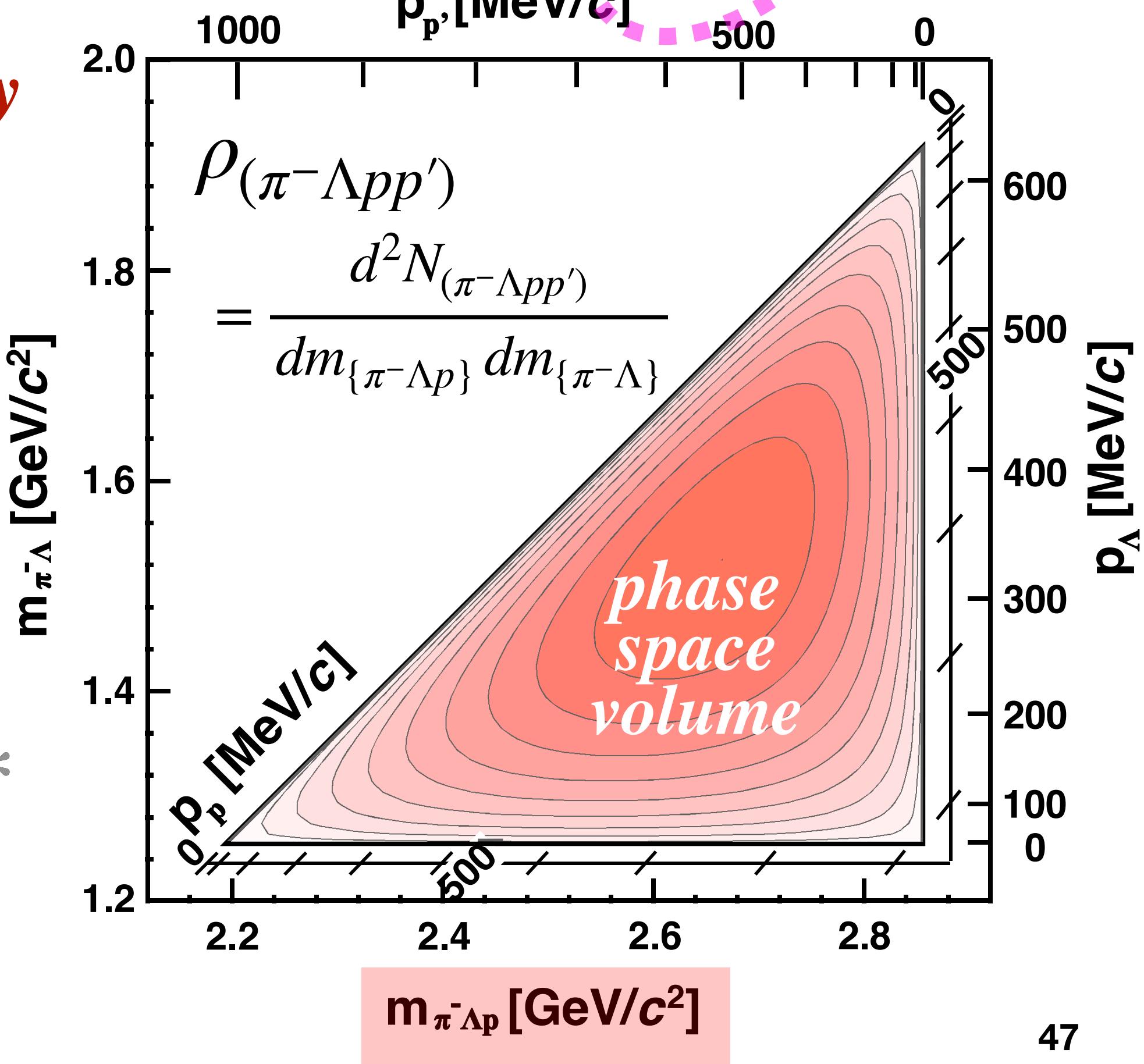
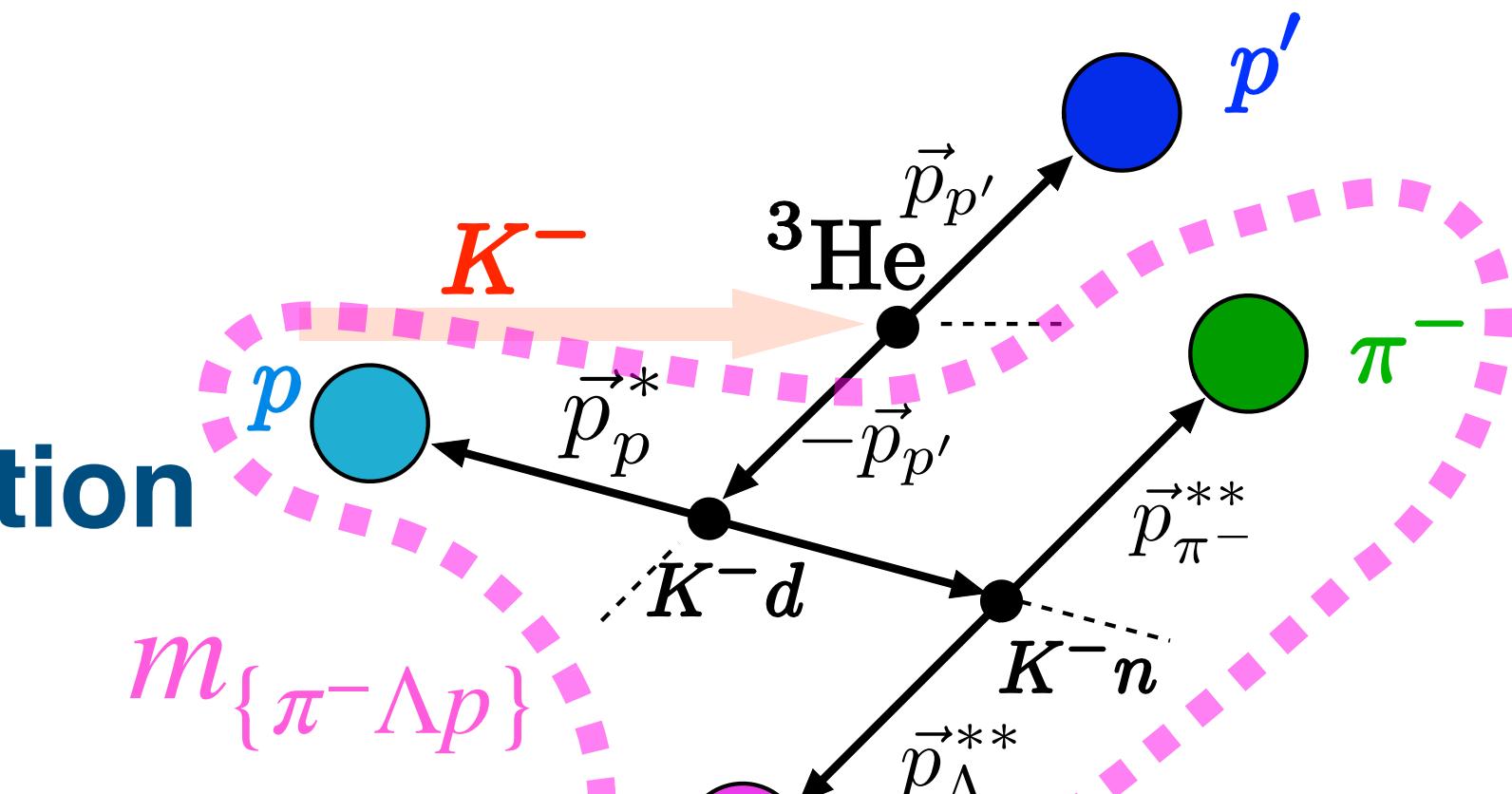
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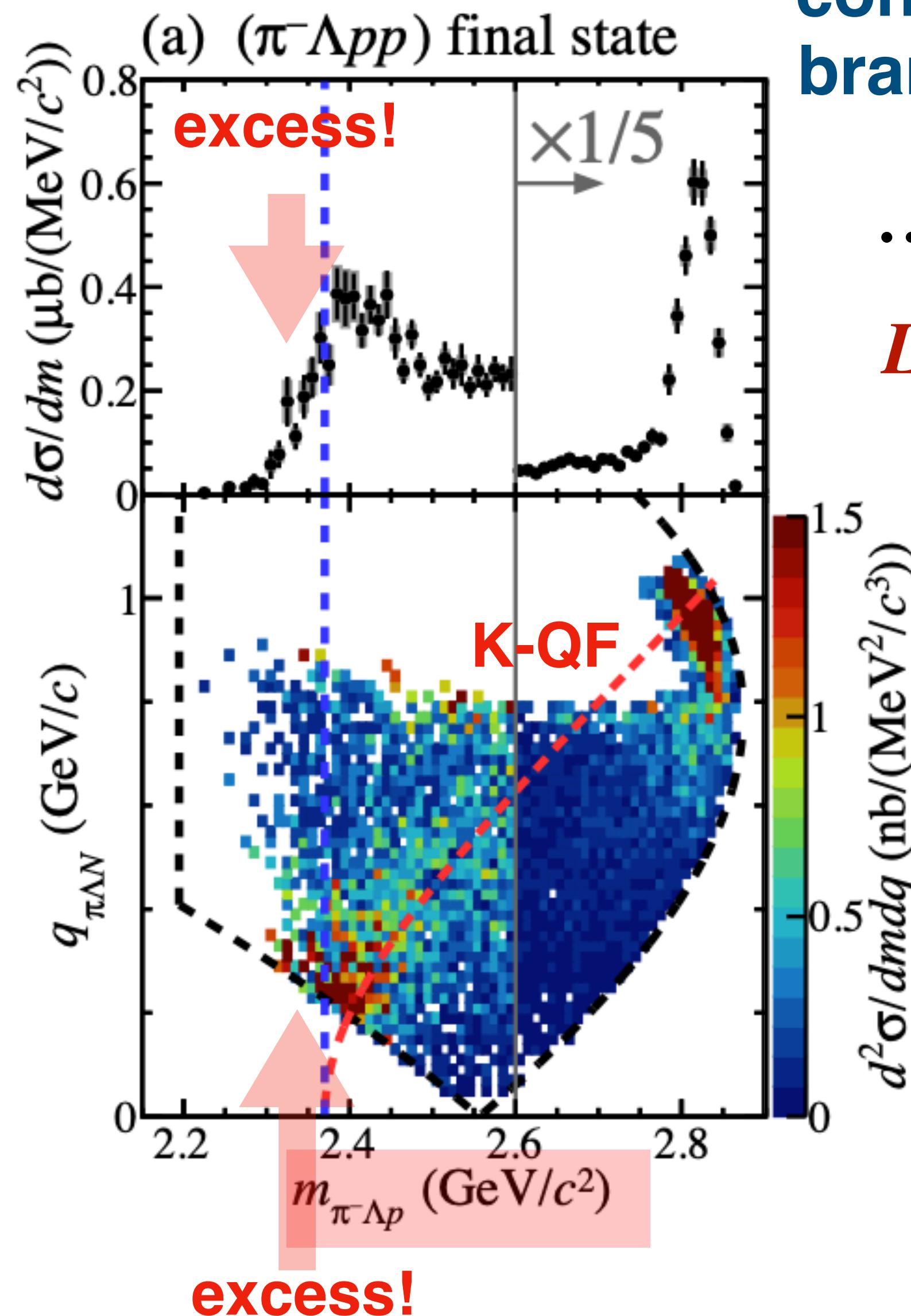
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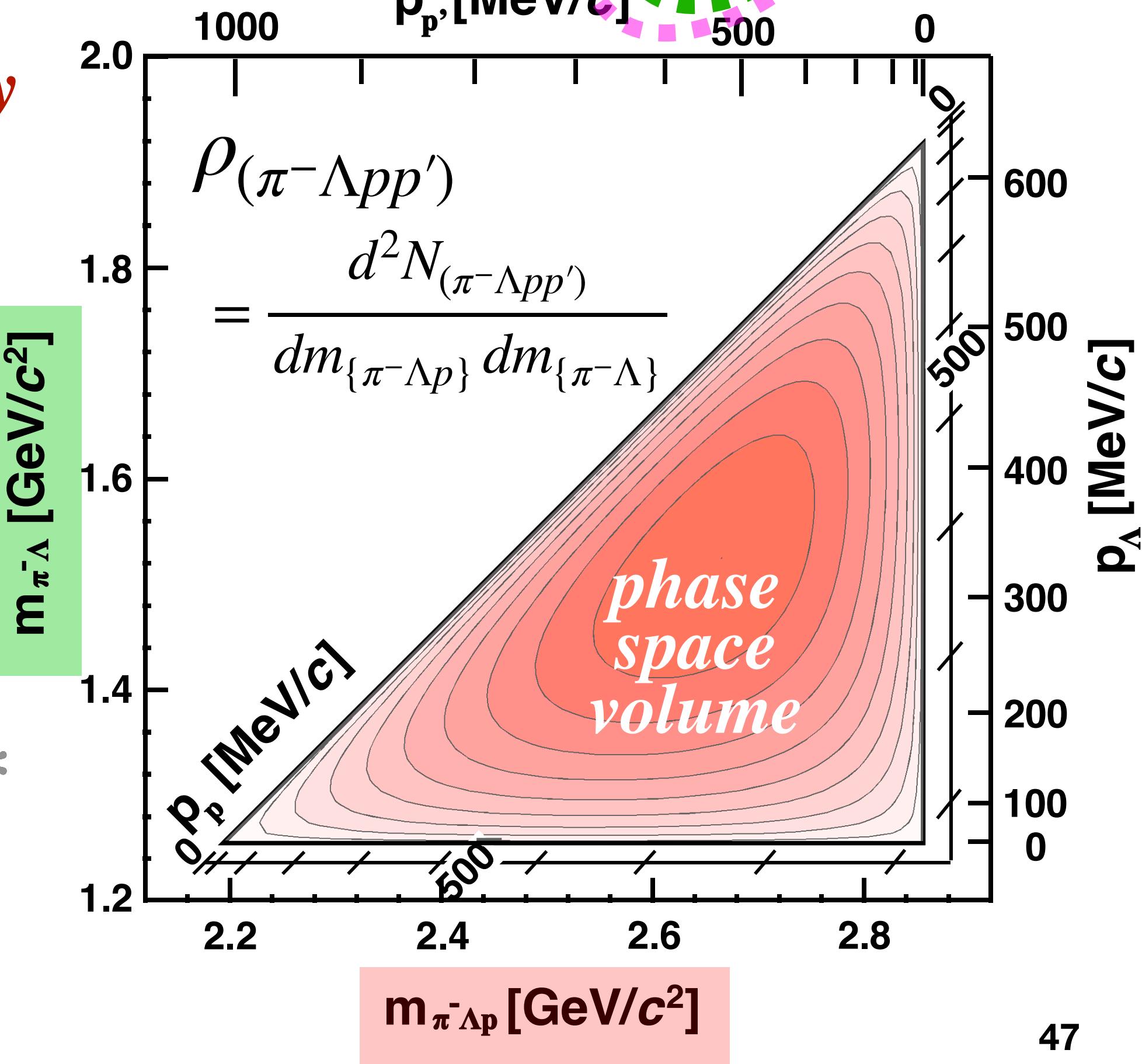
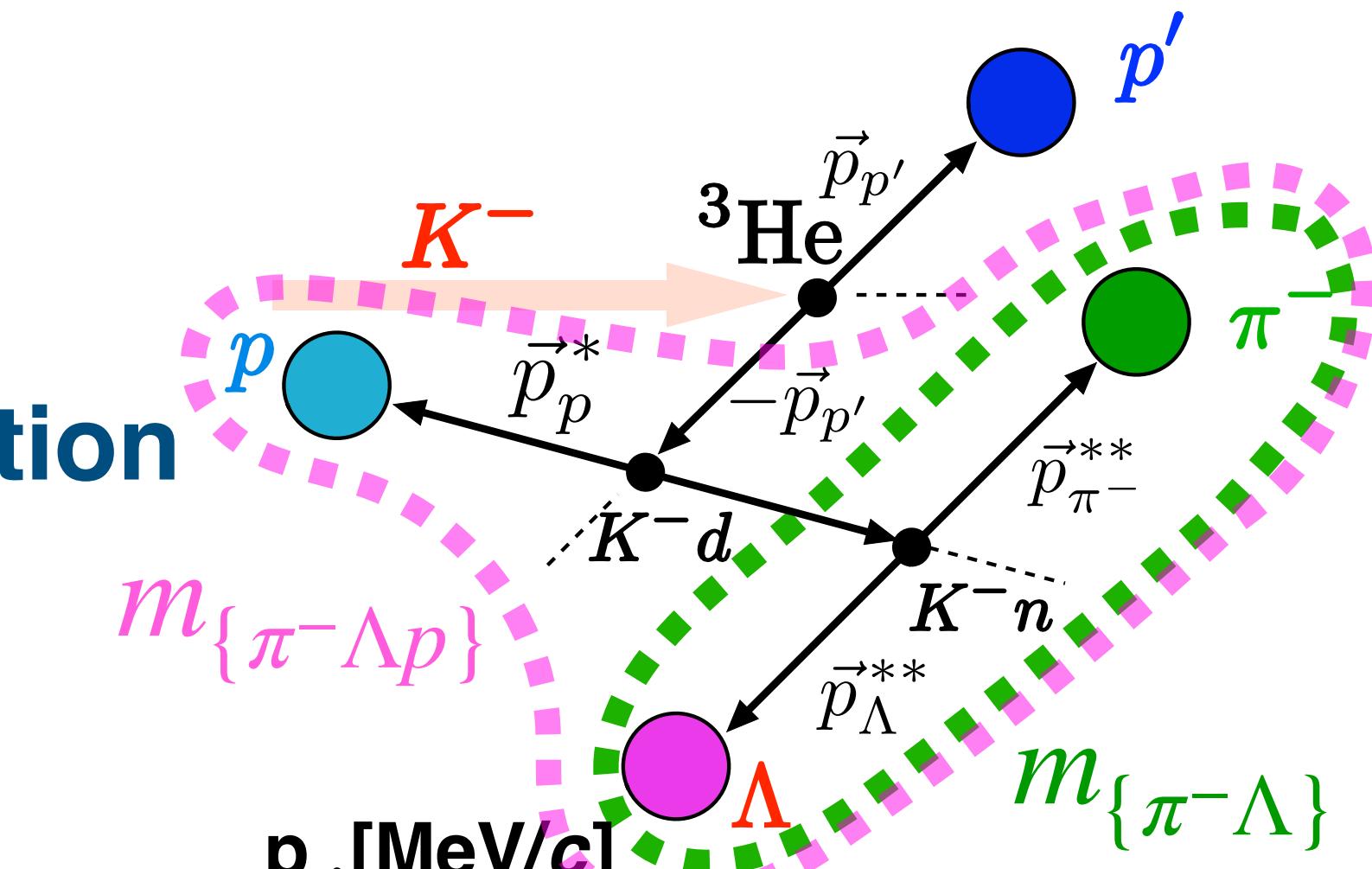
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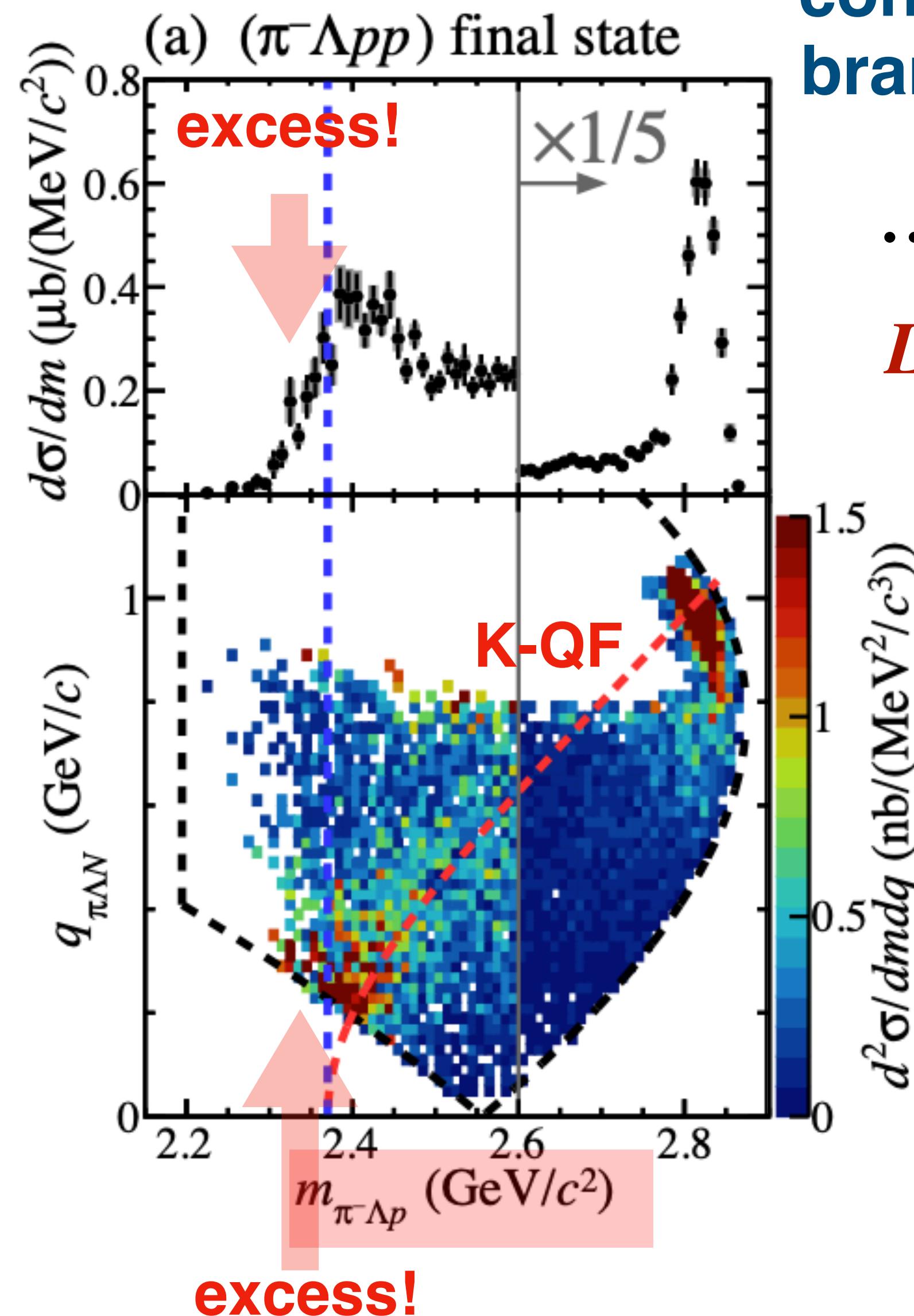
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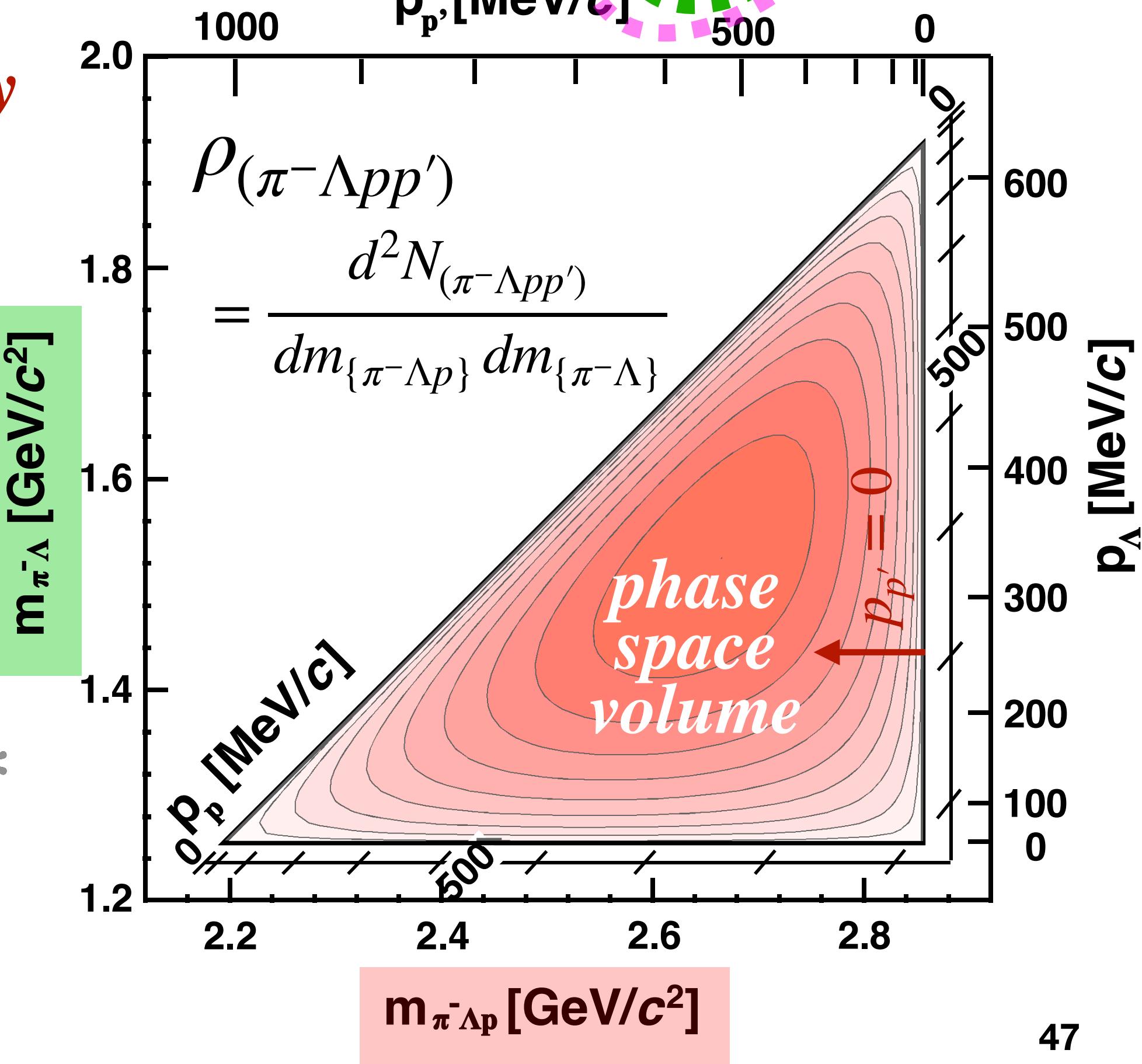
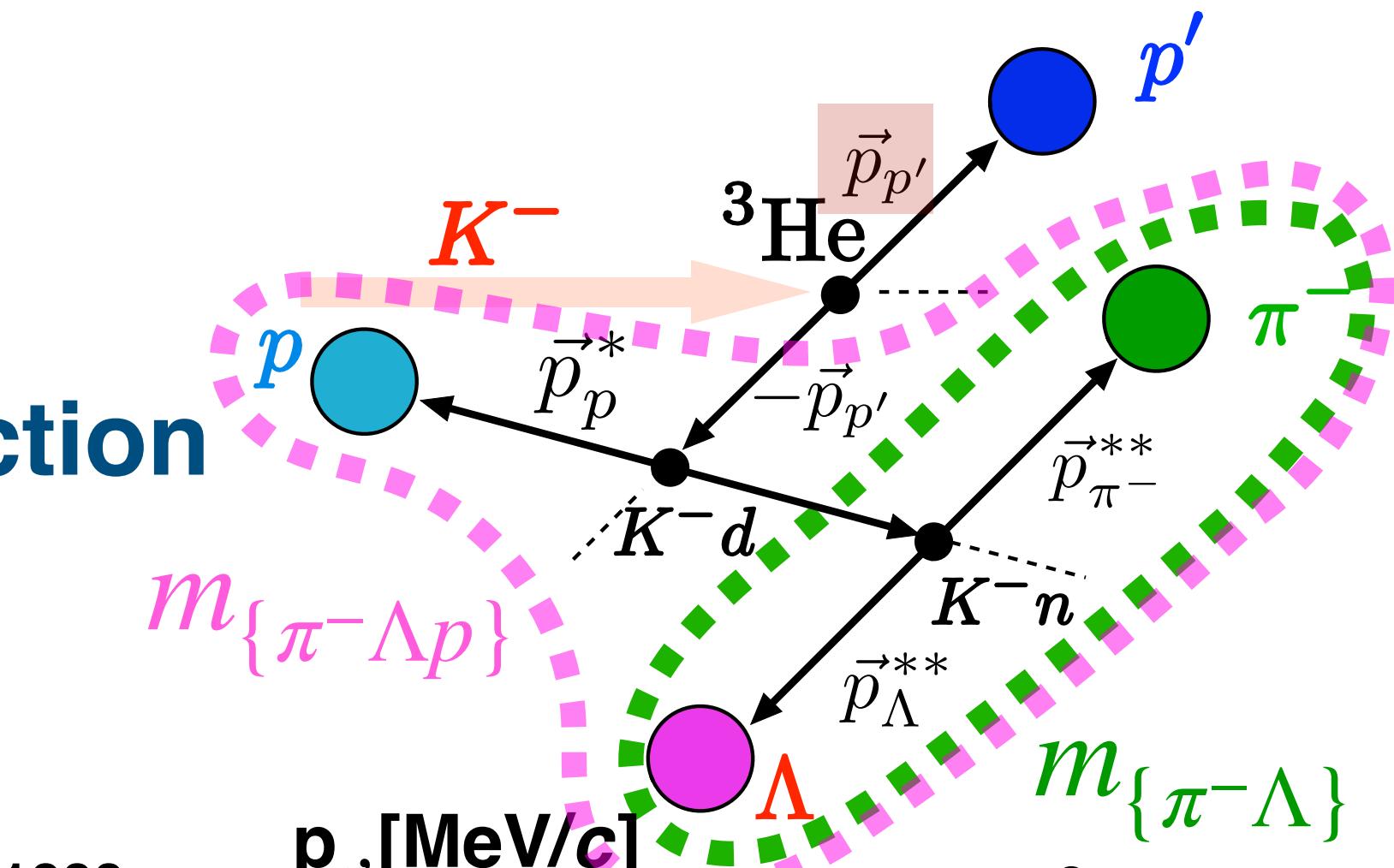
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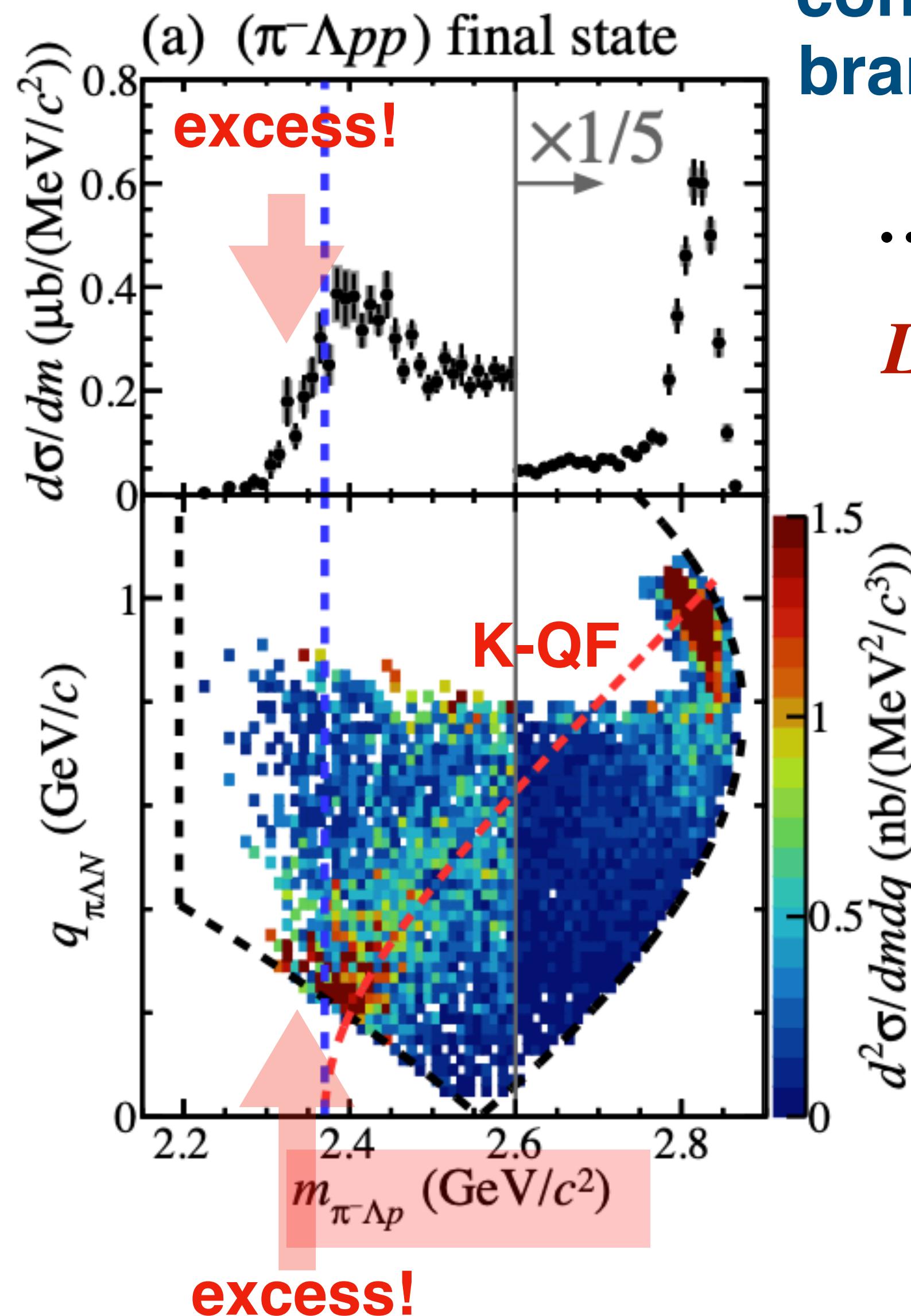
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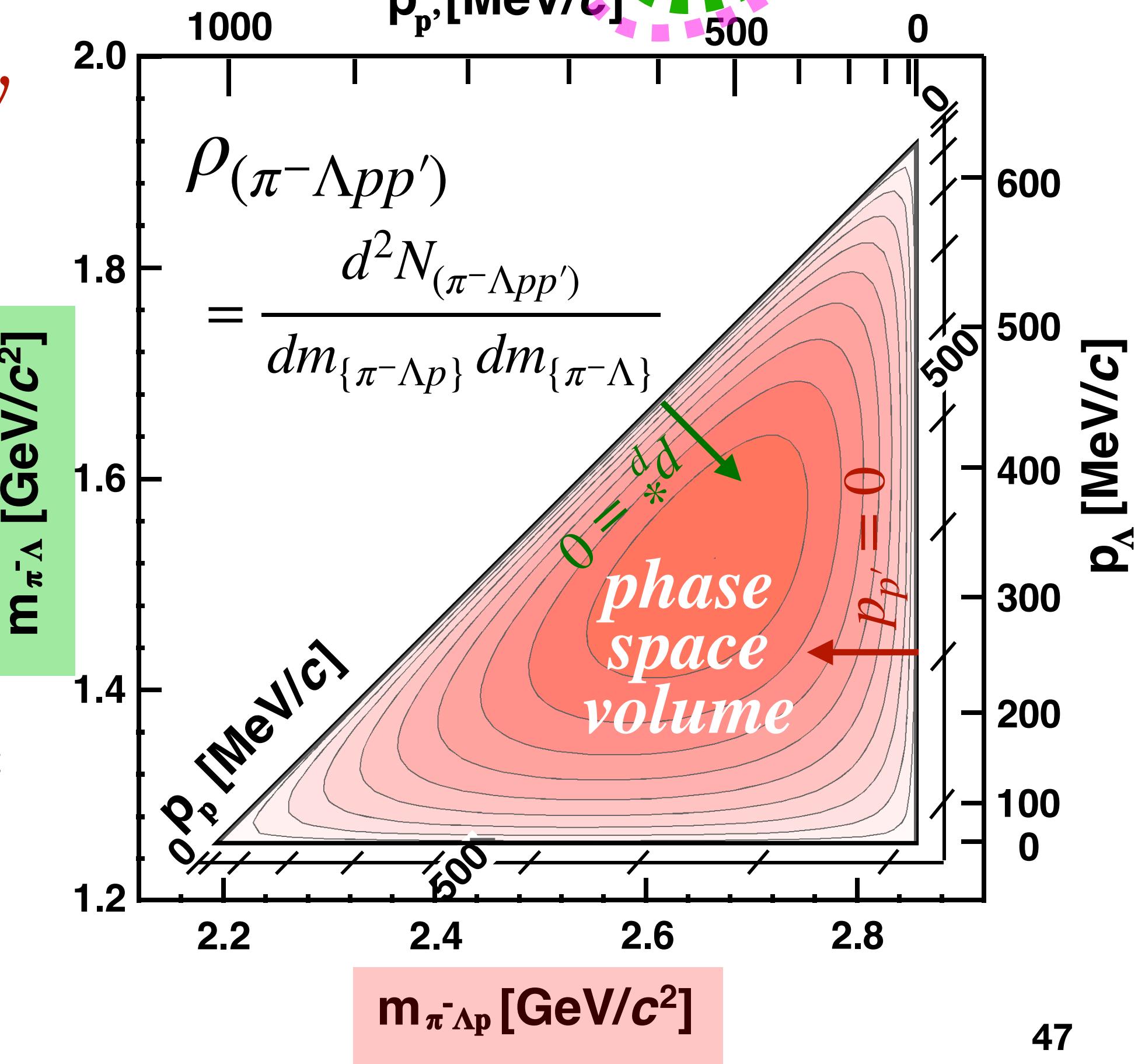
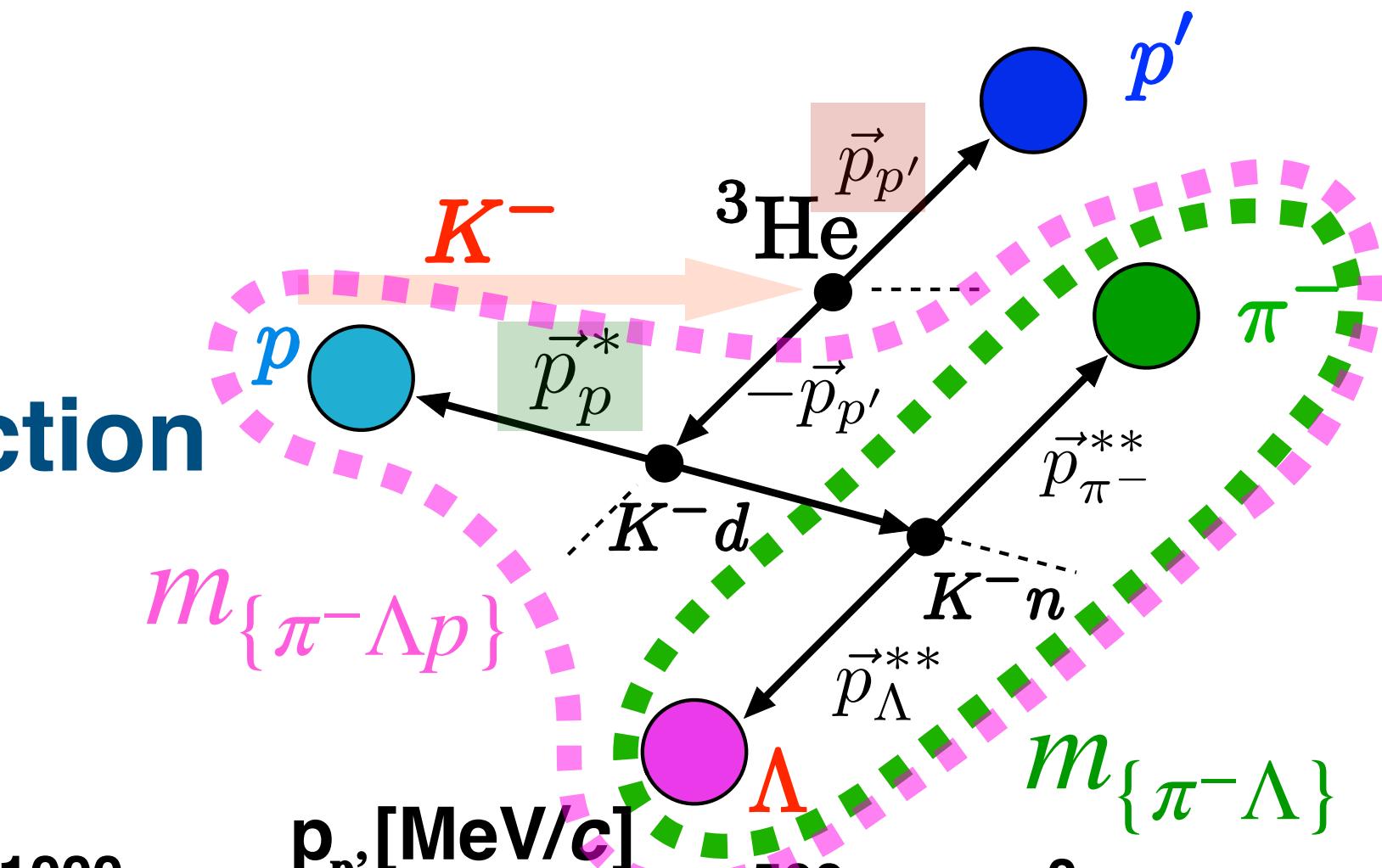
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The normalization by 4-body phase space, i.e., final-state-density

$$\rho_{(\pi^-\Lambda pp')} = \frac{d^2N_{(\pi^-\Lambda pp')}}{dm_{\{\pi^-\Lambda p\}} dm_{\{\pi^-\Lambda\}}}$$

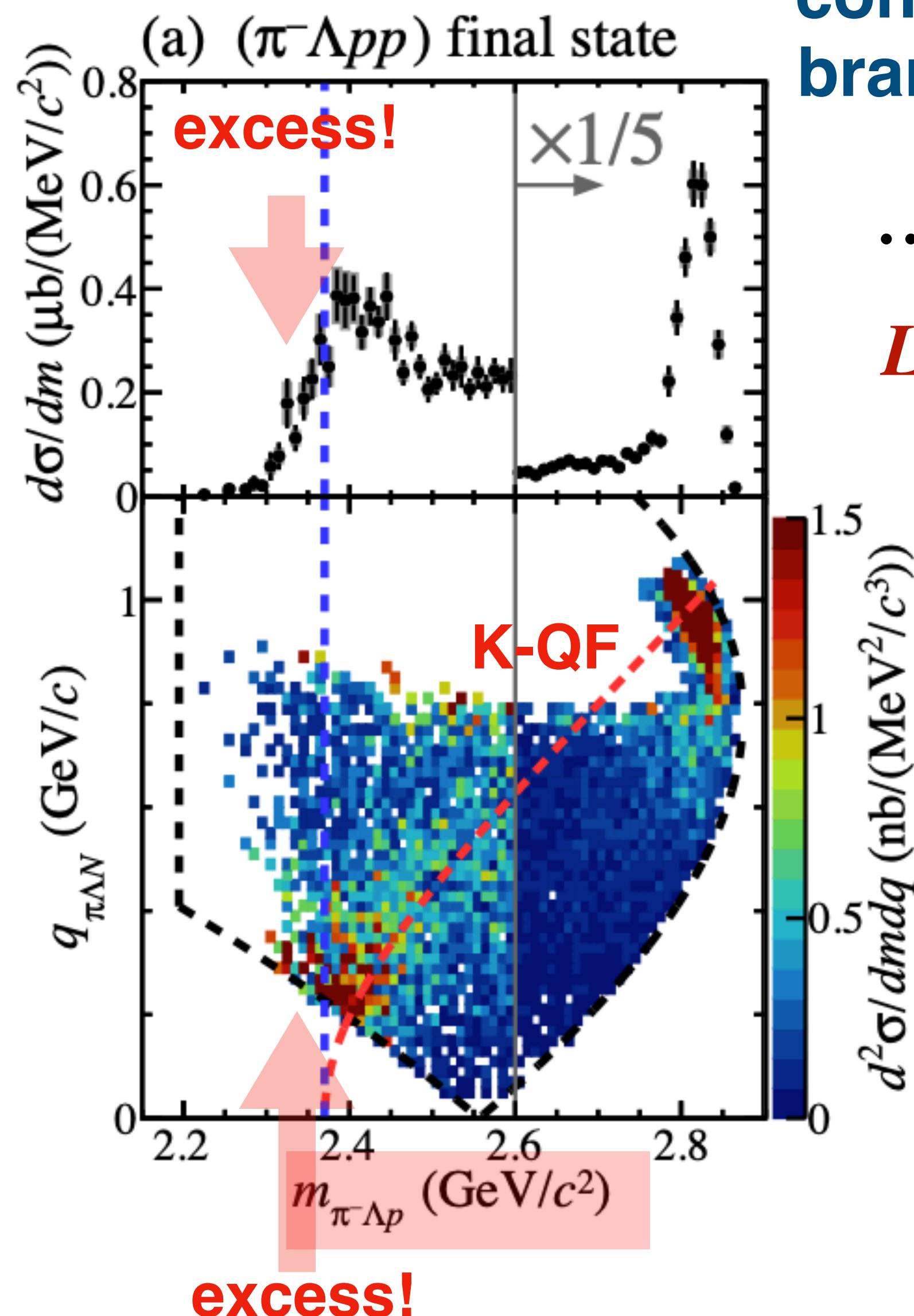
$$\propto p_{p'} \times p_p^* \times p_\Lambda^{**}$$

... analyzed by T. Yamaga



$K^- + {}^3He \rightarrow (\pi^- \Lambda p) + p$ reaction

consistent with $K^- + {}^3He \rightarrow \Lambda pn$ reaction
branch seems to be oder bigger



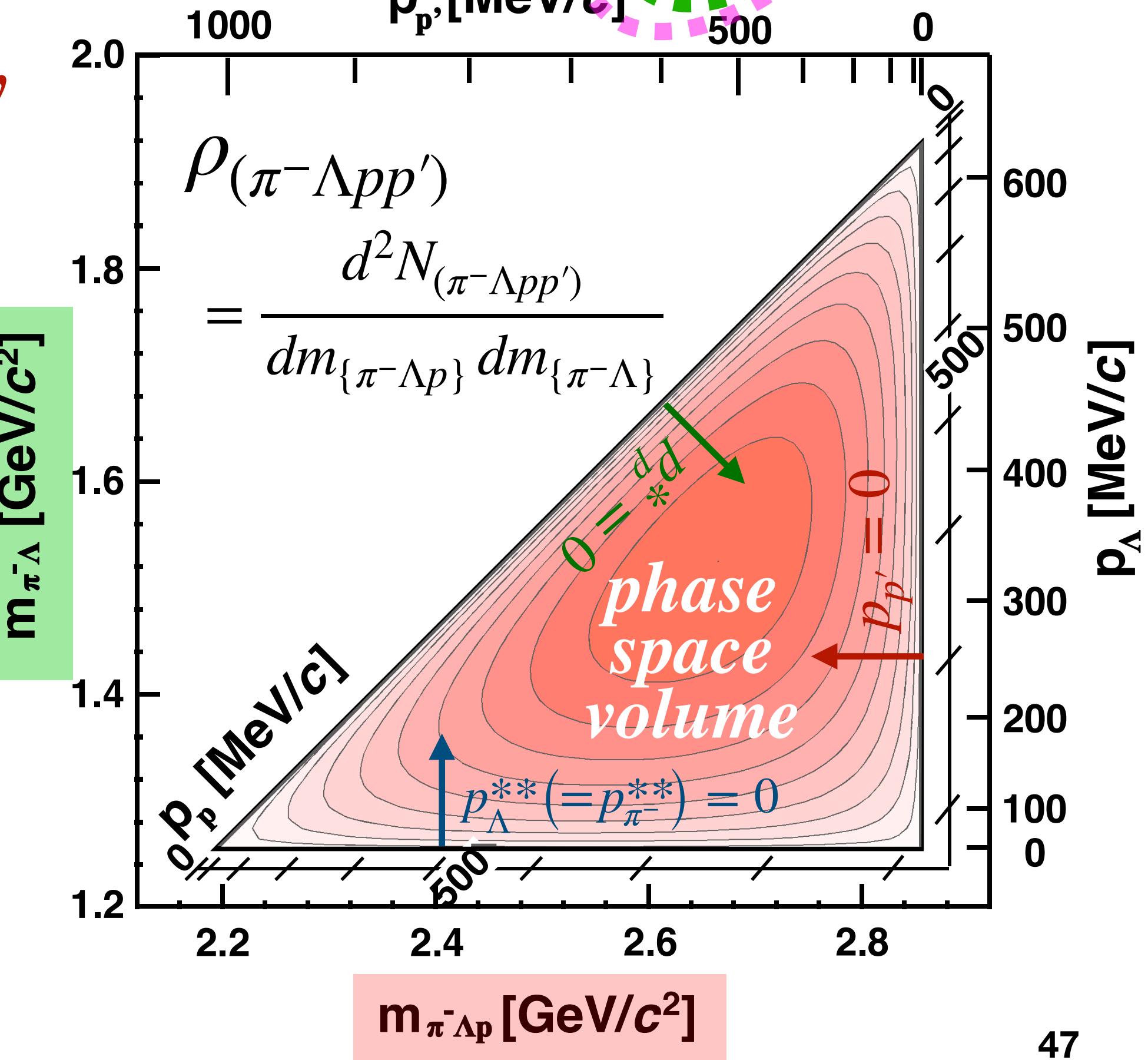
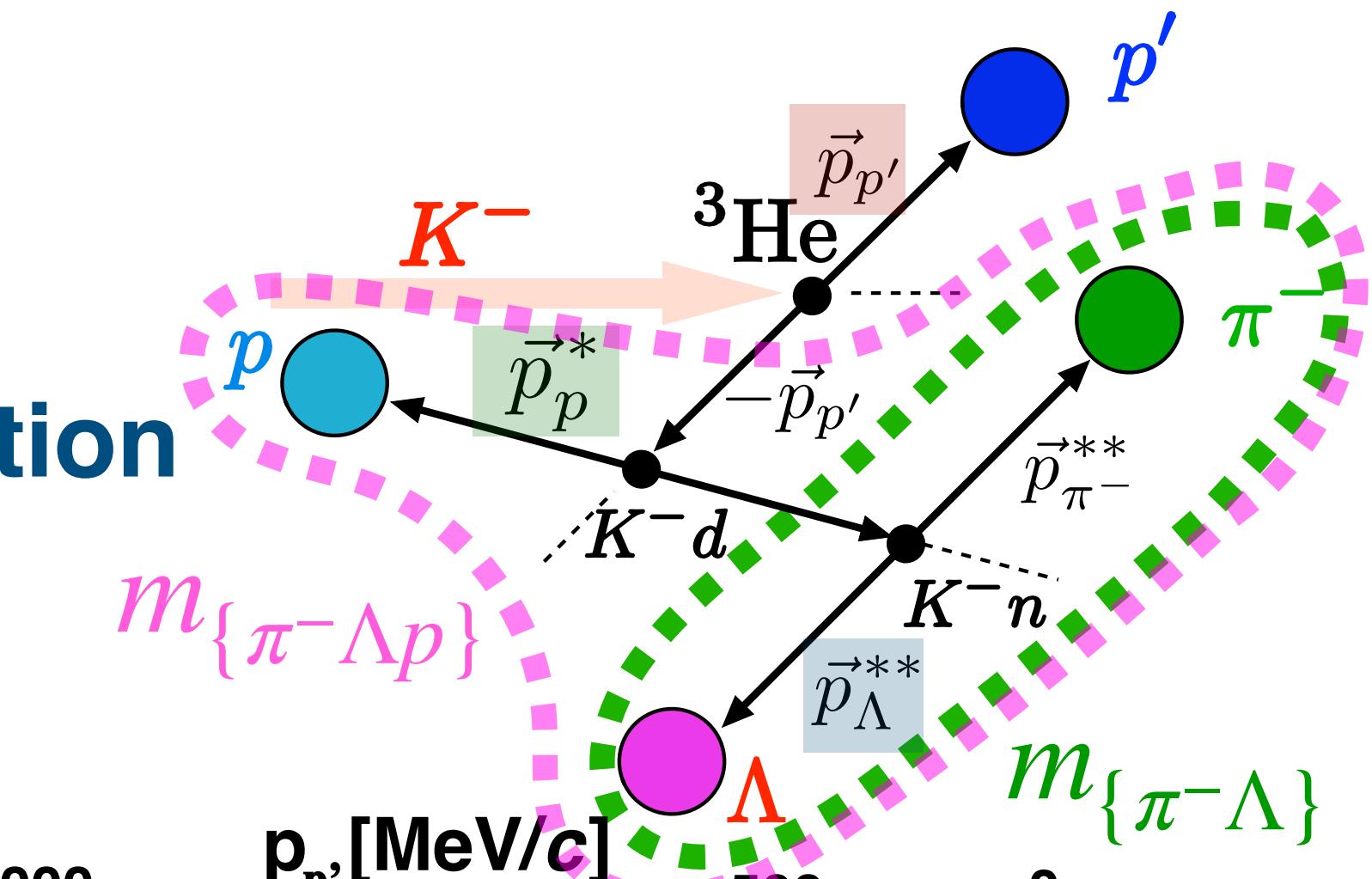
... excess is not easy to see ...

Let's normalize event density by 4-body phase space

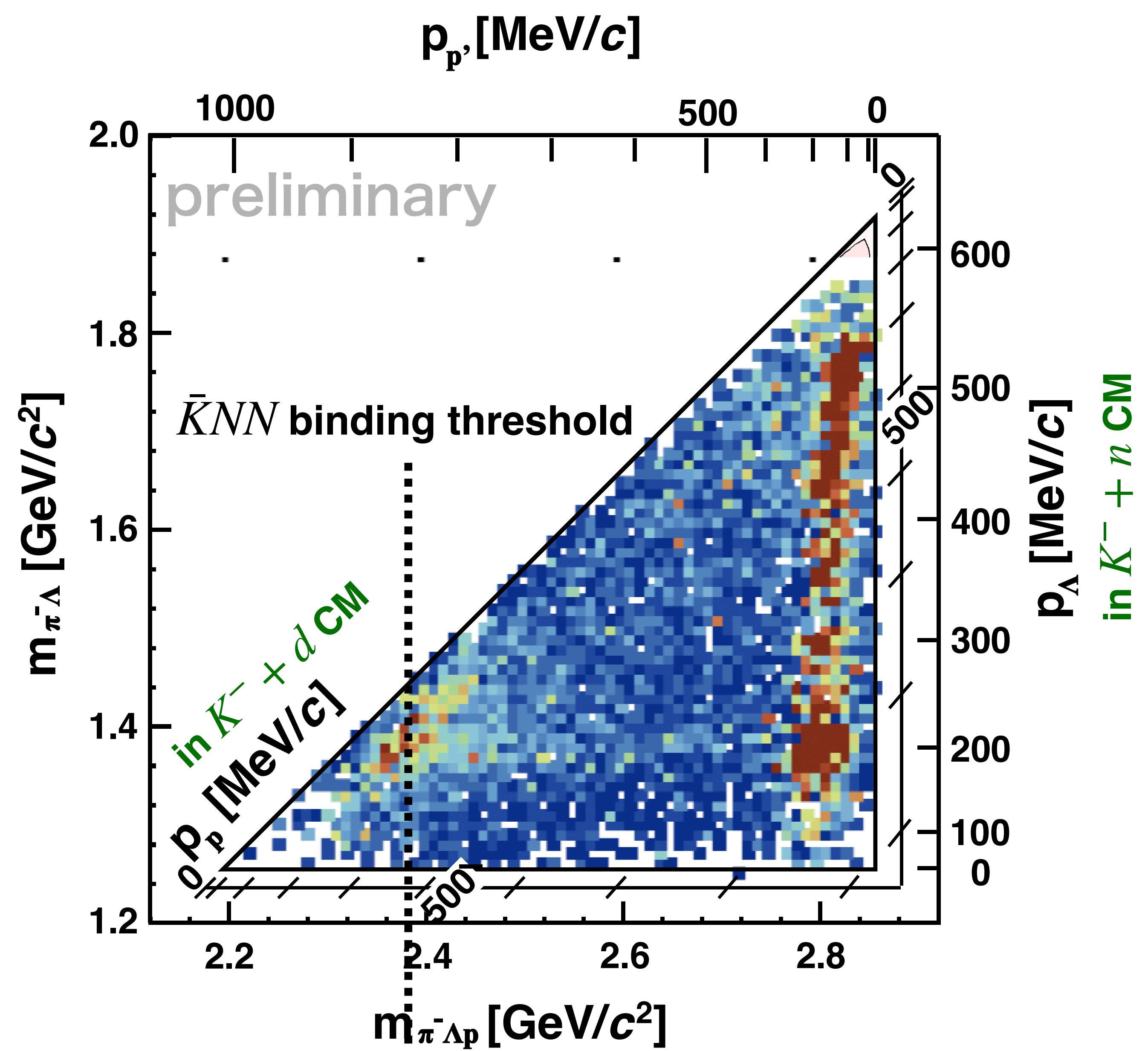
The normalization by 4-body phase space, i.e., final-state-density

$$\rho_{(\pi^-\Lambda pp')} = \frac{d^2N_{(\pi^-\Lambda pp')}}{dm_{\{\pi^-\Lambda p\}} dm_{\{\pi^-\Lambda\}}}$$

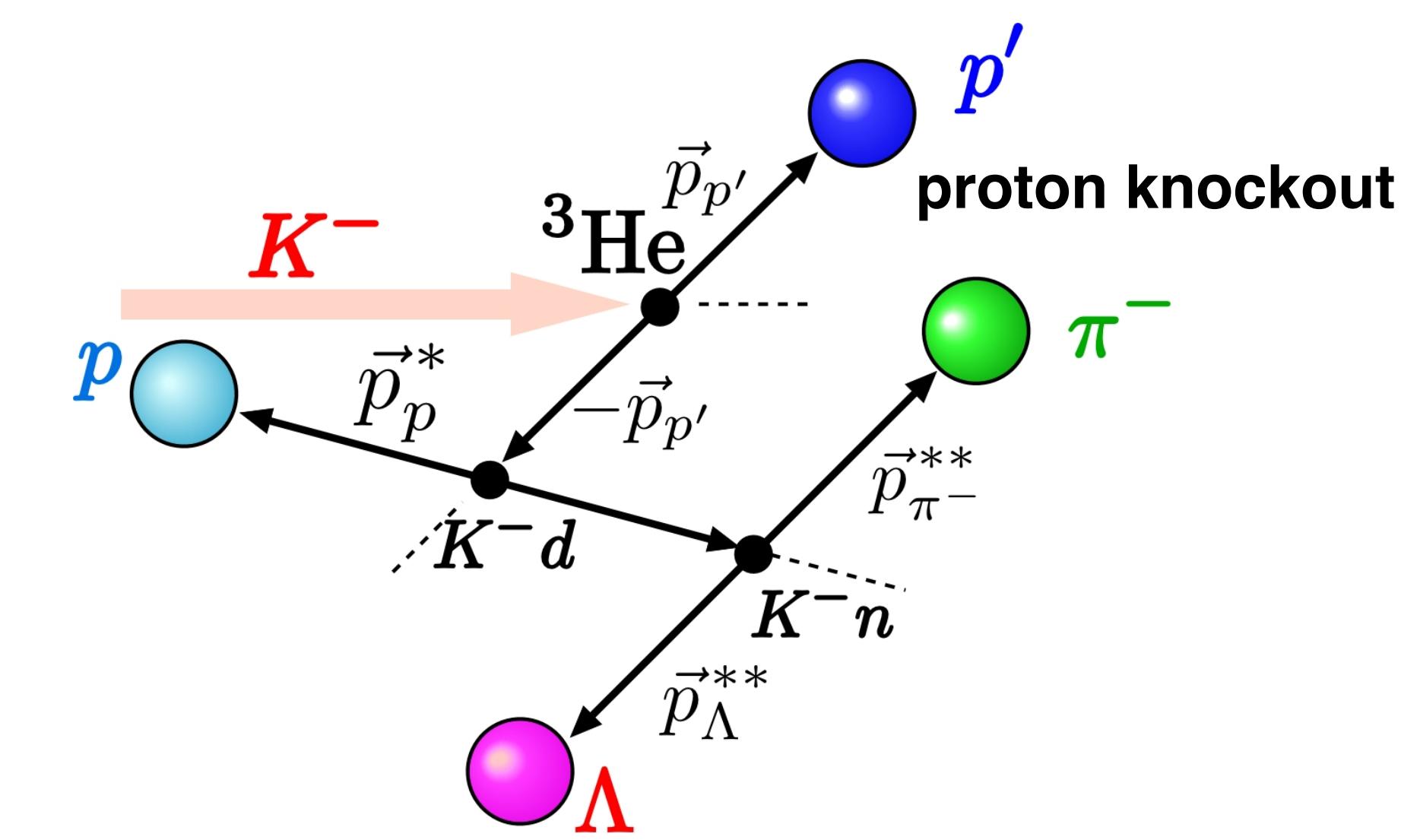
$$\propto p_{p'} \times p_p^* \times p_\Delta^{**}$$

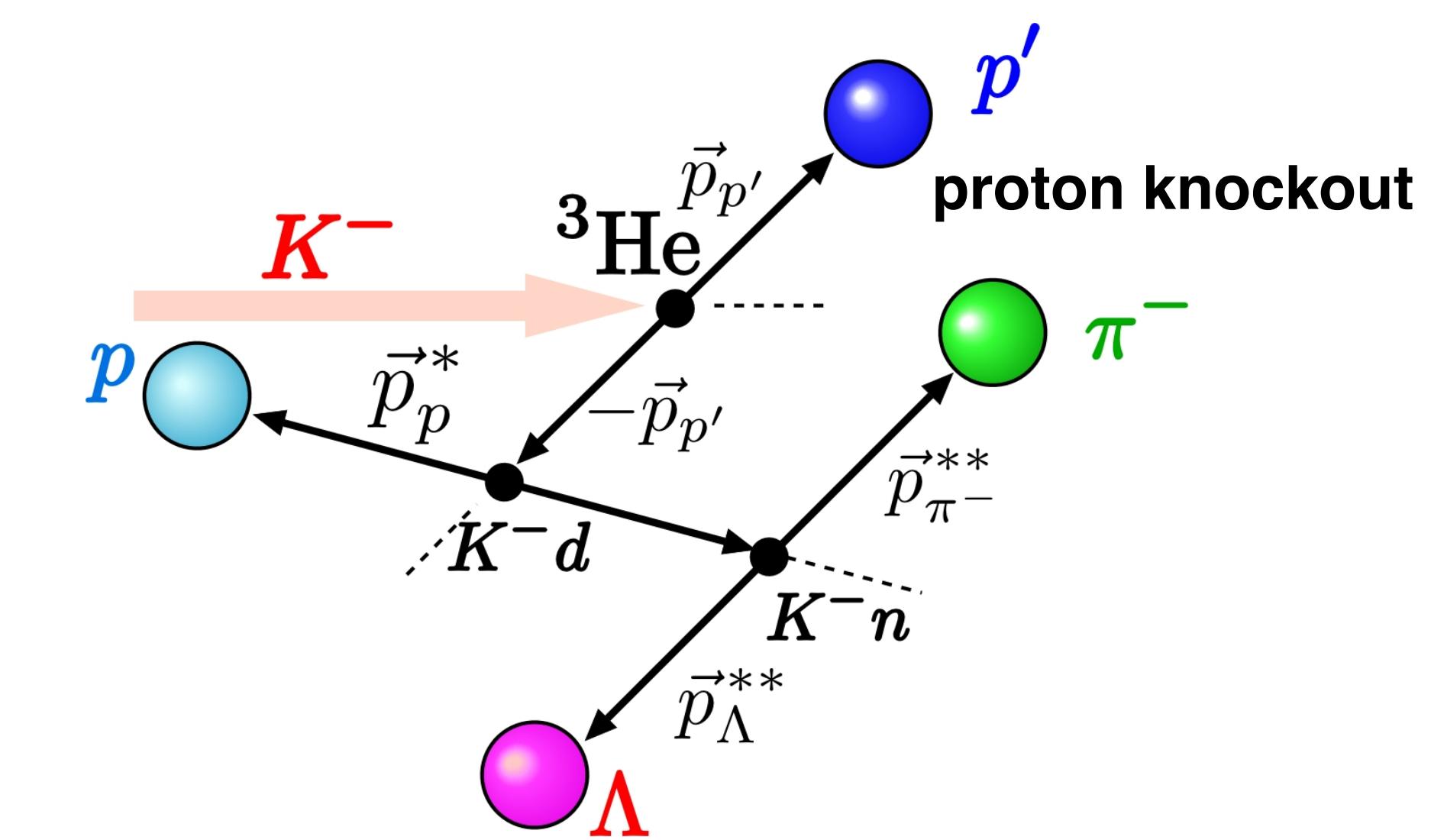
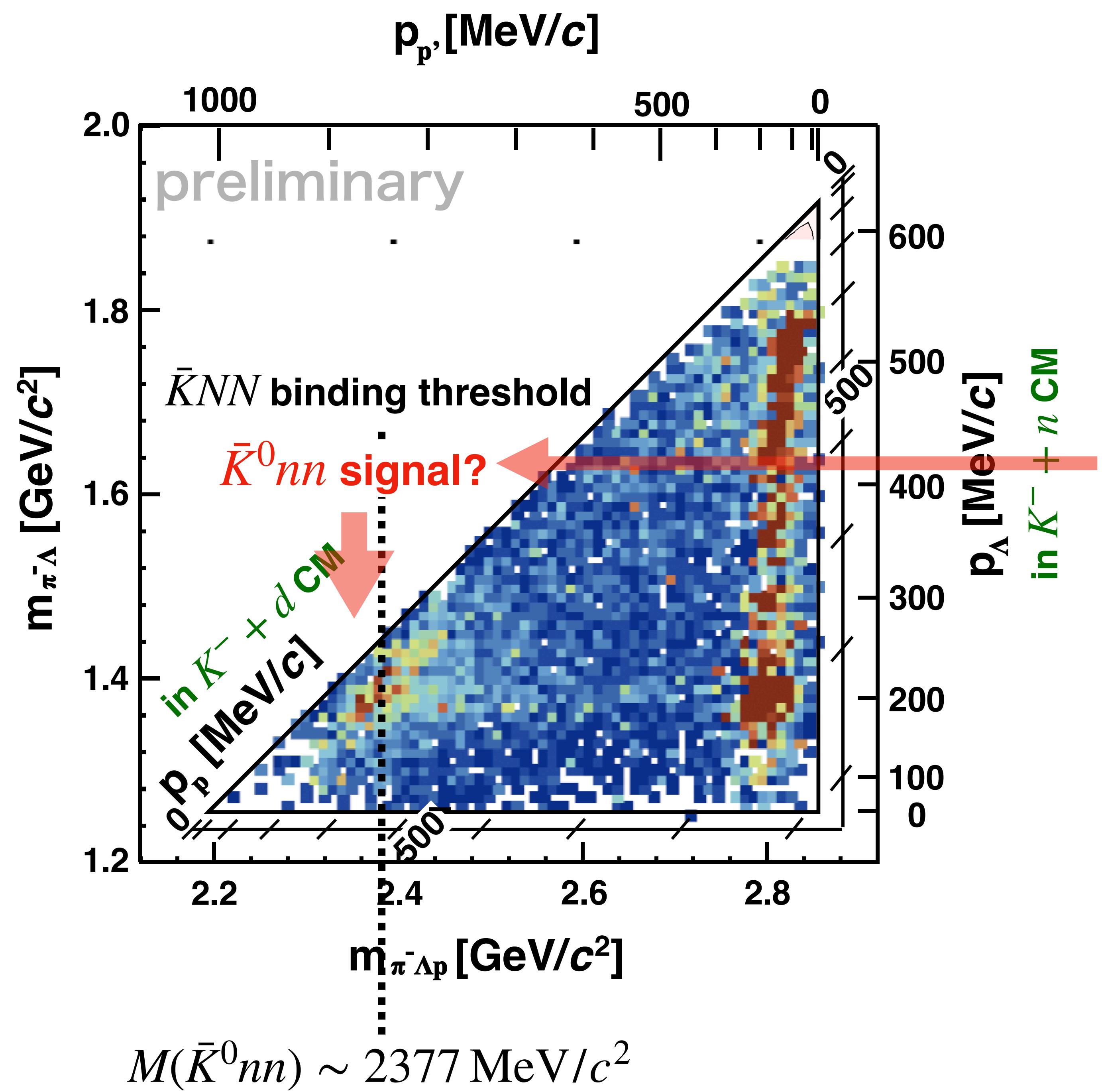


... analyzed by T. Yamaga



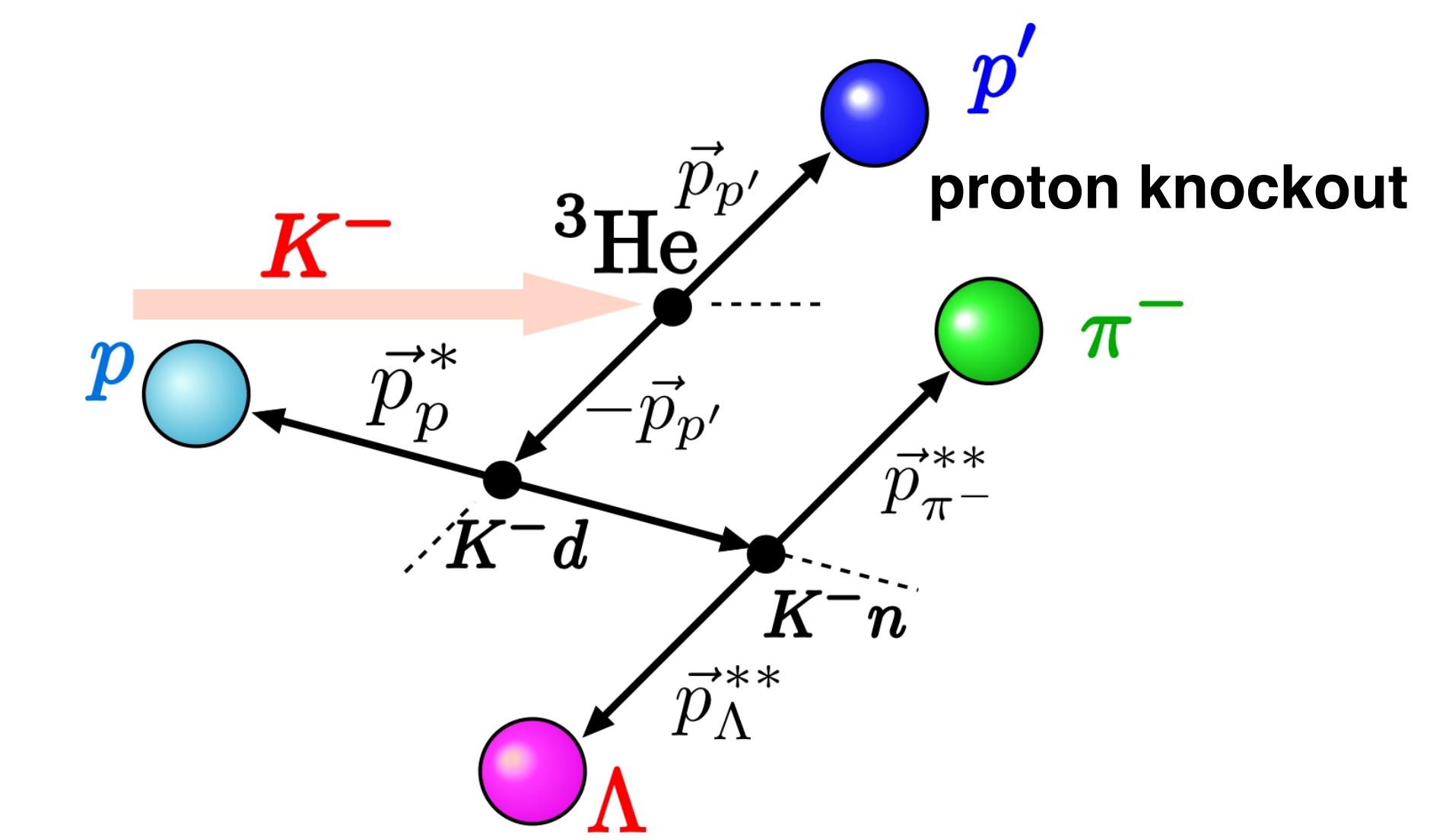
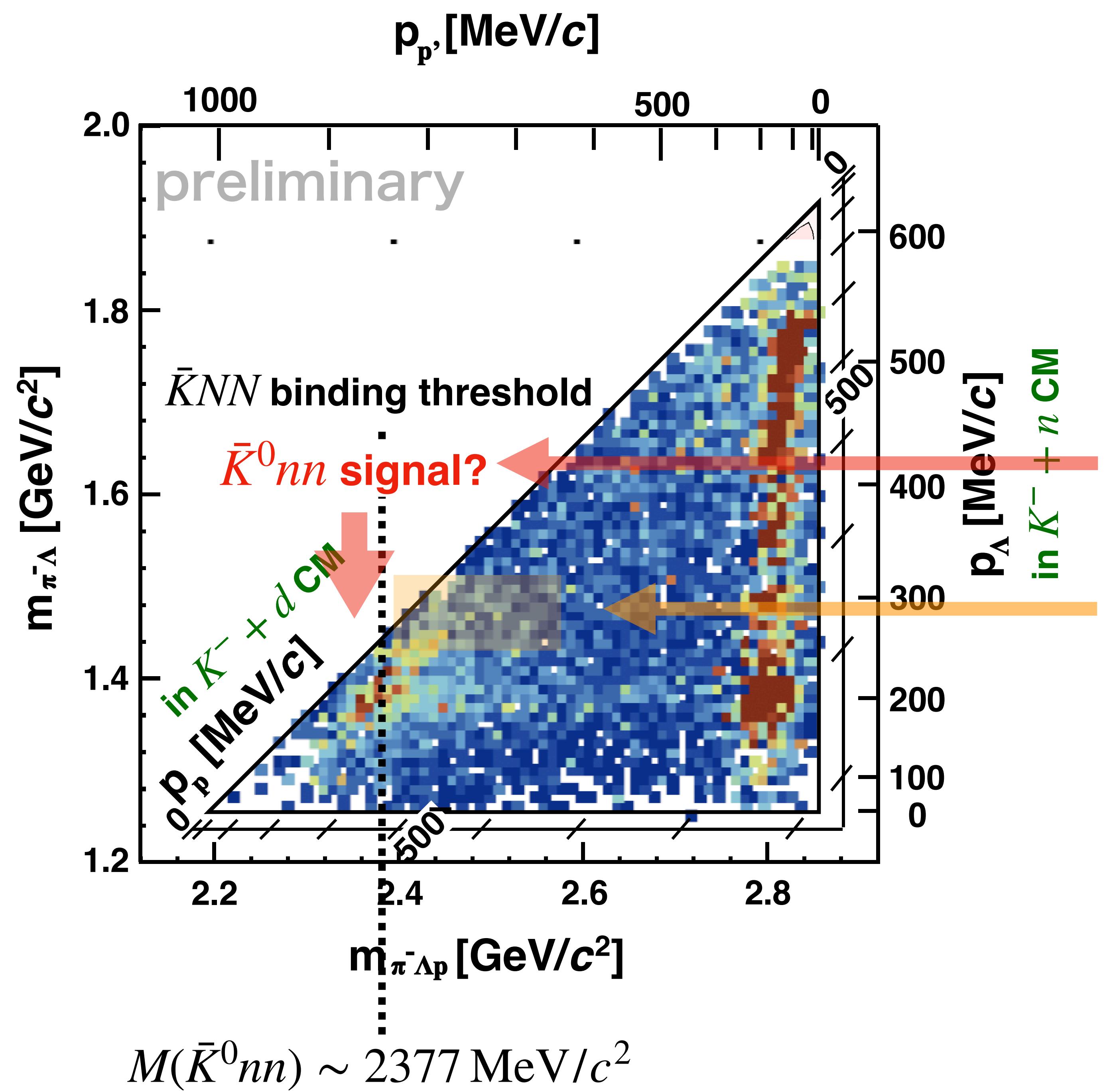
$$M(\bar{K}^0 nn) \sim 2377 \text{ MeV}/c^2$$



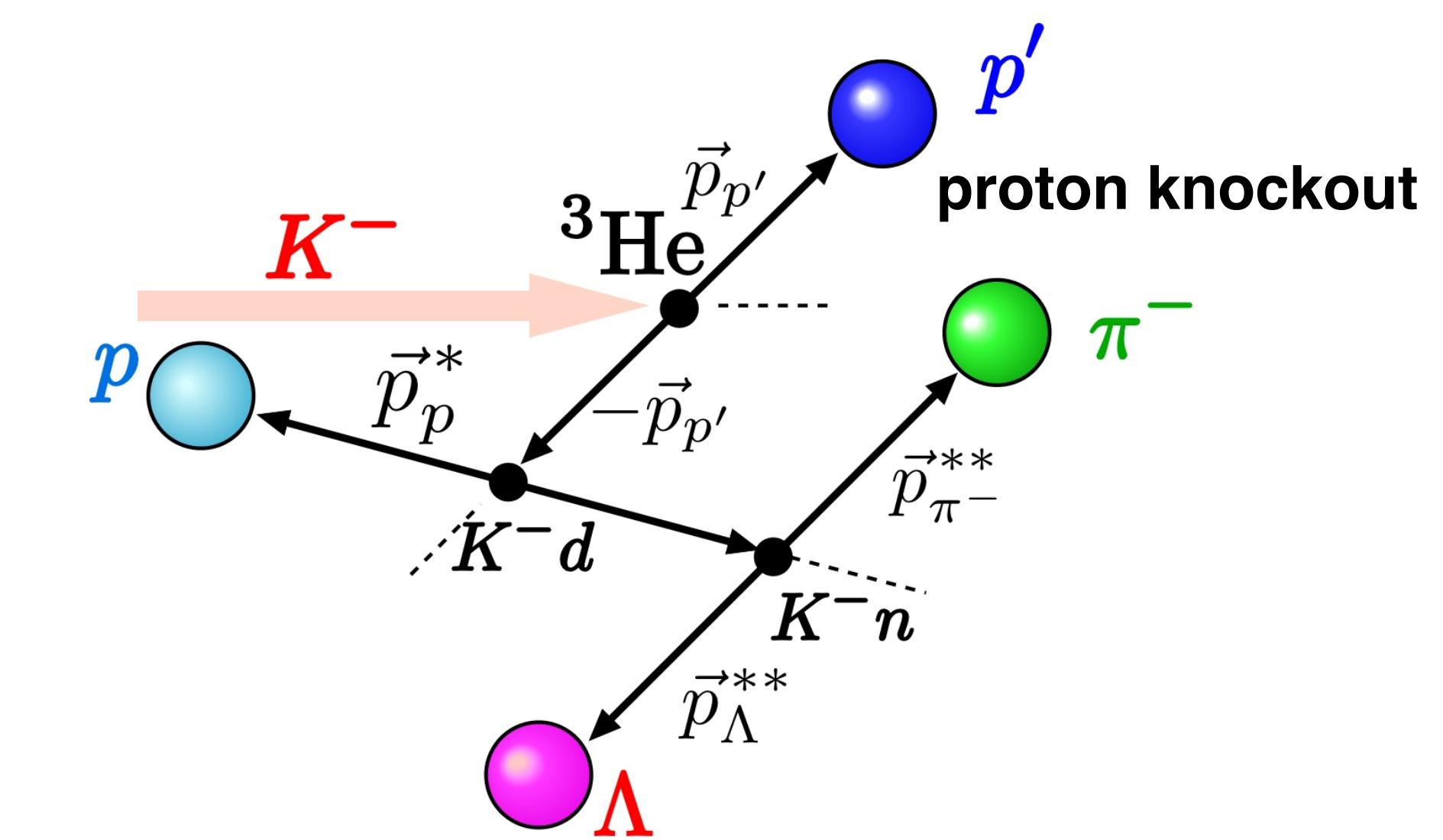
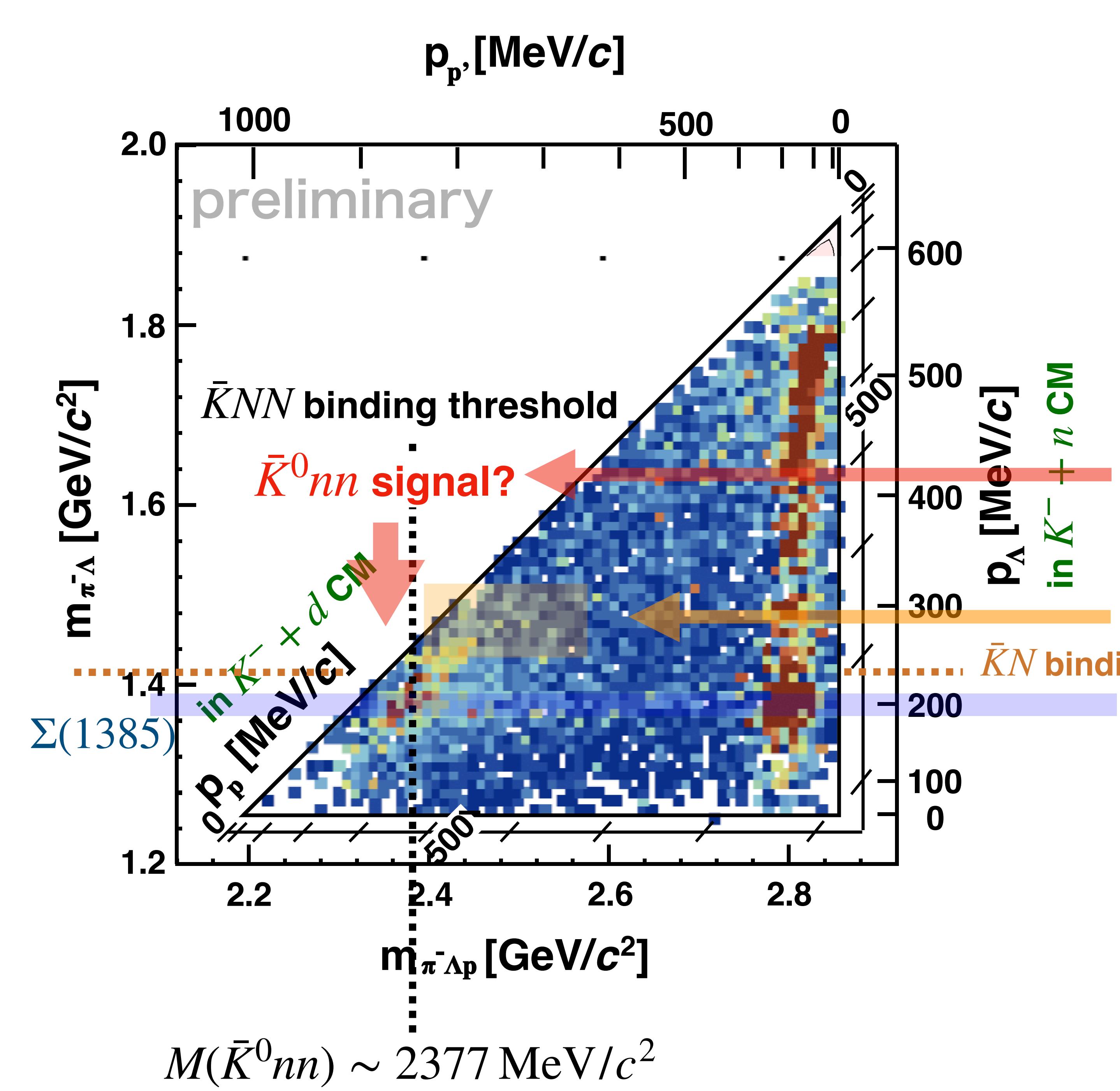


$\bar{K}^0 nn$ signal-like event concentration
below \bar{K} -bound threshold is seen?

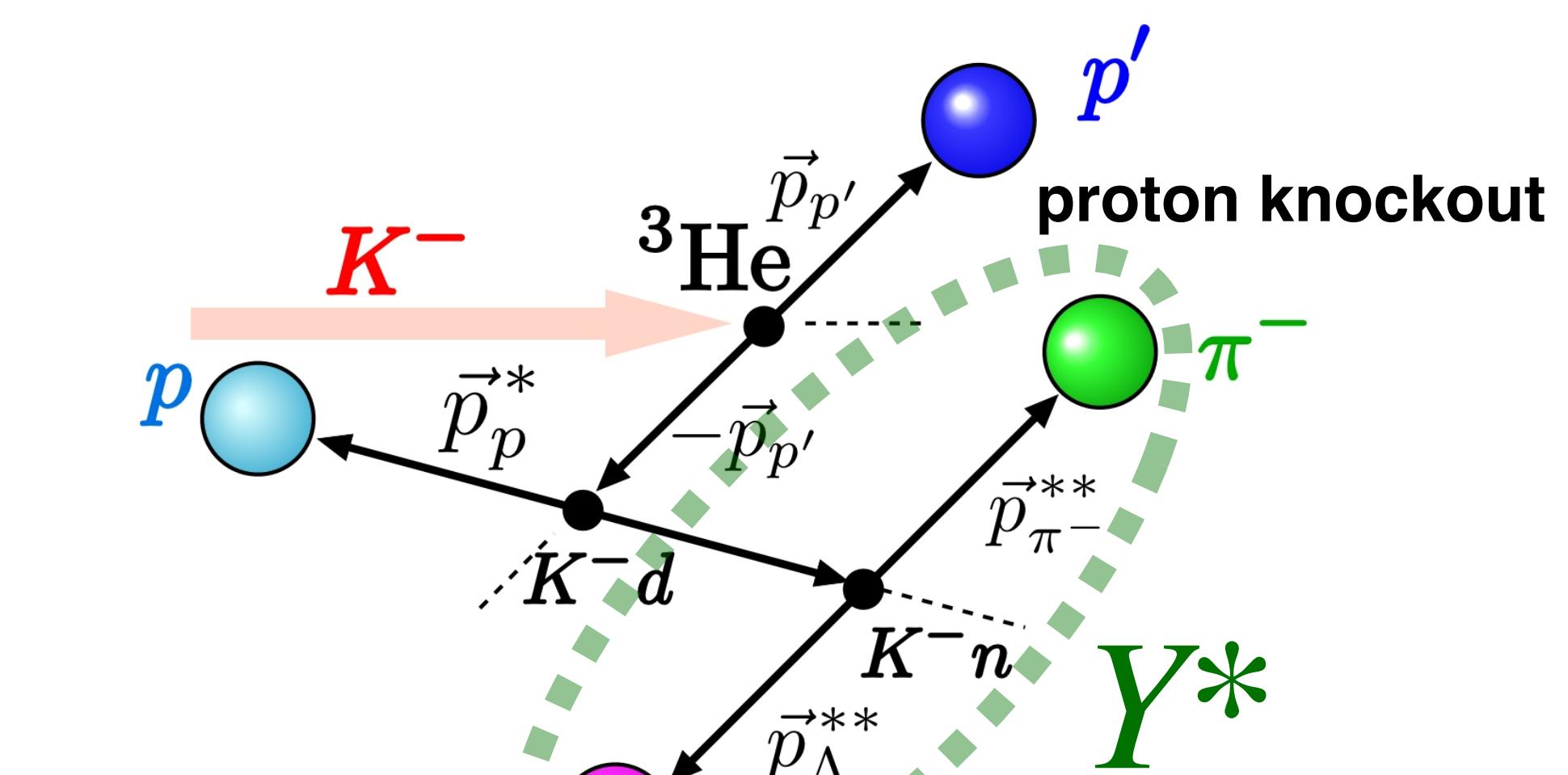
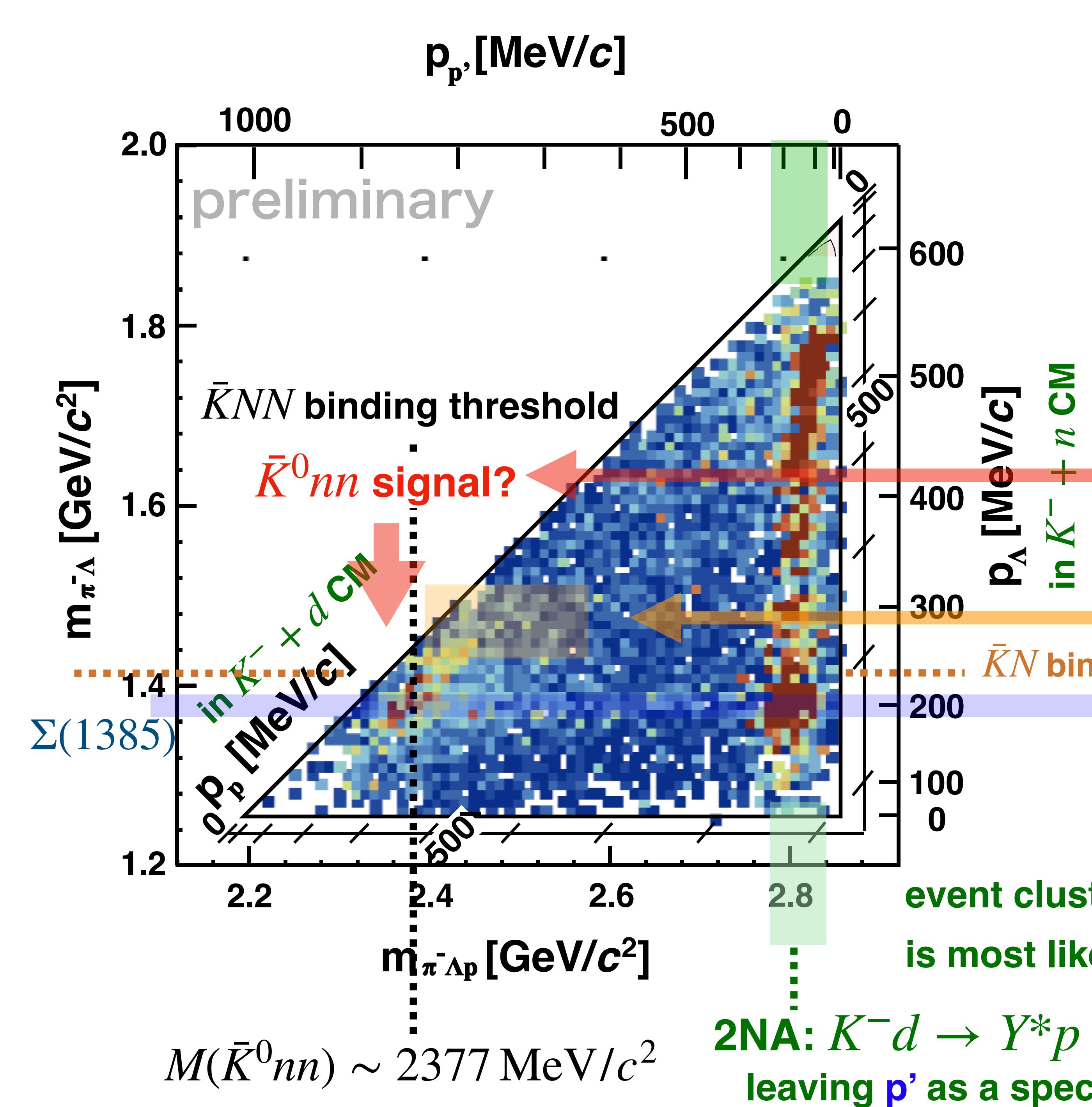
— twice more data become available in April —



$\bar{K}^0 nn$ signal-like event concentration
below \bar{K} -bound threshold is seen?
— twice more data become available in April —
QF-K induced reaction?



$\bar{K}^0 nn$ signal-like event concentration
below \bar{K} -bound threshold is seen?
— twice more data become available in April —
QF-K induced reaction?



$\bar{K}^0 nn$ signal-like event concentration below \bar{K} -bound threshold is seen?

— twice more data become available in April —

QF-K induced reaction?

$\Sigma(1385)$ contribution is not negligible compared to $(\Lambda p) + n$ final state.

event cluster at $m_{\pi^-\Lambda p} \sim \sqrt{s_{K-d}} \approx 2.83 \text{ GeV}$
is most likely ...

2NA: $K^-d \rightarrow Y^*p$ reaction leaving p' as a spectator

- K^- seems to be sensitive to the deuteron cluster in ${}^3\text{He}$ –